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On the question of the possible rotation of the local cluster.

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Dedicated to Professor František Nušl on the occasion of his seventieth anniversary December 3, 1937.

On the base of the study of the space velocities of 910 stars of B type the author shows that the suplementary rotation of the so-called local cluster about a centre situated at \( l = 237^\circ \) in the constellation of Carina in a distance of some hundred parsecs only does not exist.

In my previous work\(^1\) I gave new observational evidence of the galactic rotation. But H. Mineur showed\(^2\) that my demonstration did not settle the question of the possible sub-rotation about the second centre, i.e. the centre of the local cluster, which must be placed some hundred parsecs only from the Sun in the direction of gal. longitude \( l = 237^\circ \). I have shown in a comment on a paper by Edmondson,\(^3\) that this question may be solved by studying the motions of stars, which are placed in belts of space oriented towards gal. longitudes \( l = 147^\circ \) and \( l = 327^\circ \). I also mentioned, that I would return to this subject.

Consider a system of coordinates \( X, Y \) in the galactic plane, the origin being the Sun. The axes \( X, Y \) are oriented in the usual sense. Let \( V_\odot \) denote the instantaneous rotational velocity of the Sun, \( V_* \) the mean rotational motion of the stars in the solar vicinity, and \( V_r \) the velocity of these stars relative to the Sun. If we take a new system of coordinates \( X', Y' \) in which \( Y' \) is in the direction of \( V_* \) then \( X' \) is directed to the galactic centre, i.e. approximately to \( l = 327^\circ \). By this supposition we obtain the following relations:

\[
\begin{align*}
V_\odot \cos \alpha_\odot &= V_* + V_r \cos \beta \\
V_r \sin \beta &= x'; \quad V_r \cos \beta = y'
\end{align*}
\]

\(^{(1)}\)

\(^1\) J. M. Mohr, M. N. 92 (1932), 583.
\(^2\) H. Mineur, Bull. astron., 7 (1934), 397.
\(^3\) Edmonson, Astr. Journal, 44 (1935), No 1016, 16.
where $\beta$ is the angular distance between the Apex and the point towards which the mean rotational motion of the stars is directed, and $\alpha_0$ is the angle between the direction of the rotational motion of the Sun and the direction of the mean rotational motion of the stars. It is also easy to see that following relations must exist:

\[
\begin{align*}
V_r^2 &= x^2 + y^2, \\
x' &= x \cos \alpha' - y \sin \alpha', \\
y' &= x \sin \alpha' + y \cos \alpha',
\end{align*}
\]

(2)

where $\alpha' = 360^\circ - 327^\circ = 33^\circ$.

From my previous discussion of $B$ stars,\textsuperscript{4) which I use also in this paper, it resulted that

\[V_r = 19.12 \text{ kmsec}^{-1}\]

If we accept the value $V_0 = 300 \text{ kmsec}^{-1}$, towards $l_0 = 55^\circ$ then $V_\alpha = 284.88 \text{ kmsec}^{-1}$, towards $l_\alpha = 56^\circ 54' 24.7"$.

If as H. Mineur supposes, a secondary rotation does exist (i.e. if there does exist a second centre of rotation in the direction $l = 237^\circ$ distant only a few hundred parsecs from the Sun), then for stars, whose coordinates of positions $Y'$ are less than zero, this rotation must be seen in the $x'$ (components of their space velocities). Besides the vector $V_1$ of the fundamental rotation, there must be present also the vector $V_2$ of the second rotation, so that the resultant of both vectors in different positions of space must be different in magnitude and in the direction.

Let $V_1$ denote the vectors of the fundamental rotation. The centre of this rotation is placed in the direction $l = 327^\circ$. $V_2$ are the vectors of the possible supplementary rotation about $l = 237^\circ$ for the group of stars with $Y' < 0$. The resultant motion of stars in the group I. of stars is $V'_1$, in the group II. $V'_2$ and in the group III. $V'_3$. The angle between the vectors $V_1$ and $V_2$ is $\alpha'_1$, between the vectors $V_1$ and $V_3$ is $\alpha'_2$, etc. Denote moreover by $V_0$ the instantaneous velocity of the Sun, which is compounded also from the vector of the fundamental rotation and the vector of sub-rotation for $Y' = 0$. (This latter vector is sensibly smaller than in the space where $Y' < 0$).

For $Y' = 0$, i.e. in the vicinity of the Sun, we get

\[V_0 \sin \alpha_0 + V_2 \sin \alpha_2 = x',\]

where $\alpha_2$ is the angle between the vectors $V_1$ and $V_2$.

But for $Y' < 0$ we get similarly

\[V_0 \sin \alpha_0 + V'_2 \sin \alpha'_2 = x'_2.\]

\textsuperscript{4) Publ. de l'Institut Astron. de l'Université Charles de Prague, No. 19, 1936.}
and
\[ V \sin \alpha + V_1 \sin \alpha_1 = x_1; \quad V \sin \alpha + V_2 \sin \alpha_2 = x_2, \]
where
\[ \alpha_1 \leq \alpha_2 \leq \alpha_3, \]
and
\[ V_1 \leq V_2 \leq V_3, \]
either if the group I. is one of groups of stars situated in positive \( X' \) and the group III. in negative \( X' \), or vice versa. But if the groups I. and III. of stars are at the same distance from the axis \( Y' \), so that
\[ |X'_1| = |X'_3|, \]
then the ratio of the change of \( \alpha'_1 \) compared with \( \alpha'_3 \) and of \( V'_1 \) to \( V'_3 \) is given by the relation
\[ \frac{V'_s}{V'''_s} = \frac{\sin \alpha'_3}{\sin \alpha'_1}, \]
so that
\[ x'_1 = x'_3 > x'_2, \]
because the components \( x'_1, x'_2, x'_3 \) are negative.

If we suppose for simplicity that the vector of the fundamental rotations is independent of \( \bar{X}' \) (this supposition may be made owing to the great distance of the centre of rotation), then considering groups of stars placed in \( Y' < 0 \), the component \( x' \) of the space velocity for \( X' = 0 \) must reach a maximum negative value. For values \( \bar{X}' \geq 0 \) this component \( x' \) must gradually reach lower negative values.

The values \( x'_1, x'_2, x'_3 \) are easily obtainable if we know the components \( x, y, z \) of space velocities of the stars as shown by formulas (1) and (2).

But how is realised the inequality (3) by the computation? If we divide the space round the Sun in three parts so, that the first part contain all stars, which \( Y' > + 150 \) parsecs, the second part contain all stars, which \( Y' = \pm 150 \) parsecs, the third part contain all stars, which \( Y' < -150 \) parsecs, then the computed values of \( x' \) for different \( X' \) are given in Table I.

Table I.

| \( X' \) (parsec) | \(-502,66\) | \(-147,81\) | \(-49,59\) | \(+50,48\) | \(+189,00\) | \(+419,16\) |
| \( Y' \) (parsec) | \(+379,03\) | \(+344,31\) | \(+302,27\) | \(+284,98\) | \(+364,84\) | \(+393,96\) |
| \( x' \) (kmsec\(^{-1}\)) | \(+2,31\) | \(+5,39\) | \(+3,41\) | \(+3,48\) | \(+3,03\) | \(+17,08\) |
| Number of stars | 17 | 39 | 32 | 37 | 27 | 51 |

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The Table I. shows, that $x'$ remains nearly constant in the groups where $Y' > +150$ parsecs and $Y' = \pm 150$ parsecs. The dependence of the $x'$ on $X'$ in the group $Y' < -150$ parsecs does not show in the supplementary rotation about the centre $l = 237^\circ$. Although there are a smaller number of stars in this group, especially for large positive $X'$, it can be seen, from the values $x'$ between $X' = -491.00$ parsecs and $X' = +45.34$, where we find a homogeneous number of stars, that the behaviour of the values $x'$ is just in the opposite sense to that which would be, if the supplementary rotation did exist.

This ascertained distribution of the components $x'$ in various parts of space has of course a great significance for, for instance, the $K$ term. Because the distribution of the components $x'$ in the parts of space, the $Y'$ of which is $Y' > +150$ and $Y' = \pm 150$ parsecs, is regular, it is clear that we should have from this material of stars the possibility to obtain a value of the $K$ term not affected by the perturbing influence of the star-streaming. In the space where $Y' < -150$ parsecs the observed dependence of $x'$ on $X'$ is probably caused not only by the rotation of the Galaxy but probably by the star-streaming. By this star-streaming the velocity-ellipsoids in those parts of space are deformed.

Therefore if we would find the real value of the $K$ term we must separate all deformations from the observed velocity-ellipsoids. I myself tried a case unreal, but for the time being important for the knowledge of the real velocity-ellipsoids in different parts of space in order to find in the future what is the real form of

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the undeformed velocity-ellipsoid in different galactic longitudes and distances from the Sun.

At the end, I will not affirm that the material here used is sufficient enough. The want of known space velocities of \( B \) stars in the vicinity of \( l = 237^\circ \) is compensated however by the knowledge of a larger number of the proper motions of stars which we are able to analyze in a similar way — as it was shown in this paper. I hope therefore, that in a time as far as may be seen it would be possible to give a new proof about the unexistence of the suplementary rotation by the proper motions of stars placed in the vicinity of the supposed centre, which \( l = 237^\circ, b = 0 \) approximately.

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K otázce možné rotace lokálního shluku hvězd v okolí Slunce.

(Obsah předešlého článku.)

Vyšetřováním prostorových rychlostí 910 hvězd spektrálního typu \( B \) způsobem autorem naznačeným se ukazuje, že nemůže existovatí rotace místního shluku hvězd, jehož střed by se podle některých autorů měl nalézatí přibližně v rovině Mléčné dráhy a v okolí galaktické délky \( l = 237^\circ \) ve vzdálenosti jen několika set parseků.