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ON ISOTROPIC TENSORS

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1. Let M be an oriented Riemannian two-dimensional manifold of class C^∞ with the boundary ∂M . Let $\{U_\alpha\}$ be an open covering of M such that there is, on each U_α , a field of orthonormal frames $\{v_1, v_2\}$ with $v_1, v_2 \in T(M)$; let $\{\omega^1, \omega^2\}$ be the dual coframes. On U_α , the metric form of M is given by

$$(1) \quad g = (\omega^1)^2 + (\omega^2)^2.$$

Let the 1-form ω_1^2 be defined by

$$(2) \quad d\omega^1 = -\omega^2 \wedge \omega_1^2, \quad d\omega^2 = \omega^1 \wedge \omega_1^2.$$

Then the Gauss curvature of M is given by

$$(3) \quad d\omega_1^2 = -K\omega^1 \wedge \omega^2.$$

On M , be given a quadratic form S ; its expression in U_α be

$$(4) \quad S = S_{11}(\omega^1)^2 + 2S_{12}\omega^1\omega^2 + S_{22}(\omega^2)^2.$$

The covariant derivatives $S_{ijk} = S_{jik}$ of the symmetric tensor S_{ij} (with respect to the coframe ω^1, ω^2) be defined by

$$(5) \quad \begin{aligned} dS_{11} - 2S_{12}\omega_1^2 &= S_{111}\omega^1 + S_{112}\omega^2, \\ dS_{12} + (S_{11} - S_{22})\omega_1^2 &= S_{121}\omega^1 + S_{122}\omega^2, \\ dS_{22} + 2S_{12}\omega_1^2 &= S_{221}\omega^1 + S_{222}\omega^2. \end{aligned}$$

Then

$$(6) \quad J(S) = S_{121}(S_{112} - S_{222}) + S_{122}(S_{221} - S_{111})$$

is an invariant. We are going to prove this auxiliary result as well as the following

Theorem. *Let the data be as above. Further, suppose that: (i) $K > 0$ on M ; (ii) $J(S) \geq 0$ on M ; (iii) there is a function $\lambda : \partial M \rightarrow \mathcal{R}$ such that $S = \lambda g$ on ∂M . Then there is a function $\Lambda : M \rightarrow \mathcal{R}$ such that $S = \Lambda g$ on M .*

I am going to prove this result by means of an integral formula which is a generalization of an integral formula introduced by A. ŠVEC [1]. Švec constructed a certain 1-form τ on surfaces of E^3 ; he claims this form to be invariant without proving it. In what follows, I am going to prove this; in fact, I am going to show that τ is invariant on oriented surfaces of E^2 . Because of that, I restrict myself to the case that M is a surface of E^3 ; it is easy to see that the general proof of our Theorem is exactly the same. From a general point of view, remark that I get an integral formula for a symmetric tensor without supposing this tensor to be of the Ricci type; see, p. ex., [2].

2. Let $M \subset E^3$ be a surface of class C^∞ . Let us cover M by domains U_α such that there is, on each U_α , a field of orthonormal frames $\{m; v_1, v_2, v_3\}$ with $m \in U_\alpha \subset M$; $v_1, v_2 \in T_m(M)$. Then

$$(7) \quad dm = \omega^1 v_1 + \omega^2 v_2,$$

$$dv_1 = \omega_1^2 v_2 + \omega_1^3 v_3, \quad dv_2 = -\omega_1^2 v_1 + \omega_2^3 v_3, \quad dv_3 = -\omega_1^3 v_1 - \omega_2^3 v_2;$$

$$(8) \quad \begin{aligned} d\omega^i &= \omega^j \wedge \omega_j^i, & d\omega_j^i &= \omega_i^k \wedge \omega_k^j \\ & \text{with } \omega^3 &= 0, & \omega_i^j + \omega_j^i = 0. \end{aligned}$$

From $\omega^3 = 0$, we get the existence of functions a, b, c such that

$$(9) \quad \omega_1^3 = a\omega^1 + b\omega^2, \quad \omega_2^3 = b\omega^1 + c\omega^2;$$

further,

$$(10) \quad \begin{aligned} da - 2b\omega_1^2 &= a\omega^1 + b\omega^2, \\ db + (a - c)\omega_1^2 &= b\omega^1 + c\omega^2, \\ dc + 2b\omega_1^2 &= c\omega^1 + d\omega^2. \end{aligned}$$

As always, let

$$(11) \quad 2H = a + c, \quad K = ac - b^2$$

define the mean and Gauss curvature resp.

Let $\{m; w_1, w_2, w_3\}$ be another field of orthonormal frames on U_α , i.e.,

$$(12) \quad \begin{aligned} v_1 &= \varepsilon_1 \cos \varphi \cdot w_1 - \varepsilon_1 \sin \varphi \cdot w_2, & v_2 &= \sin \varphi \cdot w_1 + \cos \varphi \cdot w_2, \\ v_3 &= \varepsilon_2 w_3; & \varepsilon_1^2 &= \varepsilon_2^2 = 1 \end{aligned}$$

and φ a function on U_α . Further, let

$$(13) \quad dm = \tau^1 w_1 + \tau^2 w_2,$$

$$dw_1 = \tau_1^2 w_2 + \tau_1^3 w_3, \quad dw_2 = -\tau_1^2 w_1 + \tau_2^3 w_3, \quad dw_3 = -\tau_1^3 w_1 - \tau_2^3 w_2;$$

$$(14) \quad \tau_1^3 = a^* \tau^1 + b^* \tau^2, \quad \tau_2^3 = b^* \tau^1 + c^* \tau^2;$$

$$da^* - 2b^* \tau_1^2 = \alpha^* \tau^1 + \beta^* \tau^2, \quad \text{etc.}$$

From (7), (13) and (12), we get

$$(15) \quad \tau^1 = \varepsilon_1 \cos \varphi \cdot \omega^1 + \sin \varphi \cdot \omega^2, \quad \tau^2 = -\varepsilon_1 \sin \varphi \cdot \omega^1 + \cos \varphi \cdot \omega^2,$$

i.e.,

$$(16) \quad \omega^1 = \varepsilon_1 \cos \varphi \cdot \tau^1 - \varepsilon_1 \sin \varphi \cdot \tau^2, \quad \omega^2 = \sin \varphi \cdot \tau^1 + \cos \varphi \cdot \tau^2.$$

Further, it is easy to see that

$$(17) \quad \tau_1^2 - d\varphi = \varepsilon_1 \omega_1^2;$$

$$(18) \quad \tau_1^3 = \varepsilon_1 \varepsilon_2 \cos \varphi \cdot \omega_1^3 + \varepsilon_2 \sin \varphi \cdot \omega_2^3,$$

$$\tau_2^3 = -\varepsilon_1 \varepsilon_2 \sin \varphi \cdot \omega_1^3 + \varepsilon_2 \cos \varphi \cdot \omega_2^3;$$

$$(19) \quad \omega_1^3 = \varepsilon_1 \varepsilon_2 \cos \varphi \cdot \tau_1^3 - \varepsilon_1 \varepsilon_2 \sin \varphi \cdot \tau_2^3,$$

$$\omega_2^3 = \varepsilon_2 \sin \varphi \cdot \tau_1^3 + \varepsilon_2 \cos \varphi \cdot \tau_2^3$$

and

$$(20) \quad a^* = \varepsilon_2 \cos^2 \varphi \cdot a + 2\varepsilon_1 \varepsilon_2 \sin \varphi \cos \varphi \cdot b + \varepsilon_2 \sin^2 \varphi \cdot c,$$

$$b^* = -\varepsilon_2 \sin \varphi \cos \varphi \cdot a - \varepsilon_1 \varepsilon_2 \sin^2 \varphi \cdot b + \varepsilon_1 \varepsilon_2 \cos^2 \varphi \cdot b + \varepsilon_2 \sin \varphi \cos \varphi \cdot c,$$

$$c^* = \varepsilon_2 \sin^2 \varphi \cdot a - 2\varepsilon_1 \varepsilon_2 \sin \varphi \cos \varphi \cdot b + \varepsilon_2 \cos^2 \varphi \cdot c.$$

Thus

$$(21) \quad H^* = \varepsilon_2 H, \quad K^* = K,$$

the well known results. By a little more complicated calculation, we obtain

$$(22) \quad \alpha^* = \varepsilon_1 \varepsilon_2 \cos^3 \varphi \cdot \alpha + 3\varepsilon_2 \sin \varphi \cos^2 \varphi \cdot \beta + 3\varepsilon_1 \varepsilon_2 \sin^2 \varphi \cos \varphi \cdot \gamma + \varepsilon_2 \sin^3 \varphi \cdot \delta,$$

$$\beta^* = -\varepsilon_1 \varepsilon_2 \sin \varphi \cos^2 \varphi \cdot \alpha + (\varepsilon_2 \cos^3 \varphi - 2\varepsilon_2 \sin^2 \varphi \cos \varphi) \beta + (2\varepsilon_1 \varepsilon_2 \sin \varphi \cos^2 \varphi - \varepsilon_1 \varepsilon_2 \sin^3 \varphi) \gamma + \varepsilon_2 \sin^2 \varphi \cos \varphi \cdot \delta,$$

$$\begin{aligned}
\gamma^* &= \varepsilon_1 \varepsilon_2 \sin^2 \varphi \cos \varphi \cdot \alpha + (\varepsilon_2 \sin^3 \varphi - 2\varepsilon_2 \sin \varphi \cos^2 \varphi) \beta - \\
&\quad - (2\varepsilon_1 \varepsilon_2 \sin^2 \varphi \cos \varphi - \varepsilon_1 \varepsilon_2 \cos^3 \varphi) \gamma + \varepsilon_2 \sin \varphi \cos^2 \varphi \cdot \delta, \\
\delta^* &= -\varepsilon_1 \varepsilon_2 \sin^3 \varphi \cdot \alpha + 3\varepsilon_2 \sin^2 \varphi \cos \varphi \cdot \beta - 3\varepsilon_1 \varepsilon_2 \sin \varphi \cos^2 \varphi \cdot \gamma + \\
&\quad + \varepsilon_2 \cos^3 \varphi \cdot \delta
\end{aligned}$$

using (10), (14), (17) and (20).

Now, introduce the 1-form

$$\begin{aligned}
(23) \quad \Phi &= R_1 \omega^1 + R_2 \omega^2, \\
R_1 &:= (a - c) \beta + b(\gamma - \alpha), \quad R_2 := (a - c) \gamma + b(\delta - \beta),
\end{aligned}$$

this being exactly the form τ introduced in [1]. From (14) and (16), we get

$$(24) \quad R_1^* = \cos \varphi \cdot R_1 + \varepsilon_1 \sin \varphi \cdot R_2, \quad R_2^* = -\sin \varphi \cdot R_1 + \varepsilon_1 \cos \varphi \cdot R_2,$$

i.e.,

$$(25) \quad \Phi^* = \varepsilon_1 \Phi.$$

Thus we have proved that Φ is an invariant form on oriented surfaces.

3. Now, consider the quadratic form (10) on our surface M . From

$$(26) \quad S = S_{11}(\omega^1)^2 + 2S_{12}\omega^1\omega^2 + S_{22}(\omega^2)^2 = S_{11}^*(\tau^1)^2 + 2S_{12}^*\tau^1\tau^2 + S_{22}^*(\tau^2)^2$$

and (15), (16), we get

$$\begin{aligned}
(27) \quad S_{11}^* &= \cos^2 \varphi \cdot S_{11} + 2\varepsilon_1 \sin \varphi \cos \varphi \cdot S_{12} + \sin^2 \varphi \cdot S_{22}, \\
S_{12}^* &= -\sin \varphi \cos \varphi \cdot S_{11} - \varepsilon_1 \sin^2 \varphi \cdot S_{12} + \varepsilon_1 \cos^2 \varphi \cdot S_{12} + \\
&\quad + \sin \varphi \cos \varphi \cdot S_{22}, \\
S_{22}^* &= \sin^2 \varphi \cdot S_{11} - 2\varepsilon_1 \sin \varphi \cos \varphi \cdot S_{12} + \cos^2 \varphi \cdot S_{22}.
\end{aligned}$$

From (11) and similar equations

$$\begin{aligned}
(28) \quad dS_{11}^* - 2S_{12}^*\tau_1^2 &= S_{111}^*\tau^1 + S_{112}^*\tau^2, \\
dS_{12}^* + (S_{11}^* - S_{22}^*)\tau_1^2 &= S_{121}^*\tau^1 + S_{122}^*\tau^2, \\
dS_{22}^* + 2S_{12}^*\tau_1^2 &= S_{221}^*\tau^1 + S_{222}^*\tau^2
\end{aligned}$$

we obtain

$$\begin{aligned}
&\cos^2 \varphi \cdot S_{111} + 2\varepsilon_1 \sin \varphi \cos \varphi \cdot S_{121} + \sin^2 \varphi \cdot S_{221} = \\
&= \varepsilon_1 \cos \varphi \cdot S_{111}^* - \varepsilon_1 \sin \varphi \cdot S_{112}^*,
\end{aligned}$$

$$\begin{aligned}
& \cos^2 \varphi \cdot S_{112} + 2\varepsilon_1 \sin \varphi \cos \varphi \cdot S_{122} + \sin^2 \varphi \cdot S_{222} = \\
& = \sin \varphi \cdot S_{111}^* + \cos \varphi \cdot S_{112}^*, \\
& -\sin \varphi \cos \varphi \cdot S_{111} - \varepsilon_1 \sin^2 \varphi \cdot S_{121} + \varepsilon_1 \cos^2 \varphi \cdot S_{121} + \sin \varphi \cos \varphi \cdot S_{221} = \\
& = \varepsilon_1 \cos \varphi \cdot S_{121}^* - \varepsilon_1 \sin \varphi \cdot S_{122}^*, \\
& -\sin \varphi \cos \varphi \cdot S_{112} - \varepsilon_1 \sin^2 \varphi \cdot S_{122} + \varepsilon_1 \cos^2 \varphi \cdot S_{122} + \sin \varphi \cos \varphi \cdot S_{222} = \\
& = \sin \varphi \cdot S_{121}^* + \cos \varphi \cdot S_{122}^*, \\
& \sin^2 \varphi \cdot S_{111} - 2\varepsilon_1 \sin \varphi \cos \varphi \cdot S_{121} + \cos^2 \varphi \cdot S_{221} = \\
& = \varepsilon_1 \cos \varphi \cdot S_{121}^* - \varepsilon_1 \sin \varphi \cdot S_{222}^*, \\
& \sin^2 \varphi \cdot S_{112} - 2\varepsilon_1 \sin \varphi \cos \varphi \cdot S_{122} + \cos^2 \varphi \cdot S_{222} = \\
& = \sin \varphi \cdot S_{221}^* + \cos \varphi \cdot S_{222}^*,
\end{aligned}$$

i.e.,

$$\begin{aligned}
(29) \quad S_{111}^* &= \varepsilon_1 \cos^2 \varphi \cdot S_{111} + \sin \varphi \cos^2 \varphi \cdot S_{112} + 2 \sin \varphi \cos^2 \varphi \cdot S_{121} + \\
& \quad + 2\varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{122} + \varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{221} + \sin^3 \varphi \cdot S_{222}, \\
S_{112}^* &= -\varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{111} + \cos^3 \varphi \cdot S_{112} - 2 \sin^2 \varphi \cos \varphi \cdot S_{121} + \\
& \quad + 2\varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{122} - \varepsilon_1 \sin^3 \varphi \cdot S_{221} + \sin^2 \varphi \cos \varphi \cdot S_{222}, \\
S_{121}^* &= -\varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{111} - \sin^2 \varphi \cos \varphi \cdot S_{112} - \\
& \quad - \sin^2 \varphi \cos \varphi \cdot S_{121} + \cos^3 \varphi \cdot S_{121} - \varepsilon_1 \sin^3 \varphi \cdot S_{122} + \\
& \quad + \varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{122} + \varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{221} + \\
& \quad + \sin^2 \varphi \cos \varphi \cdot S_{222}, \\
S_{122}^* &= \varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{111} - \sin \varphi \cos^2 \varphi \cdot S_{112} + \sin^3 \varphi \cdot S_{121} - \\
& \quad - \sin \varphi \cos^2 \varphi \cdot S_{121} - \varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{122} + \varepsilon_1 \cos^3 \varphi \cdot S_{122} - \\
& \quad - \varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{221} + \sin \varphi \cos^2 \varphi \cdot S_{222}, \\
S_{221}^* &= \varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{111} + \sin^3 \varphi \cdot S_{112} - 2 \sin \varphi \cos^2 \varphi \cdot S_{121} - \\
& \quad - 2\varepsilon_1 \sin^2 \varphi \cos \varphi \cdot S_{122} + \varepsilon_1 \cos^3 \varphi \cdot S_{221} + \sin \varphi \cos^2 \varphi \cdot S_{222}, \\
S_{222}^* &= -\varepsilon_1 \sin^3 \varphi \cdot S_{111} + \sin^2 \varphi \cos \varphi \cdot S_{112} + 2 \sin^2 \varphi \cos \varphi \cdot S_{121} - \\
& \quad - 2\varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{122} - \varepsilon_1 \sin \varphi \cos^2 \varphi \cdot S_{221} + \cos^3 \varphi \cdot S_{222}.
\end{aligned}$$

From this

$$(30) \quad \begin{aligned} S_{121}(S_{112} - S_{222}) + S_{122}(S_{221} - S_{111}) &= \\ &= S_{121}^*(S_{112}^* - S_{222}^*) + S_{122}^*(S_{221}^* - S_{111}^*), \end{aligned}$$

i.e.,

$$(31) \quad J(S) = J^*(S).$$

Consider the 1-form

$$(32) \quad \begin{aligned} \Psi &= T_1 \omega^1 + T_2 \omega^2, \\ T_1 &:= (S_{11} - S_{22}) S_{121} + S_{12}(S_{221} - S_{111}), \\ T_2 &:= (S_{11} - S_{22}) S_{122} + S_{12}(S_{222} - S_{112}). \end{aligned}$$

From (27) and (30),

$$(33) \quad T_1^* = \cos \varphi \cdot T_1 + \varepsilon_1 \sin \varphi \cdot T_2, \quad T_2^* = -\sin \varphi \cdot T_1 + \varepsilon_1 \cos \varphi \cdot T_2,$$

i.e.,

$$(34) \quad \Psi^* = \varepsilon_1 \Psi,$$

and the form Ψ (32) is thus invariant on oriented surfaces.

The equations (11) imply

$$(35) \quad \begin{aligned} \{dS_{111} - (S_{112} + 2S_{121})\omega_1^2\} \wedge \omega^1 + \\ + \{dS_{112} + (S_{111} - 2S_{122})\omega_1^2\} \wedge \omega^2 &= 2S_{12}K\omega^1 \wedge \omega^2, \\ \{dS_{121} + (S_{111} - S_{122} - S_{221})\omega_1^2\} \wedge \omega^1 + \\ + \{dS_{122} + (S_{112} + S_{121} - S_{222})\omega_1^2\} \wedge \omega^2 &= (S_{22} - S_{11})K\omega^1 \wedge \omega^2, \\ \{dS_{221} + (2S_{121} - S_{222})\omega_1^2\} \wedge \omega^1 + \\ + \{dS_{222} + (2S_{122} + S_{221})\omega_1^2\} \wedge \omega^2 &= -2S_{12}K\omega^1 \wedge \omega^2, \end{aligned}$$

and we get the existence of functions T_1, \dots, T_9 such that

$$(36) \quad \begin{aligned} dS_{111} - (S_{112} + 2S_{121})\omega_1^2 &= T_1\omega^1 + (T_2 - S_{12}K)\omega^2, \\ dS_{112} + (S_{111} - 2S_{122})\omega_1^2 &= (T_2 + S_{12}K)\omega^1 + T_3\omega^2, \\ dS_{121} + (S_{111} - S_{122} - S_{221})\omega_1^2 &= T_4\omega^1 + (T_5 + S_{11}K)\omega^2, \\ dS_{122} + (S_{112} + S_{121} - S_{222})\omega_1^2 &= (T_5 + S_{22}K)\omega^1 + T_6\omega^2, \\ dS_{221} + (2S_{121} - S_{222})\omega_1^2 &= T_7\omega^1 + (T_8 + S_{12}K)\omega^2, \\ dS_{222} + (2S_{122} + S_{221})\omega_1^2 &= (T_8 - S_{12}K)\omega^1 + T_9\omega^2. \end{aligned}$$

By means of these formulas, we get

$$(37) \quad d\Psi = -\{2S_{121}(S_{112} - S_{222}) + 2S_{122}(S_{221} - S_{111}) + \\ + ((S_{11} - S_{22})^2 + 4S_{12}^2)K\} \omega^1 \wedge \omega^2.$$

The Stokes formula $\int_{\partial M} \Psi = \int_M d\Psi$ reads now

$$(38) \quad \int_{\partial M} \{(S_{11} - S_{22})(S_{121}\omega^1 + S_{122}\omega^2) + \\ + S_{12}(S_{221}\omega^1 - S_{111}\omega^1 + S_{222}\omega^2 - S_{112}\omega^2)\} = \\ = - \int_M \{2J(S) + ((S_{11} - S_{22})^2 + 4S_{12}^2)K\} \omega^1 \wedge \omega^2.$$

The proof of our Theorem follows easily. On the boundary ∂M , we have $S_{11} = S_{22} = \lambda$, $S_{12} = 0$, and the left-hand side of (38) is thus equal to zero. Because of $K > 0$ and $J(S) \geq 0$, we get $(S_{11} - S_{22})^2 + 4S_{12}^2 = 0$, i.e., $S_{11} - S_{22} = S_{12} = 0$. We are finished setting $S_{11} = S_{22} = \lambda$.

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