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## ON A GAMBLER'S RUIN PROBLEM

Mohamed A. El-Shehawey\* — B. N. Al-Matrafi\*\*

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ABSTRACT. The object of this paper is to specify an explicit expression for the absorption probabilities for a random walk on the integers  $-b, -b+1, \ldots, -1, 0, 1, \ldots, a$  which arises in a gambler's ruin problem proposed by the authors. There are two barriers, one absorbing at a and the other at -b, such that the random walk particle is returned to the position  $j, -b \leq j < a$ , whenever reaches it.

### 1. Introduction

The classical problem of the gambler's ruin affords a classical illustration of the simple random walk with absorbing barriers. For example, B a r n e t t (1964), Miller (1965),Feller (1968),Srinivasan Cox and and Mehata (1976), and Kannan (1979) have treated this problem. We consider here a more general problem where there are two barriers, one of which is absorbing and the other is such that the random walk particle is immediately returned to a certain position upon reaching it. We consider two players A and B with initial fortunes a and b dollars respectively. The game consists of a series of independent turns. Let  $X_u$ ,  $u \ge 1$ , denote the B's winnings in the uth turn with

$$X_u = \begin{cases} 1 & \text{with probability } p \,, \\ -1 & \text{with probability } q \,, \\ 0 & \text{with probability } r \,, \end{cases}$$

where p + q + r = 1 and 0 < p, q < 1,  $0 \le r < 1$ . Thus, we have

$$S_n = \sum_{u=1}^n X_u$$

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which represents B's net gain at the end of n turns. If at any stage  $S_n = -b$ , B is ruined, player A immediately donates j dollars of his fortune to player B,  $0 \leq j < a$ . The game ends when B has gained all A's fortune and A is ruined, i.e.,  $S_n = a$ . Thus  $S_n$  is a random walk with two barriers, one of which, at a, is absorbing, and the other, at -b, such that the random walk particle is returned to the position j,  $-b \leq j < a$ , whenever reaching it. In this paper, an explicit expression for the absorption probabilities is derived. Cox and Miller [3; p. 33] (1965) solved the classical random walk problem with absorbing barriers at the points -b and a; however, the probability of absorption at time n is found to be erroneous. The correct formula can be found in Gulati and Hill [6] (1981). The special case, when the barriers are symmetrically placed on either side of the origin has been given by Munford [6] (1981). If we replace the absorbing barrier at -b by an impenetrable (reflecting) barrier at j = 0, -1's are never accumulated if  $S_n$  reaches zero. The absorption probabilities have been deduced by El-Shehawy [4] (1992).

### 2. Generating function for the absorption probabilities

Let  $f(n, a \mid i_0)$  denote the probability that the particle is absorbed at a at time n given that its initial position was  $i_0, -b \leq i_0 \leq a$ . The absorption probabilities  $f(n, a \mid i_0)$  must satisfy

$$\begin{aligned} f(n,a \mid i_0) &= pf(n-1,a \mid i_0+1) + rf(n-1,a \mid i_0) + qf(n-1,a \mid i_0-1) \\ (n &= 1,2,\ldots; \ i_0 &= -b+1,\ldots,-1,0,1,\ldots,a-1), \end{aligned}$$

where

$$\begin{split} &f(0,a \mid i_0) = 0 \quad (i_0 = -b, \dots, -1, 0, 1, \dots, a-1) \,, \\ &f(0,a \mid a) = 1 \,, \qquad f(n,-b \mid i_0) = f(n,j \mid i_0) \quad (n=1,2,\dots) \,. \end{split}$$

Introducing the generating function

$$G_{i_0,a}(t) = \sum_{n=0}^{\infty} f(n,a \mid i_0) t^n \,.$$
<sup>(2)</sup>

Following Cox and Miller [3] (1965) we deduce that

$$G_{i_0,a}(t) = H_{i_0}(t)/H_a(t), \qquad i_0 = -b, \dots, -1, 0, 1, \dots, a,$$
(3)

where

$$H_{i_0}(t) = \left(h_1(t)\right)^{b+i_0} - \left(h_2(t)\right)^{b+i_0} + [q/p]^{b+i_0} \left[\left(h_1(t)\right)^{j-i_0} - \left(h_2(t)\right)^{j-i_0}\right],$$

and  $h_1(t)$ ,  $h_2(t)$  are the roots of the quadratic equation

$$pth^{2} - (1 - rt)h + qt = 0, \qquad (4)$$

p, q and r are the probabilities of taking one step to the right, to the left and remaining in position, respectively. Writing X(t) = 1 - rt and  $Y(t) = \sqrt{4pqt^2 - (1-rt)^2}$ , we have

$$h_{1,2}(t) = (2pt)^{-1} [X(t) \pm i Y(t)], \quad i = \sqrt{-1},$$

 $\operatorname{and}$ 

$$H_{i_0}(t) = (2pt)^{-(b+i_0)} \left[ \left( X(t) + iY(t) \right)^{b+i_0} - \left( X(t) - iY(t) \right)^{b+i_0} \right] + (2pt)^{i_0 - j} [q/p]^{b+i_0} \left[ \left( X(t) + iY(t) \right)^{j-i_0} - \left( X(t) - iY(t) \right)^{j-i_0} \right].$$
(5)

Accordingly, by expanding both the numerator and denominator of (3) in powers of Y(t) and noting that only odd powers of Y(t) occur, thus Y(t) is cancelled, leaving a ratio of two polynomials in t, each of degree at most a+b-1, and consequently a partial fraction expansion of it is available. Using complex variables notation, we have

$$X(t) = |z| \cos \phi$$
,  $Y(t) = |z| \sin \phi$ ,

where

$$|z| = \sqrt{X^2(t) + Y^2(t)} = 2\sqrt{pq}t$$
, and  $\phi = tan^{-1}(Y(t)/X(t))$ .

It is useful to observe that

$$1 - rt = 2\sqrt{pq}t\cos\phi, \quad \text{or} \quad t = \left(r + 2\sqrt{pq}\cos\phi\right)^{-1}, \tag{6}$$

and  $h_{1,2}(t) = \sqrt{q/p} e^{\pm i \phi}$ . Formula (3) becomes

$$G_{i_0,a}(t) = T_{i_0}(\phi) / T_a(\phi),$$
(7)

where

$$T_{i_0}(\phi) = \left[\sqrt{p/q}\right]^{a-i_0} \left[ \left(\sqrt{p}\right)^{b+j} \sin(b+i_0)\phi + \left(\sqrt{q}\right)^{b+j} \sin(j-i_0)\phi \right].$$

The denominator of (7) is found to have a+b-1 distinct roots. A study of the function  $\gamma(\phi)$ , where

$$\gamma(\phi) = \frac{\sin(a+b)\phi}{\sin(a-j)\phi} - \left[\sqrt{q/p}\right]^{b+j}$$
(8)

shows that if  $a - j = \left[\sqrt{p/q}\right]^{b+j}(b+a)$ , the roots of the equation  $\gamma(\phi) = 0$  give distinct roots of  $T_a(\phi) \in [0, \pi)$ . If  $a - j < \left[\sqrt{p/q}\right]^{b+j}(b+a)$ , there are a + b - 1

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distinct real roots  $\phi_{\nu}$  ( $\nu = 1, 2, ..., a + b - 1$ ) of (8). The corresponding roots of  $H_a(t)$  are then

$$t_{\nu} = \frac{1}{r + 2\sqrt{pq}\cos\phi_{\nu}}, \qquad \nu = 1, 2, \dots, a + b - 1.$$

If  $a-j > \left[\sqrt{p/q}\right]^{b+j}(b+a)$ , there are only a+b-2 distinct roots  $\phi_{\nu}$  ( $\nu = 2, 3, \ldots, a+b-1$ ) that give distinct roots  $t_{\nu}$  of  $H_a(t)$ . The remaining root of  $H_a(t)$  is given by

$$t_1 = \frac{1}{r + 2\sqrt{pq}\cosh\phi_1}\,,$$

where  $\phi_1$  is the unique root of the equation,

$$\sinh(a-j)\phi = \left[\sqrt{p/q}\right]^{b+j}\sinh(a+b)\phi$$
.

## 3. Explicit expression for the absorption probabilities

THEOREM. We have

$$f(n, a \mid i_{0}) = -2\sqrt{pq} \left[ M^{0}(\phi_{1}) + \sum_{\nu=2}^{a+b-1} T_{i_{0}}(\phi_{\nu}) \left( \frac{\partial T_{a}(\phi)}{\partial \phi} \big|_{\phi=\phi_{\nu}} \right)^{-1} \cdot \left( r + 2\sqrt{pq} \cos \phi_{\nu} \right)^{n-1} \sin \phi_{\nu} \right],$$
(9)

where

$$M^{0}(\phi_{1}) = M(\phi_{1}) \cdot \begin{cases} \left[r + 2\sqrt{pq}\cos\phi_{1}\right]^{n+1}, & a - j < \left(\sqrt{p/q}\right)^{b+j}(a+b), \\ \left[1 - \left(\sqrt{p} - \sqrt{q}\right)^{2}\right]^{n+1}, & a - j = \left(\sqrt{p/q}\right)^{b+j}(a+b), \\ \left[r + 2\sqrt{pq}\cosh\phi_{1}\right]^{n+1}, & a - j > \left(\sqrt{p/q}\right)^{b+j}(a+b), \end{cases}$$

and

$$M(\phi_1) = \begin{cases} T_{i_0}(\phi_1) \left[ \frac{\partial T_a(\phi)}{\partial \phi} \big|_{\phi=\phi_1} \right]^{-1} \left[ r + 2\sqrt{pq} \cos \phi_1 \right]^{-2} \sin \phi_1 \,, \\ a - j < \left( \sqrt{p/q} \right)^{b+j} (a+b) \,, \\ -2(a - i_0)(a+b)^{-1}(a-j)^{-1}(2a+b-j)^{-1} \left[ r + 2\sqrt{pq} \right]^{-2} \,, \\ a - j = \left( \sqrt{p/q} \right)^{b+j} (a+b) \,, \\ -T_{i_0}(\mathrm{i} \, \phi_1) \left[ \frac{\partial T_a(\phi)}{\partial \phi} \big|_{\phi=\mathrm{i} \, \phi_1} \right]^{-1} \left[ r + 2\sqrt{pq} \cosh \phi_1 \right]^{-2} \sinh \phi_1 \,, \\ a - j > \left( \sqrt{p/q} \right)^{b+j} (a+b) \,. \end{cases}$$

P r o o f. Employing the partial fraction method, formula (3) can be written in the form

$$G_{i_0,a}(t) = \sum_{n=0}^{\infty} \left( \sum_{\nu=1}^{a+b-1} \frac{\eta_{\nu}}{t_{\nu+1}^{n+1}} \right) t^n ,$$

where the constants  $\eta_{\nu}$  are given by,

$$\eta_{\nu} = -H_{i_0}(t) \left[ \frac{\partial H_a(t)}{\partial t} \right]^{-1} \Big|_{t=t_{\nu}}.$$
(10)

Differentiating the quadratic equation (4) with respect to t we get

$$\frac{\partial h_{1,2}(t)}{\partial t} = t^{-2} h_{1,2}^2(t) \left[ q - p h_{1,2}^2(t) \right]^{-1}.$$
(11)

The study of (8) with (3), (10) and (11) finally yield

$$\eta_{\nu} = \frac{-2\sqrt{pq} \left[\sin(b+i_0)\phi_{\nu} + \left(\sqrt{q/p}\right)^{b+j} \sin(j-i_0)\phi_{\nu}\right] t_{\nu}^2 \sin \phi_{\nu}}{\left(\sqrt{q/p}\right)^{a-i_0} \left[(a+b)\cos(a+b)\phi_{\nu} + \left(\sqrt{q/p}\right)^{b+j}(j-a)\cos(j-a)\phi_{\nu}\right]}, \\ \nu = 2, 3, \dots, a+b-1$$
(12)

and

$$\eta_1 = -2\sqrt{pq}M(\phi_1)\,,$$

where

$$\phi_1 = \begin{cases} \cos^{-1} \frac{1 - rt_1}{2\sqrt{pq}t_1} & \text{if } a - j \le \left(\sqrt{p/q}\right)^{b+j}(a+b) \,,\\ \cosh^{-1} \frac{1 - rt_1}{2\sqrt{pq}t_1} & \text{if } a - j > \left(\sqrt{p/q}\right)^{b+j}(a+b) \,, \end{cases}$$

 $t_1$  is the smallest root in absolute value of  $H_a(t)$ . Formula (9) follows.

We see that with the appropriate change of notation, expression (9) agrees with that of Munford [8] (1981) in the case j = 0, p > q; and with that of Weesakul's [10] (1961) and Blasi [2] (1976) for the particular case j = 0, b = 1, r = 0, replacing  $i_0$ , a,  $\phi_{\nu}$ ,  $\eta_{\nu}$  with b - u, b,  $\alpha_{\nu}$ ,  $\rho_{\nu}$  and interchanging p and q respectively.

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