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A GRADIENT ESTIMATE FOR SOLUTIONS OF THE
HEAT EQUATION II

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Abstract. The author obtains an estimate for the spatial gradient of solutions of the heat equation, subject to a homogeneous Neumann boundary condition, in terms of the gradient of the initial data. The proof is accomplished via the maximum principle; the main assumption is that the sufficiently smooth boundary be convex.

Keywords: gradient estimate, heat equation, maximum principle

MSC 2000: 35K05

1. INTRODUCTION

In [1] the writer obtained an estimate for the spatial gradient of the solution $u(x, t)$ of the following initial-boundary value problem for the heat equation:

$$(1.1) \quad \begin{cases} u_t = \Delta u & \text{in } \Omega \times (0, \infty) \\ u = 0 & \text{on } \partial\Omega \times (0, \infty) \\ u(x, 0) = f(x) & \text{in } \Omega, \end{cases}$$

where Ω is a bounded domain in R^n , $n \geq 2$. Assuming that $f(x) \in C^1(\bar{\Omega})$ and vanished on $\partial\Omega$; and that $\partial\Omega$ was C^3 and satisfied an appropriate mean curvature condition (see (1.6) in [1]), the estimate

$$(1.2) \quad |\text{grad } u(x, t)| \leq \max_{\bar{\Omega}} |\text{grad } f(x)|, \quad (x, t) \in \partial\Omega \times (0, \infty)$$

was obtained as a consequence of the maximum principle. (Here $\text{grad } u(x, t)$ denotes the gradient with respect to the spatial variables x).

The purpose of this paper is to obtain the same estimate for solutions of the problem (1.1) in which u satisfies a homogeneous Neumann boundary condition rather than a homogeneous Dirichlet boundary condition.

In order to obtain this result we need a stronger assumption on $\partial\Omega$ than the mean curvature assumption (1.6) made in [1]. In fact we need to assume that $\partial\Omega$ satisfies a convexity condition.

To describe this condition let p be a typical point on $\partial\Omega$ and suppose that after suitable rotation and translation of our coordinate system placing p at the origin of the system, the portion of $\partial\Omega$ lying in a neighbourhood of p is the surface corresponding to the function

$$(1.3) \quad x_n = g(x_1, \dots, x_{n-1})$$

where (x_1, \dots, x_{n-1}) varies over a neighbourhood of $(x_1 = 0, \dots, x_{n-1} = 0)$, with $g(0, \dots, 0) = 0$ and with the positive x_n direction corresponding to the outward normal direction from $\partial\Omega$ at p . Then the convexity condition that we shall assume $\partial\Omega$ to satisfy is that

$$(1.4) \quad \sum_{1 \leq j, k \leq n-1} g_{x_j x_k}(0, \dots, 0) \eta_j \eta_k \leq 0$$

for any $\eta = (\eta_1, \dots, \eta_{n-1}) \in R^{n-1}$.

We can now state the result we wish to prove as follows:

Theorem 1. *Assume*

$$\begin{cases} u_t = \Delta u & \text{in } \Omega \times (0, \infty) \\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial\Omega \times (0, \infty) \\ u(x, 0) = f(x) & \text{in } \Omega \end{cases}$$

with $f(x) \in C^1(\overline{\Omega})$ and satisfying the boundary condition

$$\frac{\partial f}{\partial n} = 0 \quad \text{on } \partial\Omega.$$

Suppose further that $\partial\Omega \in C^3$ and satisfies the convexity condition (1.4). Then

$$(1.5) \quad |\text{grad } u(x, t)| \leq \max_{\overline{\Omega}} |\text{grad } f(x)|, \quad (x, t) \in \Omega \times (0, \infty).$$

The proof of the theorem will be presented in the following section of the paper.

SECTION 2

The proof of Theorem 1 will be conducted along the same general lines as the proof of the same estimate (1.2) for problem (1.1) given in [1]. As in that proof it suffices, in view of the maximum principle (see Proposition 2.1 and Theorem 2.2 of [1]), to show that

$$(2.1) \quad \frac{\partial}{\partial n} |\text{grad } u|^2 \Big|_{\partial\Omega \times (0, \infty)} \leq 0.$$

However, unlike that proof, where to establish (2.1) we used the fact that u was a solution of the heat equation in $\Omega \times (0, \infty)$, we don't use the equation here. Rather, the conclusion (2.1) stems in the present case from the boundary condition $\frac{\partial u}{\partial n} = 0$ satisfied by u on $\partial\Omega \times (0, \infty)$ and the convexity condition (1.4) satisfied by $\partial\Omega$. This result is of independent interest and we state it separately as:

Theorem 2. *Suppose that $u(x)$ is a $C^2(\overline{\Omega})$ function which satisfies the boundary condition*

$$(2.2) \quad \frac{\partial u}{\partial n} \Big|_{\partial\Omega} = 0;$$

and suppose that $\partial\Omega$ is C^3 and satisfies the convexity condition (1.4). Then $|\text{grad } u(x)|^2$ satisfies the boundary condition

$$(2.3) \quad \frac{\partial}{\partial n} |\text{grad } u(x)|^2 \Big|_{\partial\Omega} \leq 0.$$

Preliminaries. To prove Theorem 2 we are going to show that for a typical point p of $\partial\Omega$

$$(2.4) \quad \frac{\partial}{\partial n} |\text{grad } u(x)|^2 \Big|_p \leq 0.$$

For this purpose we introduce the same coordinate change used in [1] and delineated in Section 3 of that paper.

Recapitulating, that coordinate change was based on the function

$$x_n = g(x_1, \dots, x_{n-1})$$

which described the surface constituting that portion of $\partial\Omega$ lying in a sufficiently small neighbourhood of the point p , with p placed at the origin of our coordinate system, and so

$$(2.5) \quad g(0, \dots, 0) = 0.$$

We also assumed the positive x_n direction to correspond to the outward normal direction on $\partial\Omega$ at p , which implies that $x_n = 0$ is the tangent plane to $\partial\Omega$ at p ; so that necessarily

$$(2.6) \quad g_{x_j}(0, \dots, 0) = 0 \quad \text{for } j = 1, \dots, n-1.$$

Starting from the point $(\xi_1, \dots, \xi_{n-1}, g(\xi_1, \dots, \xi_{n-1}))$ on the surface describing $\partial\Omega$, we then proceeded ξ_n units in the outward normal direction arriving at the point (x_1, \dots, x_n) in R^n . Accordingly, the coordinates of the resulting point $x = (x_1, \dots, x_n)$ are connected to the coordinates of $\xi = (\xi_1, \dots, \xi_n)$ through the formulas

$$(2.7) \quad \begin{cases} x_j = \xi_j - g_{\xi_j}(\xi_1, \dots, \xi_{n-1}) \left(1 + \sum_{k=1}^{n-1} g_{\xi_k}^2(\xi_1, \dots, \xi_{n-1})\right)^{-\frac{1}{2}} \xi_n \\ j = 1, \dots, n-1, \text{ and} \\ x_n = g(\xi_1, \dots, \xi_{n-1}) + \left(1 + \sum_{k=1}^{n-1} g_{\xi_k}^2(\xi_1, \dots, \xi_{n-1})\right)^{-\frac{1}{2}} \xi_n. \end{cases}$$

And it is these equations, abbreviated as $x = x(\xi)$, which describe the coordinate change from ξ to x that we are going to use prove (2.4).

Clearly, from the way we arrived at (2.7), the outward normal derivative in the x coordinates on $\partial\Omega$ corresponds to differentiation with respect to ξ_n in the ξ coordinates when $\xi_n = 0$. More precisely if $\varphi(x)$ represents a function in the x coordinates and $\psi(\xi)$ represents the corresponding function in the ξ coordinates, i.e. $\psi(\xi) = \varphi(x(\xi))$, then

$$(2.8) \quad \frac{\partial\varphi(x)}{\partial n} \Big|_{\partial\Omega} = \frac{\partial\psi(\xi)}{\partial\xi_n} \Big|_{\xi_n=0};$$

in particular

$$(2.9) \quad \frac{\partial\varphi(x)}{\partial n} \Big|_p = \frac{\partial\psi(\xi)}{\partial\xi_n} \Big|_{\xi=0}.$$

The differentiability properties of the transformation $x = x(\xi)$ defined by (2.7) are described in Propositions 3.1 and 3.2 of [1] and we summarize them here.

Most importantly, if $g(\xi_1, \dots, \xi_{n-1})$ is C^2 in a neighbourhood of $(\xi_1 = 0, \dots, \xi_{n-1} = 0)$, then $x = x(\xi)$ is a C^1 transformation in a neighbourhood of $\xi = 0$, sending $\xi = 0$ into $x = 0$, whose Jacobian at the origin is the identity matrix:

$$(2.10) \quad \frac{\partial x}{\partial \xi} \Big|_{\xi=0} = I.$$

Consequently, the inverse transformation $\xi = \xi(x)$ exists in a neighbourhood of $x = 0$, is C^1 there and its Jacobian at the origin is also the identity matrix:

$$(2.11) \quad \left. \frac{\partial \xi}{\partial x} \right|_{x=0} = I.$$

Moreover, if $g(\xi_1, \dots, \xi_{n-1})$ is C^3 in a neighbourhood of $(\xi_1 = 0, \dots, \xi_{n-1} = 0)$, then both $x = x(\xi)$ and $\xi = \xi(x)$ are C^2 transformations in neighbourhoods of $\xi = 0$ and $x = 0$, respectively; with the following identities holding for their second derivatives at the origin $\xi = x = 0$;

$$(2.12) \quad \left. \frac{\partial}{\partial \xi_m} \left(\frac{\partial \xi_j}{\partial x_l} \right) \right|_{\xi=0} = - \left. \frac{\partial}{\partial \xi_m} \left(\frac{\partial x_j}{\partial \xi_l} \right) \right|_{\xi=0},$$

$j, l, m = 1, \dots, n$ (see equation (3.9) of [1]).

P r o o f of Theorem 2. We are now prepared to establish Theorem 2 by showing that the function $u(x)$ which that theorem concerns satisfies the condition (2.4). Our first step in doing so is to introduce the coordinate transformation $x = x(\xi)$ defined by (2.7) and then to consider the function $u(x)$ referred to ξ coordinates which we denote by $v(\xi)$, i.e. $v(\xi) = u(x(\xi))$. Expressing $|\text{grad } u(x)|^2$ in terms of $v(\xi)$ we obtain

$$|\text{grad } u(x)|^2 = \sum_{1 \leq j, k \leq n} b_{jk} \frac{\partial v}{\partial \xi_j} \frac{\partial v}{\partial \xi_k}$$

where

$$(2.13) \quad b_{jk} = \sum_{i=1}^n \frac{\partial \xi_j}{\partial x_i} \frac{\partial \xi_k}{\partial x_i} \quad j, k = 1, \dots, n.$$

Hence, in view of the correspondence (2.8) between differentiation in the normal direction on $\partial\Omega$ in the x coordinates and differentiation with respect to ξ_n when $\xi_n = 0$ in the ξ coordinates, we have

$$(2.14) \quad \begin{aligned} \left. \frac{\partial}{\partial n} |\text{grad } u|^2 \right|_{\partial\Omega} &= \left. \frac{\partial}{\partial \xi_n} \left(\sum_{1 \leq j, k \leq n} b_{jk} \frac{\partial v}{\partial \xi_j} \frac{\partial v}{\partial \xi_k} \right) \right|_{\xi_n=0} \\ &= \sum_{1 \leq j, k \leq n} \left. \frac{\partial}{\partial \xi_n} (b_{jk}) \frac{\partial v}{\partial \xi_j} \frac{\partial v}{\partial \xi_k} \right|_{\xi_n=0} + \sum_{1 \leq j, k \leq n} 2b_{jk} \left. \frac{\partial^2 v}{\partial \xi_n \partial \xi_j} \frac{\partial v}{\partial \xi_k} \right|_{\xi_n=0}. \end{aligned}$$

But now in terms of $v(\xi)$, our hypotheses $\left. \frac{\partial u}{\partial n} \right|_{\partial\Omega} = 0$, asserts, again because of (2.8), that $\left. \frac{\partial v}{\partial \xi_n} \right|_{\xi_n=0} = 0$; and consequently $\left. \frac{\partial^2 v}{\partial \xi_n \partial \xi_j} \right|_{\xi_n=0} = 0$ for $j \neq n$; thus the preceding becomes

$$(2.15) \quad \left. \frac{\partial}{\partial n} |\text{grad } u|^2 \right|_{\partial\Omega} = \sum_{1 \leq j, k \leq n-1} \left. \frac{\partial}{\partial \xi_n} (b_{jk}) \frac{\partial v}{\partial \xi_j} \frac{\partial v}{\partial \xi_k} \right|_{\xi_n=0} + \sum_{k=1}^{n-1} 2b_{nk} \left. \frac{\partial^2 v}{\partial \xi_n^2} \frac{\partial v}{\partial \xi_k} \right|_{\xi_n=0}.$$

Specializing down to the point $x = p$ on $\partial\Omega$, which corresponds to $\xi = 0$, we then find, on account of $b_{nk}|_{\xi=0} = 0$ for $k \neq n$ (see equation (4.6) of [1]), that

$$(2.16) \quad \frac{\partial}{\partial n} |\text{grad } u|^2 \Big|_p = \sum_{1 \leq j, k \leq n-1} \frac{\partial}{\partial \xi_n} (b_{jk}) \frac{\partial v}{\partial \xi_j} \frac{\partial v}{\partial \xi_k} \Big|_{\xi=0}.$$

Finally, from the evaluation

$$(2.17) \quad \frac{\partial}{\partial \xi_n} (b_{jk}) \Big|_{\xi=0} = 2g_{\xi_j \xi_k}(0, \dots, 0), \quad 1 \leq j, k \leq n-1,$$

which we will establish in a moment, (2.16) then yields

$$\frac{\partial}{\partial n} |\text{grad } u|^2 \Big|_p = \sum_{1 \leq j, k \leq n-1} 2g_{\xi_j \xi_k}(0, \dots, 0) \frac{\partial v}{\partial \xi_j} \frac{\partial v}{\partial \xi_k} \Big|_{\xi=0} \leq 0$$

because of the assumed convexity condition (1.4) regarding $\partial\Omega$. This proves (2.4) and with it Theorem 2.

It remains to establish the evaluation (2.17). For this purpose we differentiate the defining formula (2.13) for b_{jk} with respect to ξ_n and evaluate at $\xi = 0$:

$$\frac{\partial}{\partial \xi_n} (b_{jk}) \Big|_{\xi=0} = \sum_{i=1}^n \frac{\partial}{\partial \xi_n} \left(\frac{\partial \xi_j}{\partial x_i} \right) \frac{\partial \xi_k}{\partial x_i} \Big|_{\xi=0} + \sum_{i=1}^n \frac{\partial \xi_j}{\partial x_i} \frac{\partial}{\partial \xi_n} \left(\frac{\partial \xi_k}{\partial x_i} \right) \Big|_{\xi=0}.$$

In view of (2.11), $\frac{\partial \xi_k}{\partial x_i} \Big|_{\xi=0} = \delta_{ki}$, where δ_{ki} is the Kronecker delta, *i.e.* $\delta_{ki} = 1$ if $k = i$ and is zero otherwise. Hence

$$\frac{\partial}{\partial \xi_n} (b_{jk}) \Big|_{\xi=0} = \frac{\partial}{\partial \xi_n} \left(\frac{\partial \xi_j}{\partial x_k} \right) \Big|_{\xi=0} + \frac{\partial}{\partial \xi_n} \left(\frac{\partial \xi_k}{\partial x_j} \right) \Big|_{\xi=0}.$$

Making use of (2.12) this becomes

$$\frac{\partial}{\partial \xi_n} (b_{jk}) \Big|_{\xi=0} = -\frac{\partial}{\partial \xi_n} \left(\frac{\partial x_j}{\partial \xi_k} \right) \Big|_{\xi=0} - \frac{\partial}{\partial \xi_n} \left(\frac{\partial x_k}{\partial \xi_j} \right) \Big|_{\xi=0}.$$

The derivatives on the right are then evaluated directly by differentiating the expressions (2.7) defining the x_j 's in terms of the ξ_k 's; taking (2.6) into account, this yields

$$\frac{\partial}{\partial \xi_n} \left(\frac{\partial x_j}{\partial \xi_k} \right) \Big|_{\xi=0} = \frac{\partial}{\partial \xi_k} \left(\frac{\partial x_j}{\partial \xi_n} \right) \Big|_{\xi=0} = -g_{\xi_j \xi_k}(0, \dots, 0) \quad \text{for } j, k = 1, \dots, n-1;$$

and (2.17) follows. □

References

- [1] *C. Kahane*: A gradient estimate for solutions of the heat equation. Czechoslovak Math. J. 48(123) (1998), 711–725.

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