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# ON CONDITIONING OF SCHUR COMPLEMENTS OF H-TFETI CLUSTERS FOR 2D PROBLEMS GOVERNED BY LAPLACIAN 

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#### Abstract

Bounds on the spectrum of the Schur complements of subdomain stiffness matrices with respect to the interior variables are key ingredients in the analysis of many domain decomposition methods. Here we are interested in the analysis of floating clusters, i.e. subdomains without prescribed Dirichlet conditions that are decomposed into still smaller subdomains glued on primal level in some nodes and/or by some averages. We give the estimates of the regular condition number of the Schur complements of the clusters arising in the discretization of problems governed by 2D Laplacian. The estimates depend on the decomposition and discretization parameters and gluing conditions. We also show how to plug the results into the analysis of H-TFETI methods and compare the estimates with numerical experiments. The results are useful for the analysis and implementation of powerful massively parallel scalable algorithms for the solution of variational inequalities.


Keywords: two-level domain decomposition; hybrid FETI; Schur complement; bounds on the spectrum

MSC 2010: 34B16, 34C25

## 1. Introduction

Variants of FETI (finite element tearing and interconnecting) methods introduced by Farhat and Roux [7], [8] belong to the most powerful methods for massively parallel solution of large discretized elliptic partial differential equations. The basic

[^0]idea is to decompose the domain into subdomains which are joined by Lagrange multipliers. After eliminating the primal variables, the original problem reduces to a smaller global problem and a number of local problems that can be solved in parallel. Moreover, if applied to any variational inequality, the duality transforms the inequality constraints into bound constraints. The conditioning of the global problem depends on the conditioning of Schur complements of local stiffness matrices with respect to interior variables. The bounds on the spectra of local Schur complements are well known (see e.g. Brenner [1] or Pechstein [13]) and are the essential ingredients of the analysis of any FETI-type method (see e.g. Farhat, Mandel, and Roux [6] or Tosseli and Widlund [14]).

Here we are interested in the estimates that are necessary for the analysis of H-TFETI (hybrid total FETI, see e.g. [4]) without preconditioner. This method exploits the decomposition of domains into floating clusters, i.e. the subdomains without Dirichlet conditions that are decomposed into still smaller subdomains glued on the primal level in some nodes and/or averages. The two-level structure of the stiffness matrices arising from applications of H-TFETI complies well with the hierarchical organization of modern supercomputers - the clusters and their subdomains can be naturally associated with the nodes and their cores, respectively. We give the estimates in terms of the discretization and decomposition parameters of the regular condition number of the Schur complements of clusters arising from the discretization of problems governed by 2D Laplacian. We consider various gluing conditions, plug the results into the analysis of H-TFETI methods, and compare the results with numerical experiments.

Let us mention that the interest in the conditioning of FETI without preconditioners is motivated by the possibility to work with bound or separable inequality constraints. The preconditioning transforms these constraints into more general constraints (see e.g. Dostál et al. [4]). Let us mention that H-TFETI has been successfully applied to the solution of nonlinear elliptic problems discretized by some hundred billions of nodal variables [4].

## 2. Domain decomposition, subdomains and clusters

Let us consider a problem governed by the Laplacian on the 2 D unit square $\Omega$ with the boundary $\Gamma$, such as the Poisson equation

$$
\begin{equation*}
\Delta u=f \tag{2.1}
\end{equation*}
$$

with homogeneous Dirichlet, Neumann, or Signorini conditions. For the application of TFETI, let us decompose $\Omega$ into square subdomains $\Omega_{i}$ of equal side-length $H_{s}$ with the boundaries $\Gamma_{i}$ as in Fig. $1, i=1, \ldots, n_{s}, n_{s}=1 / H_{s}^{2}$.


Figure 1. Decomposition of $\Omega$ into subdomains $\Omega_{i}$.
In each $\Omega_{i}$, we introduce a regular triangularization with the discretization parameter $h$ and the pattern depicted in Fig. 2. We assume matching discretization of subdomains, so the nodes should coincide on the interface. The total number of nodes is denoted by $n$.


Figure 2. Subdomain $\Omega_{i}$ and its triangularization.
For each subdomain $\Omega_{i}$, we introduce the standard finite element linear basis functions $\varphi_{j}^{i}(x)$ and set up local stiffness matrices $K^{i}$, local nodal displacement vectors $u^{i}$, and local load vectors $f^{i}$,

$$
K=\operatorname{diag}\left(K^{1}, \ldots, K^{n_{s}}\right), \quad u=\left[\begin{array}{c}
u^{1} \\
\vdots \\
u^{n_{s}}
\end{array}\right], \quad f=\left[\begin{array}{c}
f^{1} \\
\vdots \\
f^{n_{s}}
\end{array}\right], \quad i=1, \ldots, n_{s} .
$$

The basis functions span the space $V_{h}\left(\Omega_{i}\right)$ with the elements

$$
u_{h}^{i}(x)=\sum_{j} u_{j}^{i} \varphi_{j}^{i}(x), \quad x \in \Omega_{i} .
$$

Notice that each local stiffness matrix $K^{i}$ is symmetric positive semidefinite (SPS) with the kernel spanned by the vector

$$
e_{i}=[1, \ldots, 1]^{\mathrm{T}} .
$$

The solution of the discretized problem (2.1) can be obtained by the solution of a constrained quadratic programming problem

$$
\begin{equation*}
\min \left(\frac{1}{2} x^{\mathrm{T}} K x-f^{\mathrm{T}} x\right) \quad \text { subject to } \quad B_{E} x=o \quad \text { and } \quad B_{I} x \leqslant o, \tag{2.2}
\end{equation*}
$$

where $B_{E}$ represents gluing of the subdomains and the Dirichlet boundary conditions, $B_{I}$ enforces the Signorini conditions, and $o$ denotes zero vector. For example, the continuity of the solution in the corners of interior subdomains is enforced by three rows of $B$ with nonzero entries placed in four columns corresponding to the global indices of the corner variables,

$$
\left[\begin{array}{ccccccccc}
\ldots & 1 & \ldots & -1 & \ldots & 0 & \ldots & 0 & \ldots \\
\ldots & 0 & \ldots & 0 & \ldots & 1 & \ldots & -1 & \ldots \\
\ldots & 1 & \ldots & 1 & \ldots & -1 & \ldots & -1 & \ldots
\end{array}\right]
$$

Matrices $B_{E}$ and $B_{I}$ can be considered as submatrices of matrix $B$ with the column blocks complying with the block structure of $K$, i.e.

$$
B=\left[\begin{array}{c}
B_{E} \\
B_{I}
\end{array}\right]=\left[B_{1}, \ldots, B_{n_{s}}\right] .
$$

Both, original FETI algorithm (also referred to as FETI1) [7] and TFETI [3], were proposed for linear problems. The idea was to switch to the constrained dual problem in Lagrange multipliers and solve it iteratively with preconditioning by the projector

$$
P=I-G\left(G^{\mathrm{T}} G\right)^{+} G^{\mathrm{T}}, \quad G=B R,
$$

where $\left(G^{\mathrm{T}} G\right)^{+}$denotes a left generalized inverse. Recal that if $A$ is any square matrix, then

$$
A A^{+} A=A
$$

The method was later adapted to the solution of variational inequalities [2], [4]. The dual problem reads

$$
\begin{equation*}
\min \left(\frac{1}{2} \lambda^{\mathrm{T}} F \lambda-d^{\mathrm{T}} \lambda\right) \quad \text { subject to } G \lambda=o \text { and } \lambda_{I} \geqslant o, \tag{2.3}
\end{equation*}
$$

where

$$
F=B K^{+} B^{\mathrm{T}}, \quad G=R^{\mathrm{T}} B^{\mathrm{T}}, \quad \lambda=\left[\lambda_{E}^{\mathrm{T}}, \lambda_{I}^{\mathrm{T}}\right]^{\mathrm{T}},
$$

and denotes the Lagrange multipliers with the components enforcing the equality and inequality constraints, respectively, and $d$ is a vector the specification of which
is not relevant in this paper. The regular condition number $\bar{\kappa}(F)=\kappa(F \mid \operatorname{Im} F)$ was shown in [6] to satisfy

$$
\begin{equation*}
\bar{\kappa}(F) \leqslant C H / h, \tag{2.4}
\end{equation*}
$$

with a constant $C$ independent of the decomposition and discretization parameters $H$ and $h$, respectively. This estimate guarantees an optimal complexity of the method provided the cost of the action of $P$ does not dominate the cost of the iteration.

To overcome the latter limitation and to reduce the dimension of $G^{\mathrm{T}} G$ without compromising the number of subdomains, Farhat, Lesoinne, and Pierson [5] and Klawonn and Rheinbach [9], [10] proposed to enforce some gluing constraints explicitly. For example, to glue four adjacent subdomains in the only common node

$$
x \in \bar{\Omega}_{i} \cap \bar{\Omega}_{j} \cap \bar{\Omega}_{k} \cap \bar{\Omega}_{l},
$$

it is enough to replace the four columns of identity matrix which correspond to $x$ by a vector with four appropriately placed ones to get the matrix $L$ which transforms global variables $\widetilde{u}$ to $u$,

$$
u=L \widetilde{u}
$$

It can be checked that the stiffness matrix $\widetilde{K}$ of $\Omega$ and the constraint matrices $\widetilde{B}_{E}$ and $\widetilde{B}_{I}$ in remaining variables are given by

$$
\widetilde{K}=L^{\mathrm{T}} K L, \quad \widetilde{B}_{E}=B_{E} L, \quad \widetilde{B}_{I}=B_{I} L
$$

Moreover, the kernel of the stiffness matrix $\widetilde{K}^{m}$ of the cluster obtained by gluing the four subdomains in one node is only one dimensional, while the dimension of the kernel of $\operatorname{diag}\left(K^{i}, K^{j}, K^{k}, K^{l}\right)$ is four. The cluster resulting from the gluing of four adjacent subdomains in corners is depicted in Fig. 3.


Figure 3. The gluing of four subdomains in adjacent corners.
More generally, we can split the equality constraints into two blocks $B_{E P}$ and $B_{E D}$,

$$
B_{E}=\left[\begin{array}{l}
B_{E P} \\
B_{E D}
\end{array}\right]
$$

and use $B_{E P}$ to eliminate some primal variables. The constraint matrix $B_{E P}$ can define not only gluing of some nodes but also more general relations between the variables on neighbouring edges such as zero average of the sum of variables in the interior of adjacent edges. The details of implementation are a bit tricky, but well known, see e.g. Klawonn and Rheinbach [10], Lee [11], or Dostál et al. [4], Chap. 19. In this way, we can implement a decomposition of $\Omega$ into clusters $\widetilde{\Omega}^{i}, i=1, \ldots, n_{c}$. As a result, we shall obtain the H-TFETI problem to find

$$
\begin{equation*}
\min \left(\frac{1}{2} \lambda^{\mathrm{T}} \widetilde{F} \lambda-\tilde{d}^{\mathrm{T}} \lambda\right) \quad \text { subject to } \widetilde{G} \lambda=o \text { and } \lambda_{I} \geqslant o \tag{2.5}
\end{equation*}
$$

where

$$
\widetilde{F}=\widetilde{B} \widetilde{K}^{+} \widetilde{B}^{\mathrm{T}}, \quad \widetilde{K}=\operatorname{diag}\left(\widetilde{K}^{1}, \ldots, \widetilde{K}^{n_{c}}\right), \quad \widetilde{G}=\widetilde{R}^{\mathrm{T}} \widetilde{B}^{\mathrm{T}}, \quad \lambda=\left[\lambda_{E D}^{\mathrm{T}}, \lambda_{I}^{\mathrm{T}}\right]^{\mathrm{T}}
$$

and now denotes the Lagrange multipliers associated with the components enforcing the inequality and remaining equality constraints of the modified problem described by a matrix $\widetilde{B}$, respectively. The preconditioning is carried out by a modified projector

$$
\widetilde{P}=I-\widetilde{G}\left(\widetilde{G}^{\mathrm{T}} \widetilde{G}\right)^{+} \widetilde{G}
$$

In the next sections, we shall show that it is possible, at least for small clusters, to get similar optimality results for the solution of (2.5) as those reported above for TFETI based algorithms for (2.3). Let us point out that there are better results for TFETI and H-TFETI with standard preconditioners (see e.g. [13] or [14]), but these results do not support optimality of algorithms for variational inequalities. The reason is that there is no algorithm for the solution of quadratic programming problems with general inequality constraints that enjoy the rate of convergence in terms of conditioning of the Hessian.

## 3. Schur complements of clusters and general estimates

The condition number of $\widetilde{F}$ can be estimated in two steps. The first one is a subject of the following simple lemma.

Lemma 3.1. Let there be constants $0<C_{1}<C_{2}$ such that for each $\lambda \in \mathbb{R}^{m}$,

$$
\begin{equation*}
C_{1}\|\lambda\|^{2} \leqslant\left\|\widetilde{B}^{\mathrm{T}} \lambda\right\|^{2} \leqslant C_{2}\|\lambda\|^{2}, \tag{3.1}
\end{equation*}
$$

where $\|\cdot\|$ denotes the Euclidean norm. Then

$$
\begin{equation*}
\bar{\kappa}(\widetilde{P} \widetilde{F} \widetilde{P}) \leqslant \frac{C_{2} \max \left\{\left\|\widetilde{K}^{i}\right\|: i=1, \ldots, n_{c}\right\}}{C_{1} \min \left\{\bar{\lambda}_{\text {min }}\left(\widetilde{K}^{i}\right): i=1, \ldots, n_{c}\right\}} \tag{3.2}
\end{equation*}
$$

Proof. The proof of this lemma is rather trivial; it uses only the observations that if $\lambda \in \operatorname{Im} \widetilde{P}$, then $\widetilde{G}^{\mathrm{T}} \lambda=o$, i.e. $\widetilde{B}^{\mathrm{T}} \lambda \in \operatorname{Im} K$, that the nonzero eigenvalues of $\widetilde{K}$ are reciprocal to the corresponding eigenvalues of $\widetilde{K}^{+}$, and that the spectrum $\sigma(\widetilde{K})$ of $\widetilde{K}$ satisfies

$$
\sigma(\widetilde{K})=\bigcup_{i=1}^{n_{c}} \sigma\left(\widetilde{K}^{i}\right)
$$

However, the estimate given by Lemma 3.1 is a bit pessimistic. The reason is that $\operatorname{Im} \widetilde{B}$ is spanned by the vectors that have zero components corresponding to the variables in the interior of $\Omega_{i}$. To enhance this observation, let us define the (extended) skeleton $\Sigma$ of the decomposition by

$$
\Sigma:=\bigcup_{i=1}^{n_{s}} \Gamma_{i}
$$

and decompose the set of indices $\mathcal{N}=\{1, \ldots, n\}$ into the indices of skeleton nodes $\mathcal{S}$ and subdomain interior nodes $\mathcal{I}$. For any matrix $A \in \mathbb{R}^{m \times n}$ and the subsets $\mathcal{I} \subseteq\{i=1, \ldots, m\}$ and $\mathcal{J} \subseteq\{j=1, \ldots, n\}$, let $A_{\mathcal{I} \mathcal{J}}$ denote a submatrix of $A$ with the rows $i \in \mathcal{I}$ and $j \in \mathcal{J}$. Then it is easy to check that

$$
\left(\widetilde{K}^{+}\right)_{\mathcal{S S}}=\widetilde{S}^{+}, \quad \widetilde{S}=\widetilde{K}_{\mathcal{S S}}-\widetilde{K}_{\mathcal{S} \mathcal{I}} \widetilde{K}_{\mathcal{I I}}^{-1} \widetilde{K}_{\mathcal{I S}} .
$$

Matrix $\widetilde{S}$ is called the Schur complement of the block of subdomain interior variables. The same formula holds for the clusters, i.e.

$$
\left(\left(\widetilde{K}^{i}\right)^{+}\right)_{\mathcal{S}_{i} \mathcal{S}_{i}}=\left(\widetilde{S}^{i}\right)^{+}, \quad \widetilde{S}^{i}=\widetilde{K}_{\mathcal{S}_{i} \mathcal{S}_{i}}^{i}-\widetilde{K}_{\mathcal{S}_{i} \mathcal{I}_{i}}^{i}\left(\widetilde{K}_{\mathcal{I}_{i} \mathcal{I}_{i}}^{i}\right)^{-1} \widetilde{K}_{\mathcal{I}_{i} \mathcal{S}_{i}}^{i}, \quad i=1, \ldots, n_{c},
$$

where $\mathcal{S}_{i}$ and $\mathcal{J}_{i}$ denote the boundary and interior indices of the subdomains of clusters, respectively. We can enhance these observations into the following corollary.

Corollary 3.1. Let the assumptions of Lemma 3.1 be satisfied. Then

$$
\begin{equation*}
\bar{\kappa}(\widetilde{P} \widetilde{F} \widetilde{P}) \leqslant \frac{C_{2} \max \left\{\left\|\widetilde{S}^{i}\right\|: i=1, \ldots, n_{c}\right\}}{C_{1} \min \left\{\bar{\lambda}_{\min }\left(\widetilde{S}^{i}\right): i=1, \ldots, n_{c}\right\}} \tag{3.3}
\end{equation*}
$$

It is useful to observe that the local Schur complements $\widetilde{S}^{i}$ are closely related to the harmonic extensions of functions from the boundary of subdomains $\Omega_{i}$. In particular, if

$$
u_{h}^{i}(x)=\sum_{j \in \mathcal{S}_{i}} u_{j}^{i} \varphi_{j}(x)
$$

is a discrete harmonic function with the indices of the interior and boundary displacements in $\mathcal{I}_{i}$ and $\mathcal{B}_{i}$, respectively, so that

$$
K_{\mathcal{I}_{i} \mathcal{I}_{i}}^{i}\left(u_{h}^{i}\right)_{\mathcal{I}_{i}}+K_{\mathcal{I}_{i} \mathcal{B}_{i}}^{i}\left(u_{h}^{i}\right)_{\mathcal{B}_{i}}=o,
$$

then

$$
\begin{equation*}
\left(u_{\mathcal{B}}^{i}\right)^{\mathrm{T}} S^{i} u_{\mathcal{B}}^{i}=\left(u^{i}\right)^{\mathrm{T}} K^{i} u^{i}=\int_{\Omega^{i}}\left\|\nabla u_{h}^{i}\right\|^{2} \mathrm{~d} \Omega . \tag{3.4}
\end{equation*}
$$

A similar relation holds for the clusters and their skeletons. It follows that we can bound the spectrum of $\widetilde{F}$ by the analysis of $V_{h}(\Omega)$.

## 4. Gluing by corners

Let us consider a square subdomain $\Omega_{i}$ of side length $H_{s}$ with the discretization introduced in Section 2 (see Fig. 2).

Theorem 4.1. Let $u_{h} \in V_{h}\left(\Omega_{i}\right)$ denote a discrete harmonic function, i.e.

$$
\forall v_{h} \in V_{h}\left(\Omega_{i}\right) \cap H_{0}^{1}\left(\Omega_{i}\right): \int_{\Omega_{i}} \nabla u_{h} \cdot \nabla v_{h} \mathrm{~d} \Omega=0
$$

where the dot denotes the Euclidean scalar product in $\mathbb{R}^{2}$. Then

$$
\begin{equation*}
\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega \leqslant 3 \sum_{x_{k} \in \Gamma_{i}}\left(u_{h}\left(x_{k}\right)\right)^{2} \tag{4.1}
\end{equation*}
$$

(we sum the squares of values of $u_{h}$ in all vertices of the triangles on the boundary of $\Omega_{i}$ ).

Proof. Let $\tilde{u}_{h} \in V_{h}\left(\Omega_{i}\right)$ be defined in the vertices of the triangles by

$$
\tilde{u}_{h}\left(x_{k}\right):= \begin{cases}u_{h}\left(x_{k}\right), & x_{k} \in \Gamma_{i}, \\ 0, & x_{k} \in \bar{\Omega}_{i} \backslash \Gamma_{i} .\end{cases}
$$

Let $T$ be a triangle of a triangulation of $\Omega_{i}$ with the vertices denoted by $x_{i}, x_{j}, x_{k}$, so that the right angle is at the vertex $x_{k}$. Then

$$
\begin{equation*}
\int_{T}\left\|\nabla \tilde{u}_{h}\right\|^{2} \mathrm{~d} \Omega=\frac{1}{2}\left[\left(\tilde{u}_{h}\left(x_{i}\right)-\tilde{u}_{h}\left(x_{k}\right)\right)^{2}+\left(\tilde{u}_{h}\left(x_{j}\right)-\tilde{u}_{h}\left(x_{k}\right)\right)^{2}\right] . \tag{4.2}
\end{equation*}
$$

If

$$
\tilde{u}_{h}\left(x_{i}\right) \neq 0 \neq \tilde{u}_{h}\left(x_{k}\right) \quad \text { or } \quad \tilde{u}_{h}\left(x_{j}\right) \neq 0 \neq \tilde{u}_{h}\left(x_{k}\right),
$$

then substituting of

$$
\begin{array}{ll} 
& \left(\tilde{u}_{h}\left(x_{i}\right)-\tilde{u}_{h}\left(x_{k}\right)\right)^{2} \leqslant 2\left(\left(\tilde{u}_{h}\left(x_{i}\right)\right)^{2}+\left(\tilde{u}_{h}\left(x_{k}\right)\right)^{2}\right) \\
\text { or } \quad & \left(\tilde{u}_{h}\left(x_{j}\right)-\tilde{u}_{h}\left(x_{k}\right)\right)^{2} \leqslant 2\left(\left(\tilde{u}_{h}\left(x_{j}\right)\right)^{2}+\left(\tilde{u}_{h}\left(x_{k}\right)\right)^{2}\right)
\end{array}
$$

into (4.2) yields

$$
\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega \leqslant \int_{\Omega_{i}}\left\|\nabla \tilde{u}_{h}\right\|^{2} \mathrm{~d} \Omega \leqslant 3 \sum_{x_{k} \in \Gamma_{i}}\left(\tilde{u}_{h}\left(x_{k}\right)\right)^{2}=3 \sum_{x_{k} \in \Gamma_{i}}\left(u_{h}\left(x_{k}\right)\right)^{2} .
$$

Lemma 4.1. There exists a constant $C>0$ independent of $H_{s}$ and $h$ such that

$$
\left\|u_{h}-\bar{u}_{h}\right\|_{L^{\infty}\left(\Omega_{i}\right)}^{2} \leqslant C\left(1+\ln \frac{H_{s}}{h}\right) \int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega
$$

for every $u_{h} \in V_{h}\left(\Omega_{i}\right)$, where

$$
\bar{u}_{h}=\frac{1}{H_{s}^{2}} \int_{\Omega_{i}} u_{h} \mathrm{~d} \Omega
$$

Proof. See [12], Corollary 3.2.

Lemma 4.2. There exists a constant $C>0$ independent of $H_{s}$ and $h$ such that for every $u_{h} \in V_{h}\left(\Omega_{i}\right)$ and $\alpha \in\left[\min _{x \in \bar{\Omega}_{i}} u_{h}(x), \max _{x \in \bar{\Omega}_{i}} u_{h}(x)\right]$

$$
\left\|u_{h}-\alpha\right\|_{L^{\infty}\left(\Omega_{i}\right)}^{2} \leqslant C\left(1+\ln \frac{H_{s}}{h}\right) \int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega .
$$

Proof. The proof is an easy consequence of Lemma 4.1, because

$$
\begin{aligned}
\left\|u_{h}-\alpha\right\|_{L^{\infty}\left(\Omega_{i}\right)} & =\max \left\{\max _{x \in \bar{\Omega}_{i}} u_{h}(x)-\alpha, \alpha-\min _{x \in \bar{\Omega}_{i}} u_{h}(x)\right\} \\
& \leqslant \max _{x \in \bar{\Omega}_{i}} u_{h}(x)-\min _{x \in \bar{\Omega}_{i}} u_{h}(x) \\
& \leqslant 2 \max \left\{\max _{x \in \bar{\Omega}_{i}} u_{h}(x)-\bar{u}_{h}, \bar{u}_{h}-\min _{x \in \bar{\Omega}_{i}} u_{h}(x)\right\} \\
& =2\left\|u_{h}-\bar{u}_{h}\right\|_{L^{\infty}\left(\Omega_{i}\right)} .
\end{aligned}
$$

To simplify the notation, let us now denote by $\Omega$ a cluster of four equal nonoverlapping square subdomains $\Omega_{1}, \ldots, \Omega_{4}$ of side length $H_{s}$, so that $\bigcap_{i=1}^{4} \bar{\Omega}_{i}=\left\{x_{0}\right\}$ (see Fig. 4), and let $\Sigma$ denote the skeleton of $\Omega$. Define

$$
\begin{aligned}
V_{h}(\Omega) & :=\left\{u_{h}=\left(u_{h}^{1}, u_{h}^{2}, u_{h}^{3}, u_{h}^{4}\right): u_{h}^{i} \in V_{h}\left(\Omega_{i}\right)\right\}, \\
\int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega & :=\sum_{i=1}^{4} \int_{\Omega_{i}}\left\|\nabla u_{h}^{i}\right\|^{2} \mathrm{~d} \Omega \\
\sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} & :=\sum_{i=1}^{4} \sum_{x_{k} \in \Gamma_{i}}\left(u_{h}^{i}\left(x_{k}\right)\right)^{2}, \\
\left\|u_{h}\right\|_{L^{\infty}(\Omega)} & :=\max \left\{\left\|u_{h}^{1}\right\|_{L^{\infty}\left(\Omega_{1}\right)}, \ldots,\left\|u_{h}^{4}\right\|_{L^{\infty}\left(\Omega_{4}\right)}\right\}, \\
\max _{x \in \bar{\Omega}} u_{h}(x) & :=\max \left\{\max _{x \in \bar{\Omega}_{1}} u_{h}^{1}(x), \ldots, \max _{x \in \bar{\Omega}_{4}} u_{h}^{4}(x)\right\}, \\
\min _{x \in \bar{\Omega}} u_{h}(x) & :=\min \left\{\min _{x \in \bar{\Omega}_{1}} u_{h}^{1}(x), \ldots, \min _{x \in \bar{\Omega}_{4}} u_{h}^{4}(x)\right\} .
\end{aligned}
$$

In this and the following section, we use simplified notations to improve readability. The following theorem is the main result of this section.


Figure 4. Cluster $\Omega$ joined in one point.
Theorem 4.2. The following statements are true for the cluster obtained by gluing the subdomains in a corner.
(i) There exists a constant $C>0$ independent of $H_{s}$ and $h$ such that

$$
\left(\max _{x \in \bar{\Omega}} u_{h}(x)-\min _{x \in \bar{\Omega}} u_{h}(x)\right)^{2} \leqslant C\left(1+\ln \frac{H_{s}}{h}\right) \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega
$$

for every $u_{h} \in V_{h}(\Omega)$ which satisfies

$$
u_{h}^{1}\left(x_{0}\right)=u_{h}^{2}\left(x_{0}\right)=u_{h}^{3}\left(x_{0}\right)=u_{h}^{4}\left(x_{0}\right) .
$$

(ii) If $u_{h}$ satisfies also

$$
\int_{\Sigma} u_{h} \mathrm{~d} \Gamma:=\sum_{i=1}^{4} \int_{\Gamma_{i}} u_{h}^{i} \mathrm{~d} \Gamma=0
$$

then there exists a constant $C>0$ independent of $H_{s}$ and $h$ such that

$$
\begin{equation*}
C \frac{h}{H_{s}\left(1+\ln \left(H_{s} / h\right)\right)} \sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} \leqslant \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega \tag{4.3}
\end{equation*}
$$

Proof of Theorem 4.2. (i) Let $i$ and $j$ be such that

$$
\max _{x \in \bar{\Omega}} u_{h}(x)=\max _{x \in \bar{\Omega}_{i}} u_{h}^{i}(x) \quad \text { and } \quad \min _{x \in \bar{\Omega}} u_{h}(x)=\min _{x \in \bar{\Omega}_{j}} u_{h}^{j}(x) .
$$

Using Lemma 4.2, we get

$$
\begin{aligned}
\left(\max _{x \in \bar{\Omega}} u_{h}(x)-\min _{x \in \bar{\Omega}} u_{h}(x)\right)^{2} & =\left(\left(\max _{x \in \bar{\Omega}_{i}} u_{h}^{i}(x)-\alpha\right)+\left(\alpha-\min _{x \in \bar{\Omega}_{j}} u_{h}^{j}(x)\right)\right)^{2} \\
& \leqslant 2\left(\left\|u_{h}^{i}-\alpha\right\|_{L^{\infty}\left(\Omega_{i}\right)}\right)^{2}+2\left(\left\|u_{h}^{j}-\alpha\right\|_{L^{\infty}\left(\Omega_{j}\right)}\right)^{2} \\
& \leqslant C\left(1+\ln \frac{H_{s}}{h}\right) \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega
\end{aligned}
$$

where

$$
\alpha:=u_{h}^{1}\left(x_{0}\right)=u_{h}^{2}\left(x_{0}\right)=u_{h}^{3}\left(x_{0}\right)=u_{h}^{4}\left(x_{0}\right) .
$$

(ii) If $\int_{\Sigma} u_{h} \mathrm{~d} \Gamma=0$, it follows that

$$
\min _{x \in \bar{\Omega}} u_{h}(x) \leqslant 0 \leqslant \max _{x \in \bar{\Omega}} u_{h}(x)
$$

and therefore,

$$
\left\|u_{h}\right\|_{L^{\infty}(\Omega)}^{2} \leqslant\left(\max _{x \in \bar{\Omega}} u_{h}(x)-\min _{x \in \bar{\Omega}} u_{h}(x)\right)^{2}
$$

Now it suffices to use (i) and

$$
\sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} \leqslant 16 \frac{H_{s}}{h}\left\|u_{h}\right\|_{L^{\infty}(\Omega)}^{2}
$$

Applying (4.1) and (4.3) to the harmonic function $u_{h}$ satisfying $\int_{\Sigma} u_{h} \mathrm{~d} \Gamma=0$, we get

$$
\begin{equation*}
\frac{C_{1} h}{H_{s}}\left(1+\ln \frac{H_{s}}{h}\right)^{-1} \sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} \leqslant \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \leqslant 3 \sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} \tag{4.4}
\end{equation*}
$$

Using the above inequalities, (3.4) and simple manipulations, we get the following corollary.

Corollary 4.1. Let the assumption of Lemma 3.1 be satisfied. Then there are constants $C_{1}$ and $C_{2}$ independent of $H_{s}$ and $h$ such that

$$
\begin{equation*}
C_{1} \leqslant \bar{\lambda}_{\min }(\widetilde{P} \widetilde{F} \widetilde{P}) \leqslant\|\widetilde{P} \widetilde{F} \widetilde{P}\| \leqslant C_{2} \frac{H_{s}}{h}\left(1+\ln \frac{H_{s}}{h}\right) \tag{4.5}
\end{equation*}
$$

Remark. If $B B^{\mathrm{T}}=I$, then (4.5) holds with $C_{1}=1 / 3$.
The following example shows that we can neither exclude the term $\left(1+\ln \left(/ H_{s} h\right)\right)$ nor replace it with a term of lower order in inequality (4.3).

Example 4.1. Let us consider the domains

$$
\begin{array}{ll}
\Omega_{1}=\left(0, H_{s}\right) \times\left(0, H_{s}\right), & \Omega_{2}=\left(-H_{s}, 0\right) \times\left(0, H_{s}\right), \\
\Omega_{3}=\left(-H_{s}, 0\right) \times\left(-H_{s}, 0\right), & \Omega_{4}=\left(0, H_{s}\right) \times\left(-H_{s}, 0\right)
\end{array}
$$

and $u_{h}=\left(u_{h}^{1}, \ldots, u_{h}^{4}\right) \in V_{h}(\Omega)$ such that

$$
\begin{aligned}
u_{h}^{1}(i h, j h) & :=\ln (i+j+1) \quad \text { for } i, j \in\left\{0, \ldots, \frac{H_{s}}{h}\right\} \\
u_{h}^{2}(x, y) & :=0 \quad \text { for }(x, y) \in \bar{\Omega}_{2} \\
u_{h}^{3}(x, y) & :=-u_{h}^{1}(-x,-y) \quad \text { for }(x, y) \in \bar{\Omega}_{3} \\
u_{h}^{4}(x, y) & :=0 \quad \text { for }(x, y) \in \bar{\Omega}_{4}
\end{aligned}
$$

Then $u_{h}$ satisfies the assumptions of Theorem 4.2 and for every (admissible) $h>0$ we have (for simplicity we denote $m=H_{s} / h$ )

$$
\begin{equation*}
\sum_{x_{j} \in \Gamma_{1}}\left(u_{h}^{1}\left(x_{j}\right)\right)^{2} \geqslant \ln ^{2}(m+1)+\ldots+\ln ^{2}(2 m+1) \geqslant m \ln ^{2} m \tag{4.6}
\end{equation*}
$$

Moreover, using Lagrange's Mean Value Theorem, we obtain for every $h>0$

$$
\begin{align*}
\int_{\Omega_{1}}\left\|\nabla u_{h}^{1}\right\|^{2} \mathrm{~d} \Omega & =\sum_{i, j=1}^{m}\left[(\ln (i+j+1)-\ln (i+j))^{2}+(\ln (i+j)-\ln (i+j-1))^{2}\right]  \tag{4.7}\\
& \leqslant \sum_{i, j=1}^{m}\left[\frac{1}{(i+j)^{2}}+\frac{1}{(i+j-1)^{2}}\right] \leqslant 2 \sum_{i, j=1}^{m} \frac{1}{(i+j-1)^{2}} \\
& \leqslant 2 \sum_{\substack{i, j \in \mathbb{N} \\
i+j \leqslant 2 m}} \frac{1}{(i+j-1)^{2}}=2 \sum_{k=2}^{2 m}(k-1) \cdot \frac{1}{(k-1)^{2}}=2 \sum_{k=2}^{2 m} \frac{1}{k-1} \\
& \leqslant 2(1+\ln (2 m-1)) .
\end{align*}
$$

Combining (4.6) and (4.7), we obtain

$$
\begin{aligned}
\frac{\sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2}}{H_{s} h^{-1} \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega}=\frac{2 \sum_{x_{j} \in \Gamma_{1}}\left(u_{h}^{1}\left(x_{j}\right)\right)^{2}}{2 H_{s} h^{-1} \int_{\Omega_{1}}\left\|\nabla u_{h}^{1}\right\|^{2} \mathrm{~d} \Omega} & \geqslant \frac{\ln ^{2} H_{s} / h}{2\left(1+\ln \left(2 H_{s} / h-1\right)\right)} \\
& \approx\left(1+\ln \frac{H_{s}}{h}\right) \quad \text { for } h \rightarrow 0+
\end{aligned}
$$

## 5. Gluing by averages

Let us consider a square subdomain $\Omega_{i}$ of side length $H_{s}$ and let one of its sides be denoted by $\Phi$. We denote by $x_{0}, x_{1}, \ldots, x_{n}$ the nodes of the discretization on $\Phi$ with the discretization parameter $h$ and set $n=H_{s} / h$ as in Fig. 5. For $u_{h} \in V_{h}\left(\Omega_{i}\right)$ we denote the average value in the interior of $\Phi$ by

$$
p_{\Phi}\left(u_{h}\right):=\frac{u_{h}\left(x_{1}\right)+u_{h}\left(x_{2}\right)+\ldots+u_{h}\left(x_{n-1}\right)}{n-1} .
$$



Figure 5. A subdomain with nodes on one side.

Lemma 5.1. Let $H_{s}=1$. Then there exists a constant $C>0$ independent of $h$ such that for every $u_{h} \in V_{h}\left(\Omega_{i}\right)$

$$
\int_{\Gamma_{i}}\left(u_{h}\right)^{2} \mathrm{~d} \Gamma \leqslant C\left(\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi}^{2}\left(u_{h}\right)\right) .
$$

Proof. First notice that due to the definition of $p_{\Phi}\left(u_{h}\right)$ and the triangle inequality,

$$
\begin{aligned}
\left(u_{h}\left(x_{k}\right)-p_{\Phi}\left(u_{h}\right)\right)^{2} & \leqslant\left(\max _{j \in\{1, \ldots, n-1\}}\left|u_{h}\left(x_{k}\right)-u_{h}\left(x_{j}\right)\right|\right)^{2} \leqslant\left(\sum_{i=1}^{n}\left|u_{h}\left(x_{i}\right)-u_{h}\left(x_{i-1}\right)\right|\right)^{2} \\
& \leqslant n \sum_{i=1}^{n}\left(u_{h}\left(x_{i}\right)-u_{h}\left(x_{i-1}\right)\right)^{2} \leqslant 2 n \int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega
\end{aligned}
$$

It follows that

$$
\begin{aligned}
\left(\int_{\Phi} u_{h} \mathrm{~d} \Gamma\right)^{2} & =h^{2}(\frac{1}{2} u_{h}\left(x_{0}\right)+\underbrace{u_{h}\left(x_{1}\right)+u_{h}\left(x_{2}\right)+\ldots+u_{h}\left(x_{n-1}\right)}_{=(n-1) p_{\Phi}\left(u_{h}\right)}+\frac{1}{2} u_{h}\left(x_{n}\right))^{2} \\
& =h^{2}\left(\frac{1}{2}\left(u_{h}\left(x_{0}\right)-p_{\Phi}\left(u_{h}\right)\right)+\frac{1}{2}\left(u_{h}\left(x_{n}\right)-p_{\Phi}\left(u_{h}\right)\right)+n p_{\Phi}\left(u_{h}\right)\right)^{2} \\
& \leqslant 3 h^{2}\left(\frac{1}{4}\left(u_{h}\left(x_{0}\right)-p_{\Phi}\left(u_{h}\right)\right)^{2}+\frac{1}{4}\left(u_{h}\left(x_{n}\right)-p_{\Phi}\left(u_{h}\right)\right)^{2}+n^{2} p_{\Phi}^{2}\left(u_{h}\right)\right) \\
& \leqslant 3 h^{2}\left(n \int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega+n^{2} p_{\Phi}^{2}\left(u_{h}\right)\right) \leqslant 3\left(\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi}^{2}\left(u_{h}\right)\right),
\end{aligned}
$$

i.e.

$$
\begin{equation*}
\left(\int_{\Phi} u_{h} \mathrm{~d} \Gamma\right)^{2} \leqslant 3\left(\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi}^{2}\left(u_{h}\right)\right) \tag{5.1}
\end{equation*}
$$

The remaining part of the proof is a direct consequence of the Poincaré inequality, the trace theorem and formula (5.1):

$$
\begin{aligned}
\int_{\Gamma_{i}}\left(u_{h}\right)^{2} \mathrm{~d} \Gamma & \leqslant C_{1}\left\|u_{h}\right\|_{H^{1}\left(\Omega_{i}\right)}^{2} \leqslant C_{2}\left(\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega+\left(\int_{\Phi} u_{h} \mathrm{~d} \Gamma\right)^{2}\right) \\
& \leqslant C\left(\int_{\Omega_{i}}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi}^{2}\left(u_{h}\right)\right)
\end{aligned}
$$

Now we again consider a cluster $\Omega$ comprising four equal square subdomains $\Omega_{1}, \ldots, \Omega_{4}$ of side $H_{s}$ such that $\bigcap_{i=1}^{4} \bar{\Omega}_{i}=\left\{x_{0}\right\}$ (as in Fig. 4),

$$
V_{h}(\Omega):=\left\{u_{h}=\left(u_{h}^{1}, u_{h}^{2}, u_{h}^{3}, u_{h}^{4}\right): u_{h}^{i} \in V_{h}\left(\Omega_{i}\right)\right\}
$$

and denote $\Phi_{i, j}=\bar{\Omega}_{i} \cap \bar{\Omega}_{j}$.
Theorem 5.1. There exists a constant $C>0$ independent of $H_{s}$ and $h$ such that for every $u_{h} \in V_{h}(\Omega)$ which satisfies

$$
\int_{\Sigma} u_{h} \mathrm{~d} \Gamma:=\sum_{i=1}^{4} \int_{\Gamma_{i}} u_{h}^{i} \mathrm{~d} \Gamma=0
$$

and for every $\Phi_{i, j} \in\left\{\Phi_{2,3}, \Phi_{3,4}, \Phi_{1,4}\right\}$ which satisfies

$$
p_{\Phi_{i, j}}\left(u_{h}^{i}\right)=p_{\Phi_{i, j}}\left(u_{h}^{j}\right)
$$

we have

$$
\int_{\Sigma}\left(u_{h}\right)^{2} \mathrm{~d} \Gamma \leqslant C H_{s} \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega
$$

Corollary 5.1. Under the above assumptions there is a constant $C>0$ independent of $H_{s}$ and $h$ such that

$$
\begin{equation*}
C \frac{h}{H_{s}} \sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} \leqslant \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega . \tag{5.2}
\end{equation*}
$$

Proof of Theorem 5.1. First assume that $H_{s}=1$ (the general case can be directly obtained using the substitution $x=H_{s} y$ ). If $v_{h}=\left(v_{h}^{1}, v_{h}^{2}, v_{h}^{3}, v_{h}^{4}\right) \in V_{h}(\Omega)$ and $p_{\Phi_{i, j}}\left(v_{h}^{i}\right)=p_{\Phi_{i, j}}\left(v_{h}^{j}\right)$ for $\Phi_{i, j} \in\left\{\Phi_{2,3}, \Phi_{3,4}, \Phi_{1,4}\right\}$, then by Lemma 5.1

$$
\begin{aligned}
p_{\Phi_{i, j}}^{2}\left(v_{h}^{i}\right) & =\left(\frac{v_{h}^{j}\left(x_{1}\right)+v_{h}^{j}\left(x_{2}\right)+\ldots+v_{h}^{j}\left(x_{n-1}\right)}{n-1}\right)^{2} \\
& \leqslant \frac{1}{n-1}\left(\left(v_{h}^{j}\left(x_{1}\right)\right)^{2}+\left(v_{h}^{j}\left(x_{2}\right)\right)^{2}+\ldots+\left(v_{h}^{j}\left(x_{n-1}\right)\right)^{2}\right) \\
& \leqslant \frac{3}{h(n-1)} \int_{\Gamma_{j}}\left(v_{h}^{j}\right)^{2} \mathrm{~d} \Gamma \leqslant \frac{6}{h n} \int_{\Gamma_{j}}\left(v_{h}^{j}\right)^{2} \mathrm{~d} \Gamma=6 \int_{\Gamma_{j}}\left(v_{h}^{j}\right)^{2} \mathrm{~d} \Gamma \\
& \leqslant C_{1}\left(\int_{\Omega_{j}}\left\|\nabla v_{h}^{j}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi_{3,4}}^{2}\left(v_{h}^{3}\right)\right),
\end{aligned}
$$

and therefore also

$$
\begin{aligned}
\int_{\Sigma}\left(v_{h}\right)^{2} \mathrm{~d} \Gamma \leqslant & C_{2}\left(\int_{\Omega_{1}}\left\|\nabla v_{h}^{1}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi_{1,4}}^{2}\left(v_{h}^{1}\right)+\int_{\Omega_{2}}\left\|\nabla v_{h}^{2}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi_{2,3}}^{2}\left(v_{h}^{2}\right)\right. \\
& \left.+\int_{\Omega_{3}}\left\|\nabla v_{h}^{3}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi_{3,4}}^{2}\left(v_{h}^{3}\right)+\int_{\Omega_{4}}\left\|\nabla v_{h}^{4}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi_{3,4}}^{2}\left(v_{h}^{4}\right)\right) \\
\leqslant & C\left(\int_{\Omega}\left\|\nabla v_{h}\right\|^{2} \mathrm{~d} \Omega+p_{\Phi_{3,4}}^{2}\left(v_{h}^{3}\right)\right)
\end{aligned}
$$

To finish the proof it is enough to choose $v_{h}:=u_{h}-p_{\Phi_{3,4}}\left(u_{h}^{3}\right)$ and notice that

$$
\begin{gathered}
\int_{\Sigma}\left(v_{h}\right)^{2} \mathrm{~d} \Gamma \geqslant \int_{\Sigma}\left(u_{h}\right)^{2} \mathrm{~d} \Gamma-2 p_{\Phi_{3,4}}\left(u_{h}^{3}\right) \underbrace{\int_{\Sigma} u_{h} \mathrm{~d} \Gamma}_{=0}=\int_{\Sigma}\left(u_{h}\right)^{2} \mathrm{~d} \Gamma \\
\int_{\Omega}\left\|\nabla v_{h}\right\|^{2} \mathrm{~d} \Omega=\int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \mathrm{~d} \Omega, \quad p_{\Phi_{3,4}}^{2}\left(v_{h}^{3}\right)=0 .
\end{gathered}
$$

Applying (5.2) and (4.1) to the harmonic function $u_{h}$, we get

$$
\begin{equation*}
C_{1} \frac{h}{H_{s}} \sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} \leqslant \int_{\Omega}\left\|\nabla u_{h}\right\|^{2} \leqslant 3 \sum_{x_{j} \in \Sigma}\left(u_{h}\left(x_{j}\right)\right)^{2} . \tag{5.3}
\end{equation*}
$$

Using the above inequalities, (3.4) and simple manipulations, we get the following corollary.

Corollary 5.2. Let the assumption of Lemma 3.1 be satisfied. Then there are constants $C_{1}$ and $C_{2}$ independent of $H_{s}$ and $h$ such that

$$
\begin{equation*}
C_{1} \leqslant \bar{\lambda}_{\min }(\widetilde{P} \widetilde{F} \widetilde{P}) \leqslant\|\widetilde{P} \widetilde{F} \widetilde{P}\| \leqslant C_{2} \frac{H_{s}}{h} \tag{5.4}
\end{equation*}
$$

## 6. Numerical experiments

To compare our estimate with the specific values, we computed the bounds on the nonzero eigenvalues of the Schur complement of $2 \times 2$ clusters with $H_{s}=1 / 4$, $H_{c}=1 / 2$, and varying discretization parameter $h$. We considered clusters obtained by gluing corners and edge averages. To compute the averages, the variables of each edge were transformed using the transformation matrix, see [10]. The results are in Table 1 and in Fig. 6.


Figure 6. Regular condition number $\bar{\kappa}(\widetilde{S})$ for $2 \times 2$ subdomains and changing $H_{s} / h$ : dark lower graph-averages, light upper graph-corners, dotted lines-fitted bounds $O\left(H_{s} / h\right)$ and $O\left(\left(1+\log \left(H_{s} / h\right)\right) H_{s} / h\right)$.

| Averages |  |  |  |  | Corners |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $H_{s} / h$ | $\kappa(S)$ | $\lambda_{\max }$ | $\lambda_{\min }$ | $\kappa(S)$ | $\lambda_{\max }$ | $\lambda_{\min }$ |  |  |
| 5 | 22.0798 | 2.6720 | 0.1210 | 38.3458 | 2.6720 | 0.0697 |  |  |
| 11 | 48.2900 | 2.7904 | 0.0578 | 118.8414 | 2.7904 | 0.0235 |  |  |
| 17 | 74.5157 | 2.8118 | 0.0377 | 211.3836 | 2.8118 | 0.0133 |  |  |
| 23 | 100.7273 | 2.8191 | 0.0280 | 311.3911 | 2.8191 | 0.0091 |  |  |
| 29 | 126.9323 | 2.8225 | 0.0222 | 416.8811 | 2.8225 | 0.0068 |  |  |

Table 1. Effective condition numbers and extremal eigenvalues for changing $H_{s} / h$ and $2 \times 2$ subdomains.

The results are in agreement with the theoretical results, in particular, it is possible to observe the nonlinear effect in the estimates for the clusters glued in corners. To illustrate the conditioning of larger clusters, we carried out the computations also
for $4 \times 4$ clusters (see Fig. 7 and Table 2). The results indicate fast deterioration of the regular condition number for the clusters glued by corners, but promising results for the clusters glued by averages. In the latter case, the condition number increased (as compared with $2 \times 2$ cluster) by some sixty percent, which indicates that improved parallelization may be very effective.


Figure 7. Effective condition number $\kappa(S)$ for $4 \times 4$ subdomains and changing $H_{s} / h$ : dark lower graph-averages, light upper graph-corners, dotted lines-fitted bounds $O\left(H_{s} / h\right)$ and $O\left(\left(1+\log \left(H_{s} / h\right)\right) H_{s} / h\right)$.

|  | Averages |  |  | Corners |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $H_{s} / h$ | $\kappa(S)$ | $\lambda_{\max }$ | $\lambda_{\min }$ | $\kappa(S)$ | $\lambda_{\max }$ | $\lambda_{\min }$ |
| 5 | 51.418 | 2.672 | 0.052 | 136.633 | 2.672 | 0.020 |
| 11 | 102.526 | 2.790 | 0.027 | 442.471 | 2.790 | 0.006 |
| 17 | 153.916 | 2.812 | 0.018 | 793.049 | 2.812 | 0.004 |
| 23 | 205.364 | 2.819 | 0.014 | 1170.667 | 2.819 | 0.002 |

Table 2. Effective condition numbers and extremal eigenvalues for changing $H_{s} / h$ and $4 \times 4$ subdomains.

It seems that H-TFETI without preconditioners can be very effective method for the nonlinear problems but more research is necessary.

## 7. Conclusions

We have established bounds on the regular condition number of the Schur complements of floating $2 \times 2$ clusters arising from gluing the subdomains in nodes and averages. In the first case, we showed that our estimates cannot be qualitatively improved. We considered the Schur complements of the stiffness matrices arising from the discretization of problems governed by the Laplace operator. The research was motivated by the effort to understand massively parallel H-TFETI algorithms for the solution of elliptic problems described by variational inequalities which have already proved to be effective for the problems discretized by billions of nodal variables [4]. It seems that the application of H-TFETI can overcome the bottleneck of TFETI associated with the dimension of the coarse grid and can increase the scalability from the current tens of thousands of cores to hundreds of thousands of cores.

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