

Acta Universitatis Palackianae Olomucensis. Facultas Rerum  
Naturalium. Mathematica

---

Ján Andres

Asymptotic properties of solutions of a certain third-order differential equation with an oscillatory restoring term

*Acta Universitatis Palackianae Olomucensis. Facultas Rerum Naturalium. Mathematica*, Vol. 27 (1988), No. 1, 201--210

Persistent URL: <http://dml.cz/dmlcz/120193>

**Terms of use:**

© Palacký University Olomouc, Faculty of Science, 1988

Institute of Mathematics of the Academy of Sciences of the Czech Republic provides access to digitized documents strictly for personal use. Each copy of any part of this document must contain these *Terms of use*.



This paper has been digitized, optimized for electronic delivery and stamped with digital signature within the project *DML-CZ: The Czech Digital Mathematics Library* <http://project.dml.cz>

---

Katedra matematické analýzy a numerické matematiky  
přírodovědecké fakulty Univerzity Palackého v Olomouci  
Vedoucí katedry: Doc.RNDr. Jindřich Palát, CSc.

**ASYMPTOTIC PROPERTIES OF SOLUTIONS  
OF A CERTAIN THIRD-ORDER DIFFERENTIAL  
EQUATION WITH  
AN OSCILLATORY RESTORING TERM**

JAN ANDRES

(Received January 15, 1987)

1. Considering the equation

$$x'''' + ax'' + g(x)x' + h(x) = p(t), \quad (1)$$

where  $a > 0$  is a constant,  $p(t) \in C(0, \infty)$ ,  $g(x), h(x) \in C^1(-\infty, \infty)$  and  $h(x)$  is an oscillatory function in the whole interval  $(-\infty, \infty)$  with isolated zero points  $\bar{x}$ , K.E.Swick [1] has proved under

$$\int_0^{\infty} |p(t)| dt < \infty \quad (2)$$

the following

Theorem 0. If there exist such positive constants  $b, c$  that the assumptions

- 1)  $\frac{1}{x} \int_0^x g(s) ds \geq b,$
- 2)  $h'(x) \leq c$  with  $c < ab,$
- 3)  $h(x) \operatorname{sgn} x \geq 0$

are fulfilled for all  $x \in (-\infty, \infty),$  then all solutions  $x(t)$  of (1) are bounded satisfying

$$\lim_{t \rightarrow \infty} x(t) = \bar{x}, \quad \lim_{t \rightarrow \infty} x'(t) = \lim_{t \rightarrow \infty} x''(t) = 0. \quad (3)$$

The aim of the present paper is to make the above result more precise in two directions, namely (i) condition 2) may be localized to the origin and (ii)  $h(x) \operatorname{sgn} x$  may run below the  $x$ -axis for  $g(x) \equiv b > 0,$  both, when  $|h(x)|$  is bounded everywhere.

2. Hence let us assume

$$\limsup_{|x| \rightarrow \infty} |h(x)| < \infty \quad (4)$$

and recall several well-known results at first.

Lemma 1. Let all solutions  $x(t)$  of (1) be bounded together with their derivatives  $x'(t), x''(t).$  If (3) is satisfied for those of autonomous equation (1) (i.e.  $p(t) \equiv 0$ ), then (3) is also true for all solutions  $x(t)$  of (1) with (2).

P r o o f. The above assertion is a direct consequence of the Markus-Opial-Yoshizawa theorem [2, p.59], when specified to (1).

Lemma 2. If there exists such an  $h$ -neighbourhood of the root  $\bar{x}$  of  $h(x)$  in (1) with  $p(t) \equiv 0$  that conditions

- 2')  $ag(x) - h'(x) \geq \delta > 0 \quad (\delta - \text{const.}),$
- 3')  $h'(x) > 0,$
- 4)  $g(\bar{x}) = 0$

are satisfied for  $0 < |x - \bar{x}| < h$ , then  $\bar{x}$  is asymptotically stable.

For the proof see [3].

Remark 1. It can be readily checked that the basin of attractivity due to  $\bar{x}$  is determined by  $|g'(x)x'| \leq \delta_0$  ( $\delta_0$  - small enough constant), while for  $g(x) \equiv b > 0$  even by  $h(x)\text{sgn}(x-\bar{x}) > 0$ , because of the form of Liapunov's function employed in [3].

Lemma 3. If there exist such positive constants  $b, G$  that condition

$$1') \quad b \leq g(x) \leq G < a^2$$

is satisfied for all  $x \in (-\infty, \infty)$  together with (4) and

$$5) \quad \limsup_{t \rightarrow \infty} |p(t)| < \infty,$$

$$6) \quad \limsup_{t \rightarrow \infty} \left| \int_0^t p(s) ds \right| < \infty$$

then there exists also a constant  $D'$  such that all solutions  $x(t)$  of (1) satisfy

$$\limsup_{t \rightarrow \infty} (|x'(t)| + |x''(t)|) < D'. \quad (5)$$

For the proof see [4].

Lemma 4. If there exists (finite)

$$\lim_{t \rightarrow \infty} x(t)$$

of (1) satisfying (4), (5) and 5) of Lemma 3, then there is also

$$\lim_{t \rightarrow \infty} x'(t) = \lim_{t \rightarrow \infty} x''(t) = 0. \quad (3')$$

P r o o f. This assertion follows directly from the theorem introduced in [5, p.141], because of  $\limsup_{t \rightarrow \infty} |x''''(t)| < \infty$ .

Lemma 5. If there exists (finite)  $\lim_{t \rightarrow \infty} \int_0^t h(x(s))ds$  for  $x(t)$  of (1), then

$$\lim_{t \rightarrow \infty} h(x(t)) = 0$$

and consequently  $\lim_{t \rightarrow \infty} x(t) = \bar{x}$ . (6)

P r o o f. This assertion immediately follows from the well-known lemma of Barbălat [6].

Lemma 6. Under the assumptions of Lemma 3 every bounded solution  $x(t)$  of (1) either satisfies relation (3) or there exists such a root  $\bar{x}$  of  $h(x)$  that  $(x(t) - \bar{x})$  oscillates.

P r o o f - can be performed just in the same way as in [7].

3. Assuming all solutions of (1) being bounded, we now will deduce several important consequences of the above statements.

Consequence 1. If  $h(x)\text{sgn } x \geq 0$  is satisfied for all  $x$ , then every bounded solution  $x(t)$  of (1) either satisfies (6) or oscillates (i.e.  $\limsup_{t \rightarrow \infty} |x(t)| > 0 = \liminf_{t \rightarrow \infty} |x(t)|$ ) under (5), 6).

P r o o f. If  $x(t)$  is not oscillatory, then there is either  $x(t) \geq 0$  or  $x(t) \leq 0$  for  $t$  great, say  $t \geq T$  and

$$\int_T^t h(x(s))ds$$

is a monotone function. Thus there exists finite (cf. (5), 6))

$$\lim_{t \rightarrow \infty} \int_0^t h(x(s)) ds$$

and our assertion is implied by Lemma 5 immediately.

Consequence 2. Let the assumptions of Lemma 3 be fulfilled with conditions 5), 6) replaced by (2). If  $h(x) \operatorname{sgn} x \geq 0$  is satisfied for all  $x$  and

$$2'') \quad ag(0) - h'(0) > 0,$$

$$3'') \quad h'(0) > 0,$$

$$4') \quad g'(0) = 0,$$

then (3) is satisfied for every bounded solution  $x(t)$  of (1).

P r o o f. Consequence 1 says that every bounded solution  $x(t)$  of (1) either oscillates or satisfies (6). However, conditions 2''), 3''), 4') imply the existence of such an  $h$ -neighbourhood of the origin that assumptions of Lemma 2 are valid in it and therefore a trivial solution of autonomous equation (1) (i.e.  $p(t) \equiv 0$ ) is asymptotically stable. Hence, any oscillatory solution must be attracted to the origin with respect to Remark 1 and Lemma 3 and so such a possibility is reduced to (6) with  $\bar{x} = 0$  for  $p(t) \equiv 0$ .

Thus (3) is immediately implied by Lemma 3 and Lemma 4 and the same is true even for nonautonomous equation (1) in view of Lemma 1.

Consequence 3. Let  $h'(\bar{x}) \neq 0$  be satisfied for all zero points of  $h(x)$ . If

$$2') \quad ab - h'(x) \geq \delta > 0 \quad (\delta - \text{const.})$$

holds for all  $x$  and  $a^2 > g(x) \equiv b > 0$ , then every bounded solution  $x(t)$  of (1) obeys (3), provided (2) and (4).

P r o o f. Lemma 6 asserts that every bounded solution  $x(t)$  of (1) either satisfies (3) or there exists a root  $\bar{x}$  of  $h(x)$  such that  $x(t) - \bar{x}$  oscillates, provided  $p(t) \equiv 0$  and  $a^2 > g(x) \equiv b > 0$  (i.e. assumptions of Lemma 3). However, assuming  $h'(\bar{x}) \neq 0$  and  $ab - h'(x) \geq \delta > 0$ , the roots  $\bar{x}$  of  $h(x)$  with  $h'(\bar{x}) > 0$  are asymptotically stable and consequently any nontrivial  $x(t)$  of autonomous equation (1) is attracted to some  $\bar{x}$  with  $h'(\bar{x}) > 0$  (and therefore bounded as well) with respect to Remark 1. The remainder of the proof immediately follows from Lemma 1 and Lemma 4.

Remark 2. It is clear from the ideas introduced above that assumption  $h'(\bar{x}) \neq 0$  of Consequence 3 can be replaced by a weaker one, namely  $h(x)\text{sgn}(x - \bar{x}) \neq 0$ , in a suitable reduced neighbourhood of  $\bar{x}$ , but not  $h(x)\text{sgn } x < 0$ .

4. In the final section boundedness results will be given.

Theorem 1. Under the assumptions of Consequence 2 all solutions of (1) are bounded satisfying (3).

P r o o f. If any solution  $x(t)$  of (1) would not be bounded e.g.

$$\limsup_{t \rightarrow \infty} x(t) = \infty$$

(the case of  $\liminf_{t \rightarrow \infty} x(t) = -\infty$

can be treated quite analogically), then integrating (1) from a suitable  $T$  to  $t \geq T$  and using the above assumptions, we get the following inequality

$$\begin{aligned}
b(x(t) - x(T)) \operatorname{sgn} x &\leq \left| \int_T^t p(s) ds \right| - \int_T^t h(x(s)) \operatorname{sgn} x ds + a |x'(t) - \\
&- x'(T)| + |x''(t) - x''(T)| \leq \\
&\leq \left| \int_T^t p(s) ds \right| - \int_T^t |h(x(s))| ds + 2\max(a, 1)D',
\end{aligned}$$

$$\text{i.e. } |x(t)| \leq |x(T)| + \frac{1}{b}(2\max(a, 1)D' + P),$$

where  $P$  is a constant implied by (2), contradictionally. The remaining part of our assertion is included in Consequence 2.

Theorem 2. Let the assumptions of Consequence 3 be fulfilled with  $b < a^2/4$ . If conditions (2) and (4) yield such constants  $H, P, P_0$  that  $|p(t)| \leq P$ ,

$$\left| \int_0^t p(s) ds \right| \leq P_0$$

for  $t \geq 0$  and  $|h(x)| \leq H$  for all  $x \in (-\infty, \infty)$  together with

$$\min(d(\bar{x}_k, \bar{x}_{k+1}), d(\bar{x}_k, \bar{x}_{k-1})) > \frac{2(H+P)}{b} \left( \frac{2}{a} + \frac{a}{b} \right) + \frac{P_0}{b}, \quad (7)$$

where  $\bar{x}_k$  are the roots of  $h(x)$  with  $h'(\bar{x}_k) > 0$  and  $\bar{x}_{k-1}, \bar{x}_{k+1}$  denote the couple of adjacent zero points of  $\bar{x}_k$  ( $k = 0, \pm 2, \pm 4, \dots$ ), then all solutions of (1) are bounded satisfying (3).

P r o o f. The boundedness of all solutions of (1) can be verified quite analogously to [7]. Let us note that this follows directly from the assumptions of Consequence 3 for the autonomous equation even for  $g(x) \neq b$ . Indeed, if any its solution would not be bounded i.e.  $\limsup_{t \rightarrow \infty} x(t) = \infty$



or  $\liminf_{t \rightarrow \infty} x(t) = -\infty$ , then such a zero point of  $h(x)$  exists attracting  $x(t)$  asymptotically with respect to Remark 1 and Lemma 3, contradictionally. The remaining part of our assertion is included in Consequence 3.

Remark 3. Considering equation (1) with  $a^2/4 > g(x) \equiv b > 0$  and (4), it is clear that Theorem 0 of Swick [4] can be generalized in the following way: under 2), 3) condition 1) takes the local form  $ab - h'(0) > 0$  and under 1), 3) condition 2) can be replaced by much weaker assumption of oscillatory  $h(x)$  with (7), in general, but not  $h(x)\text{sgn } x < 0$  satisfied in reduced neighbourhoods of the zero points of  $h(x)$ .

Remark 4. Further generalization could be certainly done if either  $ag(x)$  is great enough or  $|g'(x)|$  is sufficiently small. This way is very important from the technical point of view, because of considering the phase-synchronization problem [8, 9].

#### REFERENCES

- [1] S w i c k, K.E.: Asymptotic behavior of the solutions of certain third order differential equations, SIAM J.Appl. Math. 19, 1, 1970, 96-102.
- [2] Y o s h i z a w a, T.: Stability theory by Liapunov's second method. Math.Soc.Japan, Tokyo 1966.
- [3] A n d r e s, J.: On stability and instability of the roots of the oscillatory function in a certain nonlinear differential equation of the third order. Čas.pěst.mat. 3, 1986, 225-229.
- [4] V o r á č e k, J.: Über eine nichtlineare Differentialgleichung dritter Ordnung. Czech.Math.J. 20, 95, 1970, 207-219.
- [5] C o p p e l, W.A.: Stability and asymptotic behavior of differential equations. D.C.Heath, Boston 1965.
- [6] B a r b a l a t, I.: Systèmes d'équations différentielles d'oscillations non linéaires. Rev.Math.Pures Appl. 4, 2, 1959, 267-270.

- [7] A n d r e s, J.: Boundedness of solutions of the third order differential equation with the oscillatory restoring and forcing terms. Czech.Math.J, 36, 1, 1986, 1-6.
- [8] B a k a e v, Yu.N.: Synchronization properties of the automatic control phase system of the third order (in Russian). Radiotekh.Elektron. 10, 6, 1965, 1083-1087.
- [9] A n d r e s, J. and Š t r u n c, M.: Lagrange-like stability of local cycles to a certain forced phase-locked loop described by the third-order differential equation. To appear in Rev.Roum.Sci.Techn. 32, 2, 1987, 219-223.

ASYMPTOTICKÉ VLASTNOSTI ŘEŠENÍ JISTÉ DIFERENCIÁLNÍ  
ROVNICE TŘETÍHO ŘÁDU S OSCILATORICKÝM OBNOVUJÍCÍM ČLEMEM

Souhrn

V práci je upřesněn a doplněn Swickův výsledek [1], týkající se vlastnosti (3), kterou nabývají všechna řešení rovnice (1). Za předpokladu o ohraničenosti funkce  $h(x)$  je ukázáno, jak může být podmínka 2) jeho věty O lokalizována do počátku a zejména, že funkce  $h(x)\operatorname{sgn} x$  může zabíhat i pod osu  $x$ .

АСИМПТОТИЧЕСКИЕ СВОЙСТВА РЕШЕНИЙ  
ОДНОГО ДИФФЕРЕНЦИАЛЬНОГО УРАВНЕНИЯ ТРЕТЬЕГО ПОРЯДКА  
С ОСЦИЛЛИРУЮЩИМ ВОССТАНАВЛИВАЮЩИМ ЧЛЕНОМ

Резюме

В работе уточняется и дополняется результат Свика [1], относящийся к свойству (3), которому подчиняются все решения уравнения (1). Ввиду предположения ограниченности функции  $h(x)$  показано, что условие 2) теоремы 0 можно локализовать в начало координат и доказано, что функция  $h(x)\operatorname{sgn} x$  может находиться тоже под осью  $x$ .

Author's address:

RNDr. Jan Andres, CSc.  
přirodovědecká fakulta  
Univerzity Palackého  
Gottwaldova 15  
771 46 Olomouc  
ČSSR /Czechoslovakia/