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INSTANTON-ANTI-INSTANTON SOLUTIONS OF DISCRETE YANG-MILLS EQUATIONS

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Abstract. We study a discrete model of the SU(2) Yang-Mills equations on a combinatorial analog of \mathbb{R}^4 . Self-dual and anti-self-dual solutions of discrete Yang-Mills equations are constructed. To obtain these solutions we use both the techniques of a double complex and the quaternionic approach.

Keywords: Yang-Mills equations, self-dual equations, anti-self-dual equations, instanton, anti-instanton, difference equations

MSC 2010: 81T13, 39A12

1. INTRODUCTION

We study an intrinsically defined discrete model of the SU(2) Yang-Mills equations on a combinatorial analog of \mathbb{R}^4 . It is known (see, for example, [5]) that a gauge potential can be defined as a certain su(2)-valued 1-form A (the connection 1-form). Then the gauge field F (the curvature 2-form) is given by

(1.1)
$$F = \mathrm{d}A + A \wedge A,$$

where \wedge denotes the exterior multiplication. The Yang-Mills equations can be expressed in terms of the 2-forms F and *F as

(1.2)
$$dF + A \wedge F - F \wedge A = 0, \quad d*F + A \wedge *F - *F \wedge A = 0,$$

where * is the Hodge star operator.

We consider the self-dual and anti-self-dual equations

(1.3)
$$F = *F, \quad F = -*F.$$

Equations (1.3) are nonlinear matrix first order partial differential equations. In the 4-dimensional Yang-Mills theories the self-dual (instanton) and anti-self-dual (antiinstanton) solutions of (1.3) are the absolute minima of the Yang-Mills action and satisfy the second-order Yang-Mills equations (1.2) (see [4]).

The purpose of this paper is to construct the self-dual and anti-self-dual solutions of discrete SU(2) Yang-Mills equations which imitate the corresponding solutions of the continual theory. The ideas presented here are strongly influenced by the book of Dezin [2]. We develop discrete models of some objects in differential geometry, including the Hodge star operator, the differential and the exterior multiplication, in such a way that they preserve the geometric structure of their continual analogs. We continue the investigations which were originated in [3], [6]–[8]. The geometrical discretisation techniques used here extend those introduced in [2] and [6]. A combinatorial model of \mathbb{R}^4 based on the use of the double complex construction is taken from [8].

2. Quaternions and the SU(2)-connection

We begin with a brief review of some preliminaries about quaternions. The quaternions are formed from real numbers by adjoining three symbols $\mathbf{i}, \mathbf{j}, \mathbf{k}$, and an arbitrary quaternion x can be written as

(2.1)
$$x = x_1 + x_2 \mathbf{i} + x_3 \mathbf{j} + x_4 \mathbf{k},$$

where $x_1, x_2, x_3, x_4 \in \mathbb{R}$. The symbols $\mathbf{i}, \mathbf{j}, \mathbf{k}$ satisfy the identities

(2.2)
$$\mathbf{i}^2 = \mathbf{j}^2 = \mathbf{k}^2 = -1,$$
$$\mathbf{i}\mathbf{j} = -\mathbf{j}\mathbf{i} = \mathbf{k}, \quad \mathbf{j}\mathbf{k} = -\mathbf{k}\mathbf{j} = \mathbf{i}, \quad \mathbf{k}\mathbf{i} = -\mathbf{i}\mathbf{k} = \mathbf{j}.$$

It is clear that the space of quaternions is isomorphic to \mathbb{R}^4 . By analogy with the complex numbers, x_1 is called the real part of x and $x_2\mathbf{i} + x_3\mathbf{j} + x_4\mathbf{k}$ is called the imaginary part. In the sequel we will write

$$\operatorname{Im} x = x_2 \mathbf{i} + x_3 \mathbf{j} + x_4 \mathbf{k}.$$

The conjugate quaternion of x is defined by

$$\bar{x} = x_1 - x_2 \mathbf{i} - x_3 \mathbf{j} - x_4 \mathbf{k}$$

Then the norm |x| of a quaternion can be introduced as

(2.3)
$$|x|^2 = x\bar{x} = x_1^2 + x_2^2 + x_3^2 + x_4^2$$

The algebra of quaternions can be represented as a sub-algebra of the 2×2 complex matrices $M(2,\mathbb{C})$. We identify the quaternion (2.1) with a matrix $f(x) \in M(2,\mathbb{C})$ by setting

(2.4)
$$f(x) = \begin{pmatrix} x_1 + x_2 \mathbf{i} & x_3 + x_4 \mathbf{i} \\ -x_3 + x_4 \mathbf{i} & x_1 - x_2 \mathbf{i} \end{pmatrix}.$$

Here i is the imaginary unit.

It is well known that the unit quaternions, i.e., those that have the norm |x| = 1, form a group and this group is isomorphic to SU(2). The 2 × 2 complex matrices

(2.5)
$$\mathbf{i} = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix}, \quad \mathbf{j} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad \mathbf{k} = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$

realize a representation of the Lie algebra su(2) of the group SU(2). Note that multiplying by -i these tree matrices we obtain the standard Pauli matrices. Matrices (2.5) correspond to the units $\mathbf{i}, \mathbf{j}, \mathbf{k}$ given by (2.2). Thus the Lie algebra su(2) can be viewed as the pure imaginary quaternions with the basis $\mathbf{i}, \mathbf{j}, \mathbf{k}$.

Let the SU(2)-connection A be given by

(2.6)
$$A = \sum_{\mu} A_{\mu}(x) \,\mathrm{d}x^{\mu},$$

where $A_{\mu}(x) \in su(2)$ and $x = (x_1, \ldots, x_4)$ is a point of \mathbb{R}^4 . On the other hand, A can be defined also as taking values in the space of pure imaginary quaternions. Let f(x) be a function of the quaternion variable (2.1) with quaternion values. Then we can write A as

(2.7)
$$A = \operatorname{Im}(f(x) \,\mathrm{d}x),$$

where $f(x) = f_1(x) + f_2(x)\mathbf{i} + f_3(x)\mathbf{j} + f_4(x)\mathbf{k}$ and $dx = dx_1 + dx_2\mathbf{i} + dx_3\mathbf{j} + dx_4\mathbf{k}$. Using the rules of multiplication (2.2) we have

$$A_1(x) = f_2(x)\mathbf{i} + f_3(x)\mathbf{j} + f_4(x)\mathbf{k}, \quad A_2(x) = f_1(x)\mathbf{i} + f_4(x)\mathbf{j} - f_3(x)\mathbf{k},$$

$$A_3(x) = -f_4(x)\mathbf{i} + f_1(x)\mathbf{j} + f_2(x)\mathbf{k}, \quad A_4(x) = f_3(x)\mathbf{i} - f_2(x)\mathbf{j} + f_1(x)\mathbf{k}.$$

Using (2.7) we can rewrite (1.1) as

(2.8)
$$F = \operatorname{Im}(\mathrm{d}f(x) \wedge \mathrm{d}x + f(x)\,\mathrm{d}x \wedge f(x)\,\mathrm{d}x).$$

In the quaternion notation the instanton and anti-instanton solutions can be found in Atiyah [1]. In Section 4 we will construct discrete analogs of these solutions.

3. Discrete model

We will use the double complex construction described in [8]. Let the tensor product $C(4) = C \otimes C \otimes C \otimes C$ of a 1-dimensional complex C be a combinatorial model of the Euclidean space \mathbb{R}^4 (for details see also [2]). The 1-dimensional complex C is defined in the following way. Let C^0 denote the real linear space of 0-dimensional chains generated by basis elements x_j (points), $j \in \mathbb{Z}$. It is convenient to introduce the shift operators τ, σ in the set of indices by

(3.1)
$$\tau j = j + 1, \quad \sigma j = j - 1.$$

We denote the open interval $(x_j, x_{\tau j})$ by e_j . We will regard the set $\{e_j\}$ as a set of basis elements of the real linear space C^1 of 1-dimensional chains. Then the 1-dimensional complex (combinatorial real line) is the direct sum of the spaces introduced above: $C = C^0 \oplus C^1$. Together with the complex C(4) we consider its double, namely, the complex $\widetilde{C}(4)$ of exactly the same structure (for details see [8]). We need the double to define a discrete analog of the Hodge star operator.

Let K(4) be a cochain complex with $gl(2, \mathbb{C})$ -valued coefficients, where $gl(2, \mathbb{C})$ is the Lie algebra of the group $GL(2, \mathbb{C})$. Recall that $gl(2, \mathbb{C})$ consists of all complex 2×2 matrices $M(2, \mathbb{C})$ with bracket operation $[\cdot, \cdot]$. The complex K(4) is a conjugate of C(4) and we have $K(4) = K \otimes K \otimes K \otimes K$, where K is a conjugate of the 1dimensional complex C. Basis elements of K can be written as x^j , e^j . Then an arbitrary p-dimensional basis element of K(4) is given by $s_{(p)}^k = s^{k_1} \otimes s^{k_2} \otimes s^{k_3} \otimes$ s^{k_4} , where s^{k_i} is either x^{k_i} or e^{k_i} , $k_i \in \mathbb{Z}$. Note that $s_{(p)}^k$ contains exactly p of 1-dimensional elements e^{k_i} . For a p-dimensional cochain $\varphi \in K(4)$ we have

(3.2)
$$\varphi = \sum_{k} \sum_{p} \varphi_{k}^{(p)} s_{(p)}^{k},$$

where $\varphi_k^{(p)} \in gl(2,\mathbb{C})$. We will call cochains forms, emphasizing their relationship with the corresponding continual objects, differential forms. Denote by $\widetilde{K}(4)$ the complex of cochains over the double complex $\widetilde{C}(4)$. It is clear that $\widetilde{K}(4)$ has the same structure as K(4). Let us introduce the operation $\tilde{\iota}: K(4) \to \widetilde{K}(4), \tilde{\iota}: \widetilde{K}(4) \to K(4)$ by setting

(3.3)
$$\tilde{\iota}s_{(p)}^k = \tilde{s}_{(p)}^k, \qquad \tilde{\iota}\tilde{s}_{(p)}^k = s_{(p)}^k,$$

where $s_{(p)}^k$ and $\tilde{s}_{(p)}^k$ are basis elements of K(4) and $\tilde{K}(4)$. Hence for a *p*-form $\varphi \in K(4)$ we have $\tilde{\iota}\varphi = \tilde{\varphi}$.

For the definitions of d^c , \cup and * on K(4), which are discrete analogs of the differential d, exterior multiplication \wedge and the Hodge star operator respectively, we refer the reader to [8].

Let us consider a discrete 0-form with coefficients belonging to $M(2,\mathbb{C})$. We put

$$(3.4) f = \sum_{k} f_k x^k$$

where $x^k = x^{k_1} \otimes x^{k_2} \otimes x^{k_3} \otimes x^{k_4}$ is the 0-dimensional basis element of K(4). Suppose that the matrices $f_k \in M(2, \mathbb{C})$ look like (2.4). Then f_k in quaternionic form can be expressed as

(3.5)
$$f_k = f_k^1 + f_k^2 \mathbf{i} + f_k^3 \mathbf{j} + f_k^4 \mathbf{k}.$$

Hence the form (3.4) can be viewed as a discrete form with quaternionic coefficients. We will call it simply the quaternionic form when no confusion can arise.

Let us denote by e the quaternionic 1-form

(3.6)
$$e = \sum_{k} e^{k} = \sum_{k} (e_{1}^{k} + e_{2}^{k} \mathbf{i} + e_{3}^{k} \mathbf{j} + e_{4}^{k} \mathbf{k}),$$

where e_i^k are the 1-dimensional basis elements of K(4). Let $A \in K(4)$ be a discrete 1-form. We define the discrete SU(2)-connection A (discrete analog of (2.6)) to be

(3.7)
$$A = \sum_{k} \sum_{i=1}^{4} A_k^i e_i^k,$$

where $A_k^i \in su(2)$. Using (3.4) and (3.6), we write (3.7) in the quaternionic form as

(3.8)
$$A = \operatorname{Im}(f \cup e) = \operatorname{Im}\left(\sum_{k} f_{k} e^{k}\right).$$

Then the A_k^i are given by

(3.9)
$$A_{k}^{1} = f_{k}^{2}\mathbf{i} + f_{k}^{3}\mathbf{j} + f_{k}^{4}\mathbf{k}, \qquad A_{k}^{2} = f_{k}^{1}\mathbf{i} + f_{k}^{4}\mathbf{j} - f_{k}^{3}\mathbf{k}, A_{k}^{3} = -f_{k}^{4}\mathbf{i} + f_{k}^{1}\mathbf{j} + f_{k}^{2}\mathbf{k}, \qquad A_{k}^{4} = f_{k}^{3}\mathbf{i} - f_{k}^{2}\mathbf{j} + f_{k}^{1}\mathbf{k}.$$

An arbitrary discrete 2-form $F \in K(4)$ can be written as

(3.10)
$$F = \sum_{k} \sum_{i < j} F_k^{ij} \varepsilon_{ij}^k,$$

where $F_k^{ij} \in gl(2, \mathbb{C}), 1 \leq i, j \leq 4$, and ε_{ij}^k is the 2-dimensional basis element of K(4). Let F be given by

$$(3.11) F = d^c A + A \cup A.$$

For convenience we also introduce the shift operator τ_i which acts in the set of indices as $\tau_i k = (k_1, \ldots, \tau k_i, \ldots, k_4)$, where τ is given by (3.1).

By the definitions of d^c and \cup , combining (3.7) and (3.11), we obtain

(3.12)
$$F_{k}^{ij} = \Delta_{i}A_{k}^{j} - \Delta_{j}A_{k}^{i} + A_{k}^{i}A_{\tau_{i}k}^{j} - A_{k}^{j}A_{\tau_{j}k}^{i}$$

where $\Delta_i A_k^j = A_{\tau_i k}^j - A_k^j$.

It should be noted that in the continual case the curvature form F (1.1) takes values in the algebra su(2) for any su(2)-valued connection form A. Unfortunately, this is not true in the discrete case because, generally speaking, the components $A_k^i A_{\tau_i k}^j - A_k^j A_{\tau_i k}^i$ of the form $A \cup A$ (see (3.12)) do not belong to su(2).

To define an su(2)-valued discrete analog of the curvature 2-form we use the quaternionic form of A (3.8) and put it in (3.11). Then the discrete curvature form F is given by

(3.13)
$$F = \operatorname{Im} \{ \mathrm{d}^{c} f \cup e + (f \cup e) \cup (f \cup e) \}.$$

Putting (3.9) in (3.12) we find that

$$\begin{split} F_k^{12} &= (\Delta_1 f_k^1 - \Delta_2 f_k^2 - f_k^3 f_{\tau_1k}^3 - f_k^4 f_{\tau_1k}^4 - f_k^3 f_{\tau_2k}^3 - f_k^4 f_{\tau_2k}^4) \mathbf{i} \\ &+ (\Delta_1 f_k^4 - \Delta_2 f_k^3 + f_k^2 f_{\tau_1k}^3 + f_k^4 f_{\tau_1k}^1 + f_k^1 f_{\tau_2k}^4 + f_k^3 f_{\tau_2k}^2) \mathbf{j} \\ &+ (-\Delta_1 f_k^3 - \Delta_2 f_k^4 + f_k^2 f_{\tau_1k}^4 - f_k^3 f_{\tau_1k}^1 - f_k^1 f_{\tau_2k}^3 + f_k^4 f_{\tau_2k}^2) \mathbf{k} \\ &- f_k^2 f_{\tau_1k}^1 - f_k^3 f_{\tau_1k}^4 + f_k^4 f_{\tau_1k}^3 + f_k^4 f_{\tau_2k}^2 + f_k^4 f_{\tau_2k}^3 - f_k^3 f_{\tau_2k}^4, \\ F_k^{13} &= (-\Delta_1 f_k^4 - \Delta_3 f_k^2 + f_k^3 f_{\tau_1k}^2 - f_k^4 f_{\tau_1k}^4 - f_k^4 f_{\tau_3k}^4 - f_k^2 f_{\tau_3k}^3) \mathbf{i} \\ &+ (\Delta_1 f_k^1 - \Delta_3 f_k^3 - f_k^2 f_{\tau_1k}^2 - f_k^4 f_{\tau_1k}^4 - f_k^4 f_{\tau_3k}^4 - f_k^2 f_{\tau_3k}^2) \mathbf{j} \\ &+ (\Delta_1 f_k^2 - \Delta_3 f_k^4 + f_k^2 f_{\tau_1k}^1 - f_k^4 f_{\tau_1k}^2 - f_k^4 f_{\tau_3k}^4 + f_k^4 f_{\tau_3k}^3 + f_k^2 f_{\tau_3k}^2) \mathbf{k} \\ &+ f_k^2 f_{\tau_1k}^4 - f_k^3 f_{\tau_1k}^1 - f_k^4 f_{\tau_1k}^2 - f_k^4 f_{\tau_3k}^2 + f_k^4 f_{\tau_3k}^4 + f_k^4 f_{\tau_3k}^3, \\ F_k^{14} &= (\Delta_1 f_k^3 - \Delta_4 f_k^2 + f_k^3 f_{\tau_1k}^1 + f_k^4 f_{\tau_1k}^2 + f_k^2 f_{\tau_4k}^4 + f_k^1 f_{\tau_4k}^3) \mathbf{i} \\ &+ (-\Delta_1 f_k^2 - \Delta_4 f_k^3 - f_k^2 f_{\tau_1k}^2 - f_k^3 f_{\tau_1k}^3 - f_k^3 f_{\tau_4k}^3 - f_k^2 f_{\tau_4k}^2) \mathbf{j} \\ &+ (\Delta_1 f_k^1 - \Delta_4 f_k^4 - f_k^2 f_{\tau_1k}^2 - f_k^3 f_{\tau_1k}^3 - f_k^3 f_{\tau_4k}^3 - f_k^2 f_{\tau_4k}^2) \mathbf{k} \\ &- f_k^2 f_{\tau_1k}^3 + f_k^3 f_{\tau_1k}^2 - f_k^4 f_{\tau_1k}^2 - f_k^3 f_{\tau_4k}^3 - f_k^2 f_{\tau_4k}^2) \mathbf{k} \end{split}$$

$$\begin{split} F_k^{23} &= (-\Delta_2 f_k^4 - \Delta_3 f_k^1 + f_k^4 f_{\tau_2 k}^2 + f_k^3 f_{\tau_2 k}^1 + f_k^1 f_{\tau_3 k}^3 + f_k^2 f_{\tau_3 k}^4) \mathbf{i} \\ &+ (\Delta_2 f_k^1 - \Delta_3 f_k^4 - f_k^1 f_{\tau_2 k}^2 + f_k^3 f_{\tau_2 k}^4 + f_k^4 f_{\tau_3 k}^3 - f_k^2 f_{\tau_3 k}^1) \mathbf{j} \\ &+ (\Delta_2 f_k^2 + \Delta_3 f_k^3 + f_k^1 f_{\tau_2 k}^1 + f_k^4 f_{\tau_2 k}^4 + f_k^4 f_{\tau_3 k}^4 + f_k^1 f_{\tau_3 k}^1) \mathbf{k} \\ &+ f_k^1 f_{\tau_2 k}^4 - f_k^4 f_{\tau_2 k}^1 + f_k^3 f_{\tau_2 k}^2 - f_k^4 f_{\tau_3 k}^1 + f_k^1 f_{\tau_3 k}^4 - f_k^2 f_{\tau_3 k}^3, \\ F_k^{24} &= (\Delta_2 f_k^3 - \Delta_4 f_k^1 + f_k^4 f_{\tau_2 k}^1 - f_k^3 f_{\tau_2 k}^2 - f_k^2 f_{\tau_4 k}^3 + f_k^1 f_{\tau_4 k}^4) \mathbf{i} \\ &+ (-\Delta_2 f_k^2 - \Delta_4 f_k^4 - f_k^1 f_{\tau_2 k}^1 - f_k^3 f_{\tau_2 k}^3 - f_k^3 f_{\tau_4 k}^3 - f_k^1 f_{\tau_4 k}^1) \mathbf{j} \\ &+ (\Delta_2 f_k^1 + \Delta_4 f_k^3 - f_k^1 f_{\tau_2 k}^2 - f_k^4 f_{\tau_2 k}^3 - f_k^3 f_{\tau_4 k}^3 - f_k^1 f_{\tau_4 k}^3) \mathbf{k} \\ &- f_k^1 f_{\tau_2 k}^3 + f_k^4 f_{\tau_2 k}^2 + f_k^3 f_{\tau_2 k}^1 + f_k^3 f_{\tau_4 k}^1 - f_k^2 f_{\tau_4 k}^3 - f_k^1 f_{\tau_4 k}^3, \\ F_k^{34} &= (\Delta_3 f_k^3 + \Delta_4 f_k^4 + f_k^1 f_{\tau_3 k}^1 + f_k^2 f_{\tau_3 k}^2 + f_k^2 f_{\tau_4 k}^2 + f_k^1 f_{\tau_4 k}^3) \mathbf{i} \\ &+ (-\Delta_3 f_k^2 - \Delta_4 f_k^1 + f_k^4 f_{\tau_3 k}^1 + f_k^2 f_{\tau_3 k}^3 + f_k^3 f_{\tau_4 k}^2 + f_k^1 f_{\tau_4 k}^4) \mathbf{j} \\ &+ (\Delta_3 f_k^1 - \Delta_4 f_k^2 + f_k^4 f_{\tau_3 k}^2 - f_k^1 f_{\tau_3 k}^3 - f_k^3 f_{\tau_4 k}^3 + f_k^2 f_{\tau_4 k}^4 + f_k^1 f_{\tau_4 k}^4) \mathbf{k} \\ &+ f_k^4 f_{\tau_3 k}^3 + f_k^1 f_{\tau_3 k}^2 - f_k^2 f_{\tau_3 k}^3 - f_k^3 f_{\tau_4 k}^4 + f_k^1 f_{\tau_4 k}^4) \mathbf{k} \\ &+ f_k^4 f_{\tau_3 k}^3 + f_k^4 f_{\tau_3 k}^2 - f_k^2 f_{\tau_3 k}^3 - f_k^3 f_{\tau_4 k}^4 + f_k^4 f_{\tau_4 k}^4 + f_k^4 f_{\tau_3 k}^4 + f_k^2 f_{\tau_3 k}^4 + f_k^4 f_{\tau_4 k}^4 + f_k^4 f_{\tau_4 k}^4 + f_k^4 f_{\tau_3 k}^4 - f_k^2 f_{\tau_4 k}^4 + f_k^4 f_{\tau_4 k}^4 + f_k^4$$

To obtain (3.13) we must take the imaginary part of these equations.

Theorem 3.1. The discrete curvature F in (3.11) is su(2)-valued if and only if

$$\begin{split} -f_k^2 f_{\tau_1 k}^1 - f_k^3 f_{\tau_1 k}^4 + f_k^4 f_{\tau_1 k}^3 + f_k^1 f_{\tau_2 k}^2 + f_k^4 f_{\tau_2 k}^3 - f_k^3 f_{\tau_2 k}^4 = 0, \\ f_k^2 f_{\tau_1 k}^4 - f_k^3 f_{\tau_1 k}^1 - f_k^4 f_{\tau_1 k}^2 - f_k^4 f_{\tau_3 k}^2 + f_k^1 f_{\tau_3 k}^3 + f_k^2 f_{\tau_3 k}^4 = 0, \\ -f_k^2 f_{\tau_1 k}^3 + f_k^3 f_{\tau_1 k}^2 - f_k^4 f_{\tau_1 k}^1 + f_k^3 f_{\tau_4 k}^2 - f_k^2 f_{\tau_4 k}^3 + f_k^1 f_{\tau_4 k}^4 = 0, \\ f_k^1 f_{\tau_2 k}^4 - f_k^4 f_{\tau_2 k}^1 + f_k^3 f_{\tau_2 k}^2 - f_k^4 f_{\tau_3 k}^1 + f_k^1 f_{\tau_3 k}^4 - f_k^2 f_{\tau_3 k}^3 = 0, \\ -f_k^1 f_{\tau_2 k}^3 + f_k^4 f_{\tau_2 k}^2 + f_k^3 f_{\tau_2 k}^2 + f_k^3 f_{\tau_4 k}^1 - f_k^2 f_{\tau_4 k}^4 - f_k^1 f_{\tau_4 k}^3 = 0, \\ f_k^1 f_{\tau_3 k}^3 + f_k^4 f_{\tau_2 k}^2 + f_k^2 f_{\tau_3 k}^1 - f_k^3 f_{\tau_4 k}^4 - f_k^2 f_{\tau_4 k}^4 - f_k^1 f_{\tau_4 k}^3 = 0, \\ f_k^4 f_{\tau_3 k}^3 + f_k^4 f_{\tau_3 k}^2 - f_k^2 f_{\tau_3 k}^1 - f_k^3 f_{\tau_4 k}^4 - f_k^2 f_{\tau_4 k}^4 + f_k^4 f_{\tau_4 k}^2 = 0. \end{split}$$

Proof. From the above, the assertion follows immediately.

Theorem 3.2. Let e be given by (3.6) and let \bar{e} be the conjugate quaternion of e. Then the 2-form $e \cup \bar{e}$ is self-dual, i.e.,

(3.14)
$$e \cup \bar{e} = *\tilde{\iota}(e \cup \bar{e}),$$

and $\bar{e} \cup e$ is anti-self-dual, i.e.,

$$(3.15) \qquad \qquad \bar{e} \cup e = - * \tilde{\iota}(\bar{e} \cup e).$$

Proof. Denote

$$e_i = \sum_k e_i^k, \qquad \varepsilon_{ij} = \sum_k \varepsilon_{ij}^k.$$

This implies $e_i \cup e_j = \varepsilon_{ij}$ and $e_j \cup e_i = -\varepsilon_{ij}$ for all i < j. Then we have

$$e \cup \overline{e} = (e_1 + e_2 \mathbf{i} + e_3 \mathbf{j} + e_4 \mathbf{k}) \cup (e_1 - e_2 \mathbf{i} - e_3 \mathbf{j} - e_4 \mathbf{k})$$

= $-2\{(e_1 \cup e_2 + e_3 \cup e_4)\mathbf{i} + (e_1 \cup e_3 - e_2 \cup e_4)\mathbf{j} + (e_1 \cup e_4 + e_2 \cup e_3)\mathbf{k}\}$
= $-2\{(\varepsilon_{12} + \varepsilon_{34})\mathbf{i} + (\varepsilon_{13} - \varepsilon_{24})\mathbf{j} + (\varepsilon_{14} + \varepsilon_{23})\mathbf{k}\}.$

By the definition of * and using (3.3), we get

$$*\tilde{\iota}(e\cup\bar{e}) = -2\tilde{\iota}\{(\tilde{\varepsilon}_{34}+\tilde{\varepsilon}_{12})\mathbf{i}+(-\tilde{\varepsilon}_{24}+\tilde{\varepsilon}_{13})\mathbf{j}+(\tilde{\varepsilon}_{23}+\tilde{\varepsilon}_{14})\mathbf{k}\}=e\cup\bar{e}.$$

In the same way we obtain (3.15).

Corollary 3.3. For any quaternionic 0-form f, the form $f \cup e \cup \overline{e}$ is self-dual and $f \cup \overline{e} \cup e$ is anti-self-dual.

Discrete self-dual and anti-self-dual equations (discrete analogs of equations (1.3)) are defined by

(3.16)
$$F = \tilde{\iota} * F, \qquad F = -\tilde{\iota} * F.$$

Using (3.10), by the definitions of $\tilde{\iota}$ and *, the first equation (self-dual) of (3.16) can be rewritten as

$$(3.17) F_k^{12} = F_k^{34}, F_k^{13} = -F_k^{24}, F_k^{14} = F_k^{23}.$$

By analogy with the continual case the solutions of (3.16) are called instantons and anti-instantons respectively.

4. DISCRETE INSTANTON AND ANTI-INSTANTON

Again in analogy with the continual case consider (3.8), where the components of f are given by

(4.1)
$$f_k = \frac{\overline{k}}{1+|k|^2}.$$

Here $k = k_1 + k_2 \mathbf{i} + k_3 \mathbf{j} + k_4 \mathbf{k}$, $k_i \in \mathbb{Z}$, and the norm |k| is defined by (2.3). Putting this in (3.9) we obtain

(4.2)
$$A_{k}^{1} = \frac{-k_{2}\mathbf{i} - k_{3}\mathbf{j} - k_{4}\mathbf{k}}{1 + |k|^{2}}, \qquad A_{k}^{2} = \frac{k_{1}\mathbf{i} - k_{4}\mathbf{j} + k_{3}\mathbf{k}}{1 + |k|^{2}},$$
$$A_{k}^{3} = \frac{k_{4}\mathbf{i} + k_{1}\mathbf{j} - k_{2}\mathbf{k}}{1 + |k|^{2}}, \qquad A_{k}^{4} = \frac{-k_{3}\mathbf{i} + k_{2}\mathbf{j} + k_{1}\mathbf{k}}{1 + |k|^{2}}.$$

It is convenient to denote

(4.3)
$$M_k^i = \frac{1}{(1+|k|^2)(1+|\tau_i k|^2)}, \quad i = 1, 2, 3, 4.$$

Substituting (4.2) in (3.12) and using (4.3) we find the components F_k^{ij} , for example,

$$\begin{split} F_k^{12} &= \{M_k^1(1+k_2^2-k_1^2-k_1)+M_k^2(1+k_1^2-k_2^2-k_2)\}\mathbf{i} \\ &+ \{M_k^1(k_4k_1+k_2k_3)-M_k^2(k_3k_2+k_4k_1)\}\mathbf{j} \\ &+ \{M_k^1(k_2k_4-k_1k_3)+M_k^2(k_1k_3-k_2k_4)\}\mathbf{k} \\ &+ M_k^1(k_1k_2+k_2)-M_k^2(k_1k_2+k_1). \end{split}$$

Note that the last term in F_k^{ij} has the form $M_k^i(k_ik_j + k_j) - M_k^j(k_ik_j + k_i)$. Hence, by Theorem 3.1, the curvature F defined by (4.2) is su(2)-valued if and only if

(4.4)
$$M_k^i(k_ik_j + k_j) - M_k^j(k_ik_j + k_i) = 0$$

for any $k_i \in \mathbb{Z}$, i, j = 1, 2, 3, 4 and i < j. An easy computation shows that equation (4.4) has only the solutions

(4.5)
$$\mu = k_1 = k_2 = k_3 = k_4, \quad k_i \in \mathbb{Z}.$$

Thus, the su(2)-valued discrete curvature 2-form F can be written in quaternionic form as

(4.6)
$$F = \sum_{k,k_i=\mu} M_{\mu} (2-2\mu) \{ (\varepsilon_{12}^k - \varepsilon_{34}^k) \mathbf{i} + (\varepsilon_{13}^k + \varepsilon_{24}^k) \mathbf{j} + (\varepsilon_{14}^k - \varepsilon_{23}^k) \mathbf{k} \},$$

where $M_{\mu} = M_k^1 = M_k^2 = M_k^3 = M_k^4$. From (4.3) we have $M_{\mu} = \frac{1}{2(1+4\mu^2)(1+\mu+2\mu^2)}$. Since $k_i = \mu$, in (4.6) we can write ε_{ij}^{μ} instead of ε_{ij}^k . If we consider the 0-form

(4.7)
$$\omega = \sum_{\mu} M_{\mu} (1-\mu) x^{\mu}, \qquad \mu \in \mathbb{Z},$$

and use the relation (see the proof of Theorem 3.2)

$$\bar{e} \cup e = 2\{(\varepsilon_{12} - \varepsilon_{34})\mathbf{i} + (\varepsilon_{13} + \varepsilon_{24})\mathbf{j} + (\varepsilon_{14} - \varepsilon_{23})\mathbf{k}\},\$$

then F can be written as

$$F = \omega \cup \bar{e} \cup e.$$

In view of Corollary 3.3, F is anti-self-dual, i.e., $F = -\tilde{\iota} * F$. Thus under the condition (4.5), A with components (4.1) describes an anti-instanton.

In the same manner we can see that the quaternionic 1-form

$$A = \operatorname{Im}(f \cup \bar{e}),$$

where f has the components

$$f_k = \frac{k}{1+|k|^2},$$

,

leads to an instanton solution of (3.17). Indeed, in this case the discrete curvature (3.13) has the form $F = \omega \cup e \cup \overline{e}$. Consequently, F is self-dual.

References

- M. F. Atiyah: Geometry of Yang-Mills Fields. Lezione Fermiane, Scuola Normale Superiore, Pisa, 1979.
- [2] A. A. Dezin: Multidimensional Analysis and Discrete Models. CRC Press, Boca Raton, 1995.
- [3] A. A. Dezin: Models generated by the Yang-Mills equations. Differ. Uravn. 29 (1993), 846–851; English translation in Differ. Equ. 29 (1993), 724–728.
- [4] D. Freed, K. Uhlenbeck: Instantons and Four-Manifolds. Springer, New York, 1984.
- [5] C. Nash, S. Sen: Topology and Geometry for Physicists. Acad. Press, London, 1989.
- [6] V. Sushch: Gauge-invariant discrete models of Yang-Mills equations. Mat. Zametki. 61 (1997), 742–754; English translation in Math. Notes. 61 (1997), 621–631.
- [7] V. Sushch: Discrete model of Yang-Mills equations in Minkowski space. Cubo A Math. Journal. 6 (2004), 35–50.
- [8] V. Sushch: A gauge-invariant discrete analog of the Yang-Mills equations on a double complex. Cubo A Math. Journal. 8 (2006), 61–78.

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