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Mathematica Slovaca, Vol. 49 (1999), No. 1, 41--52

Persistent URL: <http://dml.cz/dmlcz/132859>

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PERIODIC SOLUTIONS IN SYSTEMS AT RESONANCES WITH SMALL RELAY HYSTERESIS

MICHAL FEČKAN

(Communicated by Milan Medved')

ABSTRACT. We study the existence of periodic solutions for certain systems of constant ordinary differential equations at resonances with relay hysteresis.

1. Introduction

In this paper, we deal with relay hysteresis [5]. So there is given a pair of real numbers $\alpha < \beta$ (thresholds) and a pair of real-valued continuous functions $h_o \in C([\alpha, \infty), \mathbb{R})$, $h_c \in C((-\infty, \beta], \mathbb{R})$ such that $h_o(u) \geq h_c(u) \forall u \in [\alpha, \beta]$. Moreover, we suppose that h_o, h_c are bounded on $[\alpha, \infty)$, $(-\infty, \beta]$, respectively.

For a given continuous input $u(t)$, $t \geq t_0$, one defines the output $v(t) = f(u)(t)$ of the relay hysteresis operator as follows

$$f(u)(t) = \begin{cases} h_o(u(t)) & \text{if } u(t) \geq \beta, \\ h_c(u(t)) & \text{if } u(t) \leq \alpha, \\ h_o(u(t)) & \text{if } u(t) \in (\alpha, \beta) \text{ and } u(\tau(t)) = \beta, \\ h_c(u(t)) & \text{if } u(t) \in (\alpha, \beta) \text{ and } u(\tau(t)) = \alpha, \end{cases}$$

where $\tau(t) = \sup\{s : s \in [t_0, t], u(s) = \alpha \text{ or } u(s) = \beta\}$. If $\tau(t)$ does not exist (i.e. $u(\sigma) \in (\alpha, \beta)$ for $\sigma \in [t_0, t]$), then $f(u)(\sigma)$ is undefined and we have initially to set the relay open or closed when $u(t_0) \in (\alpha, \beta)$. Of course, when either $h_o(\beta) > h_c(\beta)$ or $h_o(\alpha) > h_c(\alpha)$ then $f(u)$ is generally discontinuous.

Electrical engineers are interested in the periodic behaviour of circuits with hysteresis. A circuit with a relay hysteresis could be modelled by

$$L_m y = f(y),$$

AMS Subject Classification (1991): Primary 34A60, 34C25, 94C05.

Key words: periodic solution, relay hysteresis.

Supported by Grant GA-MS V2M12-G

where L_m is an m th-order differential operator.

In this paper, in order to deal with much more general equations, we are interested in the periodic oscillations of systems given by

$$\dot{x} = Ax + \mu f(x_1)b, \quad (1.1)$$

where A is a constant $n \times n$ matrix, x_1 is the first component of $x \in \mathbb{R}^n$, $b \in \mathbb{R}^n$ is a constant vector and $\mu \in \mathbb{R}$ is a small parameter. Similar systems are studied in [1] and [5]–[7].

In contrast to these papers, we assume that (1.1) is at resonance, i.e. $\dot{x} = Ax$ has a nonzero periodic solution. The aim of this paper is to find conditions ensuring the existence of periodic oscillations of (1.1) for $\mu \neq 0$ small. Since (1.1) is generally discontinuous, we consider this as a differential inclusion. The method used in this paper is a combination of [3] and [6], i.e. we apply to (1.1) a Lyapunov-Schmidt decomposition procedure together with topological degree theory for multivalued mappings [2]. Periodically forced problems of (1.1) are also investigated. We end the paper with examples of unforced and forced third-order ordinary differential equations with a small relay hysteresis.

2. The existence of periodic solutions

We suppose that the following condition holds

- i) $W = \{x \in \mathbb{R}^n : x = e^A x\} \neq \{0\}$ and there is an $x_0 \in W$ such that $Ax_0 \neq 0$.

By [4] we have

$$W^* = \{x \in \mathbb{R}^n : x = e^{-A^*} x\} \neq \{0\}, \quad \dim W^* = \dim W = d > 1.$$

Moreover, the linear equation

$$\dot{x} = Ax + h(t), \quad h \in L_2 = L_2([0, 1], \mathbb{R}^n)$$

has a solution $x \in W^{1,\infty} = W^{1,\infty}([0, 1], \mathbb{R}^n)$ satisfying $x(0) = x(1)$ if and only if

$$\forall w \in W^* \quad \int_0^1 \langle h(s), e^{-A^* s} w \rangle ds = 0.$$

Here $\langle \cdot, \cdot \rangle$ is the inner product on \mathbb{R}^n . The norm on $W^{1,\infty}$ is denoted by $\|\cdot\|$.

Let $x = \mathcal{K}h$ be the unique such solution satisfying

$$\forall z \in W \quad \int_0^1 \langle x(s), e^{As} z \rangle ds = 0.$$

We put

$$X = \left\{ x \in W^{1,\infty} : \int_0^1 \langle x(s), e^{As} z \rangle ds = 0 \quad \forall z \in W \right\}.$$

Let

$$\Pi: L_2 \rightarrow \left\{ h \in L_2 : \int_0^1 \langle h(s), e^{-A^*s} w \rangle ds = 0 \quad \forall w \in W^* \right\}$$

be the orthogonal projection. Of course, $\mathcal{K}: \text{im } \Pi \rightarrow X$ is linear and bounded.

By taking a basis $\{w_1, \dots, w_d\}$ of W , we put $\gamma_i(t) = e^{At} w_i$, $i = 1, \dots, d$.

Let

$$\gamma(\theta, t) = \sum_{i=1}^{d-1} \theta_i \gamma_i(t), \quad \theta_i \in \mathbb{R}.$$

By i), we obtain that $d \geq 2$ and $\{w_1, \dots, w_d\}$ can be chosen such that any solution of $\dot{x} = Ax$, $x(0) \in W$ has the form $\gamma(\theta, t + \omega)$, $\omega \in \mathbb{R}$, $\theta \in \mathbb{R}^{d-1}$. From now on, $\{w_1, \dots, w_d\}$ will be such a basis. Let $\gamma_1(\theta, t)$ be the first component of $\gamma(\theta, t)$. We need the following conditions to hold:

ii) There is an open bounded subset $\emptyset \neq \mathcal{O} \subset \mathbb{R}^{d-1}$ such that $\forall \theta \in \mathcal{O}$ and $\forall t_0 \in \mathbb{R}$

$$\gamma_1(\theta, t_0) = \alpha, \beta \implies \dot{\gamma}_1(\theta, t_0) \neq 0.$$

iii) $\forall \theta \in \mathcal{O} \quad \min_{t \in \mathbb{R}} \gamma_1(\theta, t) < \alpha, \max_{t \in \mathbb{R}} \gamma_1(\theta, t) > \beta$.

Now in (1.1) we make the following change of variables

$$x((1 + \mu\omega)t) = \mu z(t) + \gamma(\theta, t), \quad \omega \in \mathbb{R}.$$

The conditions ii) and iii) imply that if $z \in X$ satisfies $\|z\| \leq K$ and μ is sufficiently small, then $\mu z_1(t) + \gamma_1(\theta, t)$ crosses α and β strictly monotonically for arbitrary $\theta \in \mathcal{O}$.

We rewrite (1.1) as a differential inclusion of the form

$$\dot{x} - Ax \in \mu F(x_1) b, \tag{2.1}$$

where F is a multivalued mapping defined as follows

$$F(u)(t) = \begin{cases} f(u)(t) & \text{if } u(t) \neq \alpha, \beta, \\ h_c(\alpha) & \text{if } u(t) = \alpha, u(\tau(s)) = \alpha, \text{ for any } s < t \text{ near } t, \\ h_o(\beta) & \text{if } u(t) = \beta, u(\tau(s)) = \beta, \text{ for any } s < t \text{ near } t, \\ [h_c(\alpha), h_o(\alpha)] & \text{if } u(t) = \alpha, u(\tau(s)) = \beta, \text{ for any } s < t \text{ near } t, \\ [h_c(\beta), h_o(\beta)] & \text{if } u(t) = \beta, u(\tau(s)) = \alpha, \text{ for any } s < t \text{ near } t. \end{cases}$$

ii) and iii) imply that if $u(t) = \mu z_1(t) + \gamma_1(\theta, t)$ with $z \in X$ bounded and μ sufficiently small, then $F(u)$ is well-defined. By a solution of a differential

inclusion in this paper we mean a function which is absolute continuous and which satisfies that differential inclusion almost everywhere.

Hence (2.1) has the form

$$\dot{z}(t) - Az(t) \in (1 + \mu\omega)F(\mu z_1 + \gamma_1(\theta, \cdot))(t)b + \omega A(\mu z(t) + \gamma(\theta, t)). \quad (2.2)$$

By taking the mapping

$$\begin{aligned} G(z, \omega, \theta, \mu, \lambda) = \\ = \{h \in L_2 : \text{satisfying the relation} \\ h(t) \in (1 + \lambda\mu\omega)F(\lambda\mu z_1 + \gamma_1(\theta, \cdot))(t)b + \omega A(\lambda\mu z(t) + \gamma(\theta, t)) \\ \text{a.e. on } [0, 1]\}, \end{aligned}$$

(2.2) has the form

$$\dot{z} - Az \in G(z, \omega, \theta, \mu, 1). \quad (2.3)$$

Using Π and \mathcal{K} , we rewrite (2.3) as follows

$$\begin{cases} 0 \in H(z, \omega, \theta, \mu, 1) \\ H(z, \omega, \theta, \mu, \lambda) = \{(z - \lambda\mathcal{K}\Pi h, \mathcal{L}h) : h \in G(z, \omega, \theta, \mu, \lambda)\}, \end{cases} \quad (2.4)$$

where $\mathcal{L}: L_2 \rightarrow \mathbb{R}^d$ is defined by

$$\mathcal{L}h = \left(\int_0^1 \langle h(s), e^{-A^*s} \tilde{w}_1 \rangle ds, \dots, \int_0^1 \langle h(s), e^{-A^*s} \tilde{w}_d \rangle ds \right)$$

for a basis $\{\tilde{w}_1, \dots, \tilde{w}_d\}$ of W^* .

Since f is bounded in (1.1), for arbitrary $\Gamma > 0$ there exist $\mu_0 > 0$ and $K > 0$ such that

$$\begin{aligned} \|\mathcal{K}\Pi h\| \leq K \quad \text{for arbitrary } h \in G(z, \omega, \theta, \mu, \lambda), \\ \|z\| \leq K + 1, \quad |\omega| \leq \Gamma, \quad \theta \in \mathcal{O}, \quad |\mu| \leq \mu_0, \quad \lambda \in [0, 1]. \end{aligned}$$

Moreover, if μ_0 is sufficiently small then by ii) and iii), the mapping

$$H: \Omega \times [-\mu_0, \mu_0] \times [0, 1] \rightarrow 2^{X \times \mathbb{R}^d} \quad (2.5)$$

is well-defined and singlevalued, where

$$\Omega = \{(z, \omega, \theta) \in X \times \mathbb{R}^d : \|z\| < K + 1, (\omega, \theta) \in \mathcal{B}\}$$

and \mathcal{B} is an open bounded non-empty subset satisfying $\bar{\mathcal{B}} \subset \mathbb{R} \times \mathcal{O}$.

The arguments of [6; pp. 677-678] imply that $H: \Omega \times [-\mu_0, \mu_0] \times [0, 1] \rightarrow X \times \mathbb{R}^d$ is continuous and also compact. Similarly, the mapping given by

$$\begin{aligned} M: \mathbb{R} \times \mathcal{O} \rightarrow \mathbb{R}^d, \quad M(\omega, \theta) = \mathcal{L}h \\ h(t) = F(\gamma_1(\theta, \cdot))(t)b + \omega A\gamma(\theta, t) \quad \text{a.e. on } [0, 1] \end{aligned} \quad (2.6)$$

is continuous.

THEOREM 2.1. *Assume that i)–iii) hold. If there is a non-empty open bounded set \mathcal{B} such that $\overline{\mathcal{B}} \subset \mathbb{R} \times \mathcal{O}$ and*

- (i) $0 \notin M(\partial\mathcal{B})$,
- (ii) $\deg(M, \mathcal{B}, 0) \neq 0$,

where \deg is the Brouwer degree and M is given by (2.6), then there are constants $K_1 > 0$ and $\mu_0 > 0$ such that for arbitrary $|\mu| < \mu_0$, there exist $(\omega_\mu, \theta_\mu) \in \mathcal{B}$ and a $(1 + \mu\omega_\mu)$ -periodic solution x_μ of (1.1) satisfying

$$\sup_{t \in \mathbb{R}} |x_\mu(t) - \gamma(\theta_\mu, t/(1 + \mu\omega_\mu))| \leq K_1 |\mu|.$$

P r o o f . First we show

$$0 \notin H(\partial\Omega \times [-\mu_0, \mu_0] \times [0, 1])$$

for arbitrary $\mu_0 > 0$ sufficiently small. Assume the contrary. Then there exist

$$\begin{aligned} [0, 1] \ni \lambda_i \rightarrow \lambda_0, \quad \|z_i\| \leq K + 1, \quad \mu_i \rightarrow 0, \quad i \in \mathbb{N} \\ \partial\mathcal{B} \ni (\omega_i, \theta_i) \rightarrow (\omega_0, \theta_0) \in \partial\mathcal{B}, \quad h_i \in G(z_i, \omega_i, \theta_i, \mu_i, \lambda_i) \end{aligned}$$

such that

$$\mathcal{L}h_i = 0.$$

We can assume that $z_i \rightarrow z$ in $C([0, 1], \mathbb{R}^n)$ and h_i tends weakly to some $h_0 \in L^2$. Then by applying the standard arguments (see the proof of [2; Remarks 5.5.1]), we obtain

$$h \in G(z, \omega_0, \theta_0, 0, \lambda_0) \quad \text{and} \quad \mathcal{L}h_0 = 0,$$

i.e. $0 = M(\omega_0, \theta_0)$ for some $(\omega_0, \theta_0) \in \partial\mathcal{B}$. This contradicts (i) of this theorem.

Consequently, we compute for μ sufficiently small

$$\begin{aligned} \deg(H(\cdot, \cdot, \cdot, \mu, 1), \Omega, 0) &= \deg(H(\cdot, \cdot, \cdot, \mu, 0), \Omega, 0) \\ &= \deg(M, \mathcal{B}, 0) \neq 0. \end{aligned}$$

Thus, (2.4) has a solution $(z, \omega, \theta) \in \Omega$ for arbitrary sufficiently small μ . The proof is finished. \square

Now we return to the differential equation

$$\begin{aligned} L_m y = \sum_{i=0}^m a_i y^{(i)} = \mu f(y), \\ a_i \in \mathbb{R}, \quad a_m = 1, \quad y^{(i)} = \frac{d^i}{dt^i} y. \end{aligned} \tag{2.7}$$

Of course, (2.7) can be rewritten in the form of (1.1). We put

$$L_m^* y = \sum_{i=0}^m (-1)^i a_i y^{(i)}.$$

Let ϕ_1, \dots, ϕ_d , respectively ψ_1, \dots, ψ_d , be a basis of the space of all 1-periodic solutions of $L_m y = 0$, respectively $L_m^* y = 0$. We suppose that i)–iii) hold for (2.7) and also ϕ_d is non-constant. A tedious computation shows that the mapping (2.6) for (2.7) of the form $M: \mathbb{R} \times \mathcal{O} \rightarrow \mathbb{R}^d$ is given by

$$M(\omega, \theta) = \left(\int_0^1 h(s) \psi_1(s) ds, \dots, \int_0^1 h(s) \psi_d(s) ds \right), \quad (2.8)$$

$$h(t) = F(\eta(\theta, \cdot))(t) + \omega \sum_{i=1}^m i a_i \eta^{(i)}(\theta, t) \quad \text{a.e. on } [0, 1],$$

where $\eta(\theta, t) = \sum_{i=1}^{d-1} \theta_i \phi_i(t)$. Theorem 2.1 implies the following result.

THEOREM 2.2. *Assume that ϕ_d is non-constant and that the following conditions hold:*

- a) *There is an open bounded subset $\emptyset \neq \mathcal{O} \subset \mathbb{R}^{d-1}$ such that $\forall \theta \in \mathcal{O}$ and $\forall t_0 \in \mathbb{R}$*

$$\eta(\theta, t_0) = \alpha, \beta \implies \dot{\eta}(\theta, t_0) \neq 0.$$

- b) $\forall \theta \in \mathcal{O} \quad \min_{t \in \mathbb{R}} \eta(\theta, t) < \alpha, \max_{t \in \mathbb{R}} \eta(\theta, t) > \beta.$

If there is a non-empty open bounded set \mathcal{B} such that $\overline{\mathcal{B}} \subset \mathbb{R} \times \mathcal{O}$ and

- (i) $0 \notin M(\partial \mathcal{B})$,
(ii) $\deg(M, \mathcal{B}, 0) \neq 0$,

where M is given by (2.8), then there exist constants $K_1 > 0$ and $\mu_0 > 0$ such that for arbitrary $|\mu| < \mu_0$, there exist $(\omega_\mu, \theta_\mu) \in \mathcal{B}$ and an $(1 + \mu\omega_\mu)$ -periodic solution y_μ of (2.7) satisfying

$$\sup_{t \in \mathbb{R}} |y_\mu(t) - \eta(\theta_\mu, t/(1 + \mu\omega_\mu))| \leq K_1 |\mu|.$$

The results of [3] can be modified to give existence results of subharmonic solutions of nonautonomous periodic versions of (1.1) expressed in the following theorems.

THEOREM 2.3. *Consider*

$$\dot{x} = Ax + \mu(f(x_1)b + q(t)), \quad (2.9)$$

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where $q \in C(\mathbb{R}, \mathbb{R}^n)$ is 1-periodic and A, f, b are given in (1.1). Assume that i)–iii) hold. If there is a non-empty open bounded set \mathcal{B} such that $\bar{\mathcal{B}} \subset \mathbb{R} \times \mathcal{O}$ and

- (i) $0 \notin M(\partial\mathcal{B}),$
- (ii) $\deg(M, \mathcal{B}, 0) \neq 0,$

where M is given by

$$\begin{aligned} M: \mathbb{R} \times \mathcal{O} &\rightarrow \mathbb{R}^d, & M(\omega, \theta) &= \mathcal{L}h, \\ h(t) &= F(\gamma_1(\theta, \cdot))(t)b + q(t + \omega) && \text{a.e. on } [0, 1], \end{aligned} \quad (2.10)$$

then there exist constants $K_1 > 0$ and $\mu_0 > 0$ such that for arbitrary $|\mu| < \mu_0$, there are $(\omega_\mu, \theta_\mu) \in \mathcal{B}$ and a 1-periodic solution x_μ of (2.9) satisfying

$$\sup_{t \in \mathbb{R}} |x_\mu(t) - \gamma(\theta_\mu, t - \omega_\mu)| \leq K_1 |\mu|.$$

THEOREM 2.4. Consider

$$L_m \dot{y} = \mu(f(y) + q(t)), \quad (2.11)$$

where L_m, f are given in (2.7) and $q \in C(\mathbb{R}, \mathbb{R})$ is 1-periodic. Assume that ϕ_d is non-constant, and a) and b) of Theorem 2.2 hold. If there is a non-empty open bounded set \mathcal{B} such that $\bar{\mathcal{B}} \subset \mathbb{R} \times \mathcal{O}$ and

- (i) $0 \notin M(\partial\mathcal{B}),$
- (ii) $\deg(M, \mathcal{B}, 0) \neq 0,$

where $M: \mathbb{R} \times \mathcal{O} \rightarrow \mathbb{R}^d$ is given by

$$\begin{aligned} M(\omega, \theta) &= \left(\int_0^1 h(s)\psi_1(s) \, ds, \dots, \int_0^1 h(s)\psi_d(s) \, ds \right), \\ h(t) &= F(\eta(\theta, \cdot))(t) + q(t + \omega) && \text{a.e. on } [0, 1], \end{aligned} \quad (2.12)$$

then there exist constants $K_1 > 0$ and $\mu_0 > 0$ such that for arbitrary $|\mu| < \mu_0$, there are $(\omega_\mu, \theta_\mu) \in \mathcal{B}$ and a 1-periodic solution y_μ of (2.11) satisfying

$$\sup_{t \in \mathbb{R}} |y_\mu(t) - \eta(\theta_\mu, t - \omega_\mu)| \leq K_1 |\mu|.$$

Remark 2.5. The boundedness of h_o and h_c on $[\alpha, \infty)$, respectively $(-\infty, \beta]$, is not essential.

Remark 2.6. The smallness of μ_0 in Theorems 2.1–2.4 can be estimated.

3. Examples

Let us consider the problem

$$\ddot{y} + \ddot{y} + \dot{y} + y = \mu f(y), \quad (3.1)$$

where f is of the form

$$\alpha = -\delta, \quad \beta = \delta, \quad \delta > 0, \quad h_o = g + p, \quad h_c = g - p$$

with $p > 0$ constant and $g \in C(\mathbb{R}, \mathbb{R})$. We apply Theorem 2.2. Now we have

$$\phi_1(t) = \psi_1(t) = \sin t, \quad \phi_2(t) = \psi_2(t) = \cos t, \quad \eta(\theta, t) = \theta \sin t.$$

By taking $\mathcal{O} = (\delta, \infty)$, the conditions a) and b) of Theorem 2.2 are satisfied. Let $t_0 = \arcsin \frac{\delta}{\theta}$ for $\theta \in \mathcal{O}$. We compute (2.8) for this case

$$M(\omega, \theta) = (M_1(\omega, \theta), M_2(\omega, \theta)), \quad (3.2)$$

where

$$M_1(\omega, \theta) = \int_0^{2\pi} \omega(\theta \cos t - 2\theta \sin t - 3\theta \cos t) \sin t \, dt + \int_{t_0}^{t_0+\pi} (g(\theta \sin t) + p) \sin t \, dt \\ + \int_{t_0+\pi}^{t_0+2\pi} (g(\theta \sin t) - p) \sin t \, dt$$

$$= -2\pi\theta\omega + \int_0^{2\pi} g(\theta \sin t) \sin t \, dt + 4p \cos t_0$$

$$= -2\pi\theta\omega + 4p\sqrt{1 - \frac{\delta^2}{\theta^2}} + \int_0^{2\pi} g(\theta \sin t) \sin t \, dt,$$

$$M_2(\omega, \theta) = \int_0^{2\pi} \omega(\theta \cos t - 2\theta \sin t - 3\theta \cos t) \cos t \, dt + \int_{t_0}^{t_0+\pi} (g(\theta \sin t) + p) \cos t \, dt \\ + \int_{t_0+\pi}^{t_0+2\pi} (g(\theta \sin t) - p) \cos t \, dt$$

$$= -2\pi\theta\omega + \int_0^{2\pi} g(\theta \sin t) \cos t \, dt - 4p \sin t_0$$

$$= -2\pi\theta\omega - 4\frac{\delta p}{\theta}.$$

We have the following result.

THEOREM 3.1. *If there exist numbers $\delta < a_1 < a_2$ such that the numbers*

$$4p \left(\frac{\delta}{a_1} + \sqrt{1 - \frac{\delta^2}{a_1^2}} \right) + \int_0^{2\pi} g(a_1 \sin t) \sin t \, dt,$$

$$4p \left(\frac{\delta}{a_2} + \sqrt{1 - \frac{\delta^2}{a_2^2}} \right) + \int_0^{2\pi} g(a_2 \sin t) \sin t \, dt$$

have opposite signs, then there is a constant $K > 0$ such that for arbitrary sufficiently small μ there exist $\theta_\mu \in (a_1, a_2)$, $\omega_\mu \in (3D, D)$, $D = -\frac{\delta p}{2\pi} \left(\frac{1}{a_2^2} + \frac{1}{a_1^2} \right)$ and a $2\pi(1 + \mu\omega_\mu)$ -periodic solution y_μ of (3.1) satisfying

$$\sup_{t \in \mathbb{R}} \left| y_\mu(t) - \theta_\mu \sin \frac{t}{1 + \mu\omega_\mu} \right| \leq K|\mu|.$$

P r o o f. It is sufficient to verify (i) and (ii) of Theorem 2.2 when M is given by (3.2) and $\mathcal{B} = (3D, D) \times (a_1, a_2)$.

We put (3.2) in the homotopy

$$M(\omega, \theta, \lambda) = (M_1(\omega, \theta, \lambda), M_2(\omega, \theta, \lambda)), \quad \lambda \in [0, 1],$$

where

$$M_1(\omega, \theta, \lambda) = -2\pi\theta(\omega - 2(1 - \lambda)D) + 4p\sqrt{1 - \frac{\delta^2}{\theta^2}} \\ + \int_0^{2\pi} g(\theta \sin t) \sin t \, dt + 4\frac{\delta p}{\theta} - \lambda 4\frac{\delta p}{\theta},$$

$$M_2(\omega, \theta, \lambda) = -2\pi\theta(\omega - 2(1 - \lambda)D) - \lambda 4\frac{\delta p}{\theta}.$$

It is clear that

$$\forall \lambda \in [0, 1] \quad M(\partial\mathcal{B}, \lambda) \neq 0.$$

Consequently, we obtain

$$\deg(M(\cdot, \cdot, 1), \mathcal{B}, 0) = -\deg(M_1(2D, \cdot, 0), (a_1, a_2), 0) \neq 0.$$

The proof is finished by using Theorem 2.2. □

Let us take $g(x) = c_1 x + c_2$ with $c_{1,2}$ constant. We compute

$$4p \left(\frac{\delta}{\theta} + \sqrt{1 - \frac{\delta^2}{\theta^2}} \right) + \int_0^{2\pi} (c_1 \theta \sin t + c_2) \sin t \, dt \\ = 4p \left(\frac{\delta}{\theta} + \sqrt{1 - \frac{\delta^2}{\theta^2}} \right) + c_1 \theta \pi.$$

COROLLARY 3.2. *If $g(x) = c_1x + c_2$ in (3.1) with constant $c_{1,2}$ such that $c_1 < 0$ and $4p > -c_1\delta\pi$, then the conclusion of Theorem 3.1 holds.*

Proof. In Theorem 3.1, it is enough to take $a_1 > \delta$ near to δ and $a_2 > a_1$ sufficiently large. \square

Now we consider a forced problem of (3.1)

$$\ddot{y} + \dot{y} + y = \mu(f(y) + \sin t), \quad (3.3)$$

where f is given in (3.1). According to Theorem 2.4 and the computations for (3.2), the mapping (2.12) for (3.3) has the form

$$M(\omega, \theta) = (M_1(\omega, \theta), M_2(\omega, \theta)), \quad (3.4)$$

where

$$\begin{aligned} M_1(\omega, \theta) &= 4p\sqrt{1 - \frac{\delta^2}{\theta^2}} + \int_0^{2\pi} g(\theta \sin t) \sin t \, dt + \int_0^{2\pi} \sin(t + \omega) \sin t \, dt \\ &= 4p\sqrt{1 - \frac{\delta^2}{\theta^2}} + \int_0^{2\pi} g(\theta \sin t) \sin t \, dt + \pi \cos \omega, \end{aligned}$$

$$\begin{aligned} M_2(\omega, \theta) &= -4\frac{\delta p}{\theta} + \int_0^{2\pi} \sin(t + \omega) \cos t \, dt \\ &= -4\frac{\delta p}{\theta} + \pi \sin \omega. \end{aligned}$$

Assume that $4p = \pi$ and $\pi/2 < \omega < \pi$. Then the equations $M_1 = 0$, $M_2 = 0$ are equivalent to

$$\int_0^{2\pi} g\left(\frac{\sin t}{\sin \omega} \delta\right) \sin t \, dt = 0.$$

Theorem 2.4 implies the following result.

THEOREM 3.3. *Assume that $4p = \pi$ and $g \in C^1(\mathbb{R}, \mathbb{R})$. If the function*

$$\rho \mapsto \int_0^{2\pi} g(\delta\rho \sin t) \sin t \, dt$$

has a simple root $\rho_0 > 1$, then by putting $1/\rho_0 = \sin \omega_0$, $\pi/2 < \omega_0 < \pi$, there is a constant $K > 0$ such that for any μ sufficiently small there are (ω_μ, θ_μ) near to $(\omega_0, \delta\rho_0)$ and a 2π -periodic solution y_μ of (3.3) satisfying

$$\sup_{t \in \mathbb{R}} |y_\mu(t) - \theta_\mu \sin(t - \omega_\mu)| \leq K|\mu|.$$

P r o o f. If $\rho_0 > 1$ is a simple root of $\rho \mapsto \int_0^{2\pi} g(\delta\rho \sin t) \sin t \, dt$, then $\theta_0 = \delta\rho_0$, $1/\rho_0 = \sin \omega_0$, $\pi/2 < \omega_0 < \pi$ is a simple zero of $M = 0$ given by (3.4), i.e. $M(\omega_0, \theta_0) = 0$ and $DM(\omega_0, \theta_0)$ is invertible. The proof is finished by Theorem 2.4 when \mathcal{B} is taken as a small open neighbourhood of (ω_0, θ_0) . \square

Let us take $g(x) = c_1 x^3 + c_2 x$ with $c_{1,2}$ constant. Then

$$\int_0^{2\pi} g(\delta\rho \sin t) \sin t \, dt = \frac{3}{4}\pi c_1 \delta^3 \rho^3 + \pi \delta c_2 \rho.$$

Theorem 3.3 gives the next result.

COROLLARY 3.4. *Assume that $4p = \pi$. If $g(x) = c_1 x^3 + c_2 x$ in (3.1) with constant $c_{1,2}$ such that $c_1 c_2 < \frac{3}{4}c_1^2 \delta^2$, then the conclusion of Theorem 3.3 holds.*

P r o o f. The assumption $c_1 c_2 < \frac{3}{4}c_1^2 \delta^2$ implies the existence of a simple root $\rho_0 > 1$ of the equation

$$\frac{3}{4}\pi c_1 \delta^3 \rho^3 + \pi \delta c_2 \rho = 0.$$

\square

Now we assume that $g(x) = c_1 x$ with constant $c_1 > 0$ in (3.3). Then (3.4) has the form

$$\begin{aligned} M_1(\omega, \theta) &= 4p\sqrt{1 - \frac{\delta^2}{\theta^2}} + \pi \cos \omega + c_1 \theta \pi, \\ M_2(\omega, \theta) &= -4\frac{\delta p}{\theta} + \pi \sin \omega. \end{aligned}$$

By assuming $\pi > 4p$, the equation $M(\omega, \theta) = 0$ with $\theta > \delta$ and $\pi/2 < \omega < \pi$ is equivalent to

$$4p\sqrt{1 - \frac{\delta^2}{\theta^2}} - \pi\sqrt{1 - \frac{16\delta^2 p^2}{\theta^2 \pi^2}} + c_1 \theta \pi = 0,$$

i.e.

$$8\pi c_1 p \sqrt{\theta^2 - \delta^2} + c_1^2 \theta^2 \pi^2 = \pi^2 - 16p^2. \quad (3.5)$$

If $\pi^2 - 16p^2 > c_1^2 \delta^2 \pi^2$, then (3.5) has a unique simple root

$$\theta_0 = \sqrt{\left(\frac{-4p + \pi\sqrt{1 - \delta^2 c_1^2}}{c_1 \pi}\right)^2 + \delta^2}. \quad (3.6)$$

Like for Corollary 3.4, we obtain

THEOREM 3.5. Assume that $g(x) = c_1 x$ with constant $c_1 > 0$ such that $\pi^2 - 16p^2 > c_1^2 \delta^2 \pi^2$. Then there exists a constant $K > 0$ such that for arbitrary sufficiently small μ there exist (ω_μ, θ_μ) near to (ω_0, θ_0) given by (3.6) and $\pi/2 < \omega_0 < \pi$, $\sin \omega_0 = \frac{4\delta p}{\pi \theta_0}$, and a 2π -periodic solution y_μ of (3.3) satisfying

$$\sup_{t \in \mathbb{R}} |y_\mu(t) - \theta_\mu \sin(t - \omega_\mu)| \leq K|\mu|.$$

Similarly we have

THEOREM 3.6. Assume that $g(x) = c_1 x$ with constant $c_1 < 0$ such that $16p^2(1 - c_1^2 \delta^2) > \pi^2$. Then there exists a constant $K > 0$ such that for arbitrary sufficiently small μ there exist (ω_μ, θ_μ) near (ω_0, θ_0) given by

$$\theta_0 = \frac{1}{\pi} \sqrt{\left(\frac{\pi - 4p\sqrt{1 - \delta^2 c_1^2}}{c_1} \right)^2 + 16\delta^2 p^2},$$

$$\sin \omega_0 = \frac{4\delta p}{\pi \theta_0}, \quad \pi/2 < \omega_0 < \pi,$$

and a 2π -periodic solution y_μ of (3.3) satisfying

$$\sup_{t \in \mathbb{R}} |y_\mu(t) - \theta_\mu \sin(t - \omega_\mu)| \leq K|\mu|.$$

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Received May 24, 1996

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