

Benjamin Cahen

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# Berezin-Weyl quantization for Cartan motion groups

BENJAMIN CAHEN

*Abstract.* We construct adapted Weyl correspondences for the unitary irreducible representations of the Cartan motion group of a noncompact semisimple Lie group by using the method introduced in [B. Cahen, *Weyl quantization for semidirect products*, Differential Geom. Appl. **25** (2007), 177–190].

*Keywords:* semidirect product, Cartan motion group, unitary representation, semisimple Lie group, symplectomorphism, coadjoint orbit, Weyl quantization, Berezin quantization

*Classification:* 22E45, 22E46, 22E70, 22E15, 81S10, 81R05

## 1. Introduction

In [3] and [4], we introduced the notion of adapted Weyl correspondence as a generalization of the usual quantization rules [1], [15]. The present paper is part of a larger program to study adapted Weyl correspondences for semisimple Lie groups and for semidirect products.

Let  $G$  be a connected Lie group,  $\mathfrak{g}$  the Lie algebra of  $G$  and  $\mathfrak{g}^*$  the dual space of  $\mathfrak{g}$ . Let  $\pi$  be a unitary irreducible representation of  $G$  on a Hilbert space  $\mathcal{H}$ . Suppose that  $\pi$  is associated with an orbit  $\mathcal{O}$  for the coadjoint action of  $G$  on  $\mathfrak{g}^*$  by the Kostant-Kirillov method of orbits [18], [20]. In [6], we gave the following definition for the notion of adapted Weyl correspondence.

**Definition 1.1.** An adapted Weyl correspondence is an isomorphism  $W$  from a vector space  $\mathcal{A}$  of complex-valued smooth functions on the orbit  $\mathcal{O}$  (called symbols) onto a vector space  $\mathcal{B}$  of (not necessarily bounded) linear operators on  $\mathcal{H}$  satisfying the following properties:

- (1) the elements of  $\mathcal{B}$  preserve a fixed dense domain  $\mathcal{D}$  of  $\mathcal{H}$ ;
- (2) the constant function 1 belongs to  $\mathcal{A}$ , the identity operator  $I$  belongs to  $\mathcal{B}$  and  $W(1) = I$ ;
- (3)  $A \in \mathcal{B}$  and  $B \in \mathcal{B}$  implies  $AB \in \mathcal{B}$ ;
- (4) for each  $f$  in  $\mathcal{A}$  the complex conjugate  $\bar{f}$  of  $f$  belongs to  $\mathcal{A}$  and the adjoint of  $W(f)$  is an extension of  $W(\bar{f})$ ;
- (5) the elements of  $\mathcal{D}$  are  $C^\infty$ -vectors for the representation  $\pi$ , the functions  $\tilde{X}$  ( $X \in \mathfrak{g}$ ) defined on  $\mathcal{O}$  by  $\tilde{X}(\xi) = \langle \xi, X \rangle$  are in  $\mathcal{A}$  and  $W(i\tilde{X})v = d\pi(X)v$  for each  $X \in \mathfrak{g}$  and each  $v \in \mathcal{D}$ .

Adapted Weyl correspondences were obtained in various situations, see the introduction of [6]. In particular, we constructed adapted Weyl correspondences for the principal series representations of a noncompact semisimple Lie group in [3] and [7]. We also obtained adapted Weyl correspondences for the unitary irreducible representations of the semidirect product  $V \rtimes K$  of the real vector space  $V$  by a Lie group  $K$  acting linearly on  $V$  in the following situations:

- (1)  $K$  is a connected noncompact semisimple Lie group and the little group associated with the representation of  $V \times K$  is a maximal compact subgroup of  $K$  [6];
- (2)  $K$  is a connected compact semisimple Lie group and the little group is the centralizer of a torus of  $K$  [10].

Let us mention that adapted Weyl correspondences have various applications in Harmonic Analysis and Deformation Theory as the construction of covariant star-products on coadjoint orbits [3] and the study of contractions of Lie group unitary representations [13], [5], [8], [9]. Recently, in [11], we have studied a contraction of the principal series of a semisimple Lie group to the unitary irreducible representations of its Cartan motion group by using the deformed Weyl calculus introduced in [3].

The present paper can be considered as a sequel of [6] and [10]. Let  $G_0$  be a connected noncompact semisimple Lie group with Lie algebra  $\mathfrak{g}_0$  and let  $K$  be a maximal compact subgroup of  $G_0$ . Then we have the Cartan decomposition  $\mathfrak{g}_0 = \mathfrak{k} \oplus V$  where  $\mathfrak{k}$  is the Lie algebra of  $K$  and  $V$  is an  $\text{Ad}(K)$ -invariant subspace of  $\mathfrak{g}_0$ . The Cartan motion group associated with the pair  $(G_0, K)$  is the semidirect product  $V \rtimes K$  formed with respect to the adjoint action of  $K$  on  $V$ .

It is known for a long time that the unitary irreducible representations of  $V \rtimes K$  are similar to the principal series representations of  $G_0$  [21]. This has been illustrated by means of contractions of representations in [14] (see also [11]). Here, we exploit this similarity to construct adapted Weyl correspondences for unitary irreducible representations of  $V \rtimes K$  as it was done for principal series representations of  $G_0$  in [7]. The method is essentially the same as in [6] and the explicit computations are partially based on those of [7].

More precisely, let  $\mathcal{O}$  be a coadjoint orbit of  $V \rtimes K$  which is associated with a generic unitary irreducible representation  $\pi$  of  $V \rtimes K$ . We realize  $\pi$  on a Hilbert space of functions on  $\mathbb{R}^n$  where  $n = (1/2)\dim \mathcal{O}$  and we compute the corresponding derived representation  $d\pi$  (Section 3). We dequantize  $d\pi$  by using a combination of the usual Weyl calculus on  $\mathbb{R}^{2n}$  and of the Berezin calculus on the little group orbit  $\mathcal{O}'$  (Section 4). Then we obtain an explicit symplectomorphism from  $\mathbb{R}^{2n} \times \mathcal{O}'$  onto a dense open subset of  $\mathcal{O}$  and an adapted Weyl correspondence on  $\mathcal{O}$  (Section 5). In the case when  $G_0$  is a complex Lie group, we verify that the adapted Weyl correspondence coincide with that of [10].

## 2. Preliminaries

In this section, we introduce some general facts on noncompact semisimple Lie groups and Cartan motion groups. Our main references are [16, Chapter VI], [19, Chapter V] and [22].

Let  $G_0$  be a connected noncompact semisimple real Lie group with finite center. Let  $\mathfrak{g}_0$  be the Lie algebra of  $G_0$ . We denote by  $\beta$  the Killing form of  $\mathfrak{g}_0$  defined by  $\beta(X, Y) = \text{Tr}(\text{ad } X \text{ ad } Y)$  for  $X$  and  $Y$  in  $\mathfrak{g}_0$ . Let  $\theta$  be a Cartan involution of  $\mathfrak{g}_0$  and let  $\mathfrak{g}_0 = \mathfrak{k} \oplus V$  be the corresponding Cartan decomposition of  $\mathfrak{g}_0$ . Let  $K$  be the connected compact (maximal) subgroup of  $G_0$  with Lie algebra  $\mathfrak{k}$ . Let  $\mathfrak{a}$  be a maximal abelian subalgebra of  $V$  and let  $M$  be the centralizer of  $\mathfrak{a}$  in  $K$ . Let  $\mathfrak{m}$  denote the Lie algebra of  $M$ . Let  $\Delta := \Delta(\mathfrak{a}, \mathfrak{g}_0)$  be the set of restricted roots and let

$$\mathfrak{g}_0 = \mathfrak{a} \oplus \mathfrak{m} \oplus \sum_{\lambda \in \Delta} \mathfrak{g}_\lambda$$

be the root space decomposition of  $\mathfrak{g}_0$ . We fix a Weyl chamber in  $\mathfrak{a}$  and we denote by  $\Delta^+$  the corresponding set of positive roots. We set  $\mathfrak{n} = \sum_{\lambda \in \Delta^+} \mathfrak{g}_\lambda$  and  $\bar{\mathfrak{n}} = \sum_{\lambda \in \Delta^+} \mathfrak{g}_{-\lambda}$ . Then  $\bar{\mathfrak{n}} = \theta(\mathfrak{n})$ . Let  $A, N$  and  $\bar{N}$  denote the analytic subgroups of  $G$  with algebras  $\mathfrak{a}, \mathfrak{n}$  and  $\bar{\mathfrak{n}}$ , respectively.

Recall that  $\bar{N}MAN$  is an open dense subset of  $G$ . We denote by  $g = \bar{n}(g)m(g)a(g)n(g)$  the decomposition of  $g \in \bar{N}MAN$ . Also, recall that we have the Iwasawa decomposition  $G = KAN$ . We denote by  $g = \tilde{k}(g)\tilde{a}(g)\tilde{n}(g)$  the decomposition of  $g \in G$ .

The Cartan motion group associated with the pair  $(G_0, K)$  is the semidirect product  $G := V \rtimes K$ . The group law of  $G$  is given by

$$(v, k).(v', k') = (v + \text{Ad}(k)v', kk')$$

for  $v, v'$  in  $V$  and  $k, k' \in K$ . The Lie algebra  $\mathfrak{g}$  of  $G$  is the space  $V \times \mathfrak{k}$  endowed with the Lie bracket

$$[(w, U), (w', U')] = ([U, w']_0 - [U', w]_0, [U, U']_0)$$

where  $[\cdot, \cdot]_0$  denotes the Lie bracket of  $\mathfrak{g}_0$ .

Recall that  $\beta$  is positive defined on  $V$  and negative defined on  $\mathfrak{k}$  [16, p. 184]. Then, by using  $\beta$ , we can identify  $V^*$  with  $V$  and  $\mathfrak{k}^*$  with  $\mathfrak{k}$ , hence  $\mathfrak{g}^* \simeq V^* \times \mathfrak{k}^*$  with  $V \times \mathfrak{k}$ . Under this identification, the coadjoint action of  $G$  on  $\mathfrak{g}^* \simeq V \times \mathfrak{k}$  is given by

$$(v, k) \cdot (w, U) = (\text{Ad}(k)w, \text{Ad}(k)U + [v, \text{Ad}(k)w]_0)$$

for  $v, w$  in  $V, k$  in  $K$  and  $U$  in  $\mathfrak{k}$ . This is a particular case of the general formula for the coadjoint action of a semidirect product, see for instance [22].

Let  $p_{\mathfrak{k}}^c$  and  $p_V^c$  be the projections of  $\mathfrak{g}_0$  on  $\mathfrak{k}$  and  $V$  associated with the decomposition  $\mathfrak{g}_0 = \mathfrak{k} \oplus V$ . Recall that an element  $\xi_1$  of  $\mathfrak{a}$  is said to be regular if  $\lambda(\xi_1) \neq 0$  for each  $\lambda \in \Delta$ . We shall need the following lemma.

**Lemma 2.1.** *For each regular element  $\xi_1$  of  $\mathfrak{a}$ , the space  $\text{ad } \xi_1(V)$  is the orthogonal complement of  $\mathfrak{m}$  in  $\mathfrak{k}$ .*

PROOF: For each  $\lambda \in \Delta^+$ , let  $E_\lambda \neq 0$  be in  $\mathfrak{g}_\lambda$ . Note that the space  $p_{\mathfrak{k}}^c(\mathfrak{n}) = p_{\mathfrak{k}}^c(\bar{\mathfrak{n}})$  is generated by the elements  $E_\lambda + \theta(E_\lambda)$  and hence orthogonal to  $\mathfrak{m}$ . Now, by applying successively  $p_{\mathfrak{k}}^c$  and  $p_V^c$  to the decomposition  $\mathfrak{g}_0 = \mathfrak{m} + \mathfrak{a} + \mathfrak{n} + \bar{\mathfrak{n}}$  we get the decompositions  $\mathfrak{k} = \mathfrak{m} + p_{\mathfrak{k}}^c(\mathfrak{n})$  and  $V = \mathfrak{a} + p_V^c(\mathfrak{n})$ . This shows that  $p_{\mathfrak{k}}^c(\mathfrak{n})$  is the orthogonal complement of  $\mathfrak{m}$  in  $\mathfrak{k}$ . On the other hand, since  $p_V^c(\mathfrak{n})$  is generated by the elements  $E_\lambda - \theta(E_\lambda)$ , we see that the space  $\text{ad } \xi_1(V)$  is generated by the elements

$$\text{ad } \xi_1(E_\lambda - \theta(E_\lambda)) = \lambda(\xi_1)(E_\lambda + \theta(E_\lambda))$$

where  $\lambda(\xi_1) \neq 0$  for  $\lambda \in \Delta$ . Hence  $\text{ad } \xi_1(V) = p_{\mathfrak{k}}^c(\mathfrak{n})$  is the orthogonal complement of  $\mathfrak{m}$  in  $\mathfrak{k}$ . □

The coadjoint orbits of the semidirect product of a Lie group by a vector space were described in [22]. For each  $(w, U) \in \mathfrak{g}^* \simeq \mathfrak{g}$ , we denote by  $O(w, U)$  the orbit of  $(w, U)$  under the coadjoint action of  $G$ . The following lemma shows that, for almost all  $(w, U)$ , the orbit  $O(w, U)$  is of the form  $O(\xi_1, \xi_2)$  with  $\xi_1 \in \mathfrak{a}$  and  $\xi_2 \in \mathfrak{m}$ .

**Lemma 2.2.** (1) *Let  $\mathcal{O}$  be a coadjoint orbit for the coadjoint action of  $G$  on  $\mathfrak{g}^* \simeq \mathfrak{g}$ . Then there exists an element of  $\mathcal{O}$  of the form  $(\xi_1, U)$  with  $\xi_1 \in \mathfrak{a}$ . Moreover, if  $\xi_1$  is regular then there exists  $\xi_2 \in \mathfrak{m}$  such that  $(\xi_1, \xi_2) \in \mathcal{O}$ .*  
 (2) *Let  $\xi_1$  be a regular element of  $\mathfrak{a}$ . Then  $M$  is the stabilizer of  $\xi_1$  in  $K$ .*

PROOF: (1) Let  $(w, U) \in \mathcal{O}$ . For each  $k \in K$  we have

$$(0, k) \cdot (w, U) = (\text{Ad}(k)w, \text{Ad}(k)U).$$

By [19, p. 120], we have  $\text{Ad}(K)\mathfrak{a} = V$  and then one can choose  $k \in K$  so that  $\text{Ad}(k)w \in \mathfrak{a}$ . We set  $\xi_1 := \text{Ad}(k)w$ . If we assume that  $\xi_1$  is regular then, by Lemma 2.1, we can write  $U = \xi_2 + [\xi_1, v]$  where  $\xi_2 \in \mathfrak{m}$  and  $v \in V$ . Then  $(\xi_1, U) = (v, e) \cdot (\xi_1, \xi_2)$ . Hence  $\mathcal{O} = O(\xi_1, \xi_2)$ .

(2) By [7, Lemma 4.2], the stabilizer of  $\xi_1$  in  $\mathfrak{g}_0$  is  $MA$ . Then, the stabilizer of  $\xi_1$  in  $K$  is  $MA \cap K = M$ . □

Let  $\xi_1 \in \mathfrak{a}$  be a regular element. Denote by  $O_V(\xi_1)$  the orbit of  $\xi_1$  in  $V$  under the action of  $K$ . In the next section, we shall need the chart on  $O_V(\xi_1) \simeq K/M$  which is given by the following lemma.

**Lemma 2.3** ([26, Lemma 7.6.8]). *The map  $\tau : y \rightarrow \text{Ad}(\tilde{k}(y))\xi_1$  is a diffeomorphism from  $\bar{N}$  onto a dense open subset of  $O_V(\xi_1)$ . Let us consider the action of  $k \in K$  on  $y \in \bar{N}$  defined by  $\tau(k \cdot y) = \text{Ad}(k)\tau(y)$  or, equivalently, by  $k \cdot y = \bar{n}(ky)$ . Then the  $K$ -invariant measure on  $\bar{N}$  is given by  $e^{-2\rho(\log \tilde{a}(y))} dy$  where  $dy$  is a Haar measure on  $\bar{N}$ .*

In the rest of the paper, we fix the normalization of  $dy$  as follows. Let  $(E_i)_{1 \leq i \leq n}$  be an orthonormal basis of  $\bar{\mathfrak{n}}$  with respect to the scalar product defined by

$(Y, Z) := -\beta(Y, \theta(Z))$ . Denote by  $(y_1, y_2, \dots, y_n)$  the coordinates of  $Y \in \bar{\mathfrak{n}}$  in this basis and let  $dY = dy_1 dy_2 \dots dy_n$  be the Euclidean measure on  $\bar{\mathfrak{n}}$ . The exponential map  $\exp$  is a diffeomorphism from  $\bar{\mathfrak{n}}$  onto  $\bar{N}$  and we set  $dy := (\exp^{-1})^*(dY)$ .

We shall also denote by  $k \cdot Y$  the action of  $k \in K$  on  $Y \in \bar{\mathfrak{n}}$  defined by  $\exp(k \cdot Y) = k \cdot \exp(Y)$ .

### 3. Representations

We retain the notation of Section 2. Let  $\xi_1 \in \mathfrak{a}$  be a regular element and let  $\xi_2 \in \mathfrak{m}$ . We denote by  $o(\xi_2)$  the orbit of  $\xi_2$  under the adjoint action of  $M$  on  $\mathfrak{m}$ . Let  $\sigma$  be a unitary irreducible representation of  $M$  on a complex (finite-dimensional) vector space  $E$ . In the rest of the paper, we assume that  $\sigma$  is associated with the orbit  $o(\xi_2)$  in the following sense, see [27, Section 4]. Given a maximal torus  $T$  of  $M$  with Lie algebra  $\mathfrak{t}$  and a set of positive roots in  $\Delta(\mathfrak{t}^{\mathbb{C}}, \mathfrak{m}^{\mathbb{C}})$ , the element  $i\beta(\xi_2, \cdot)$  of  $i\mathfrak{t}^*$  is the highest weight of  $\sigma$ . Under these assumptions, the orbit  $O(\xi_1, \xi_2)$  is associated with the unitarily induced representation

$$\hat{\pi} = \text{Ind}_{V \times M}^G \left( e^{i\beta(\xi_1, \cdot)} \otimes \sigma \right)$$

(see [20], [22] and [23]). By a result of Mackey,  $\hat{\pi}$  is irreducible since  $\sigma$  is irreducible [24, p. 149]. Moreover, in the terminology of [22] and [23], the group  $M$  is called the little group and the orbit  $o(\xi_2)$  the little orbit.

The representation  $\hat{\pi}$  is usually realized on the space of square-integrable sections of a Hermitian vector bundle over  $O_V(\xi_1)$ , see [10], [22]. Following [23] and using Lemma 2.3 and the section  $y \rightarrow \tilde{k}(y)$  from  $\bar{N}$  to  $K$ , we immediately obtain the realization  $\pi_0$  of  $\hat{\pi}$  defined by

$$(\pi_0(v, k)\psi)(y) = e^{i\beta(\text{Ad}(\tilde{k}(y))\xi_1, v)} \sigma(\tilde{k}(y)^{-1}k\tilde{k}(k^{-1} \cdot y))\psi(k^{-1} \cdot y)$$

on the Hilbert space  $\mathcal{H}_0$  which is the completion of the space of compactly supported smooth functions  $\psi : \bar{N} \rightarrow E$  with respect to the norm

$$\|\psi\|_0^2 = \int_{\bar{N}} \langle \psi(y), \psi(y) \rangle_E e^{-2\rho(\log \bar{a}(y))} dy.$$

For the Weyl calculus, it is more convenient to realize  $\hat{\pi}$  on the Hilbert space  $\mathcal{H} := L^2(\bar{\mathfrak{n}}, E)$  which is the completion of the space  $C_0(\bar{\mathfrak{n}}, E)$  of compactly supported smooth functions  $\phi : \bar{\mathfrak{n}} \rightarrow E$  with respect to the norm

$$\|\phi\|^2 = \int_{\bar{\mathfrak{n}}} \langle \phi(Y), \phi(Y) \rangle_E dY.$$

To this end, we use the unitary isomorphism  $B$  from  $\mathcal{H}$  onto  $\mathcal{H}_0$  defined by  $B(\phi)(\exp Y) = e^{\rho(\log \bar{a}(y))}\phi(Y)$  and we set  $\pi(g) := B^{-1}\pi_0(g)B$  for  $g \in G$ . We immediately obtain, for  $(v, k) \in G$ ,

$$(\pi(v, k)\phi)(Y) = e^{i\beta(\text{Ad}(\tilde{k}(y))\xi_1, v) + \rho(\log \bar{a}(k^{-1} \cdot y) - \log \bar{a}(y))} \sigma(m(k, y))\phi(k^{-1} \cdot Y)$$

where we have set  $y = \exp Y$  and  $m(k, y) := \tilde{k}(y)^{-1}k\tilde{k}(k^{-1} \cdot y) \in M$ . This formula can be simplified as follows. Let  $k \in K$  and  $y \in \bar{N}$ . Write  $k^{-1}y = \bar{n}(k^{-1}y)m(k^{-1}y)a(k^{-1}y)n(k^{-1}y)$ . Then  $k^{-1}\tilde{k}(y) = \tilde{k}(\bar{n}(k^{-1}y))m(k^{-1}y)$ . Thus  $m(k, y) = m(k^{-1}y)^{-1}$ . Also, we have that

$$\tilde{a}(y) = \tilde{a}(k^{-1}y) = \tilde{a}(\bar{n}(k^{-1}y))a(k^{-1}y) = \tilde{a}(k^{-1} \cdot y)a(k^{-1}y).$$

This gives

$$\begin{aligned} (\pi(v, k)\phi)(Y) &= e^{-\rho(\log a(k^{-1} \exp Y)) + i\beta(\text{Ad}(\tilde{k}(\exp Y))\xi_{1, v})} \sigma(m(k^{-1} \exp Y))^{-1} \\ &\quad \times \phi(\log \bar{n}(k^{-1} \exp Y)). \end{aligned}$$

Now, we compute the derived representation  $d\pi$ . Let  $p_{\mathfrak{a}}$ ,  $p_{\mathfrak{m}}$  and  $p_{\bar{\mathfrak{n}}}$  be the projections of  $\mathfrak{g}_0$  onto  $\mathfrak{a}$ ,  $\mathfrak{m}$  and  $\bar{\mathfrak{n}}$  associated with the direct decomposition  $\mathfrak{g}_0 = \mathfrak{a} + \mathfrak{m} + \mathfrak{n} + \bar{\mathfrak{n}}$ . For  $X \in \bar{\mathfrak{n}}$  we denote by  $X^+$  the right invariant vector field on  $\bar{N}$  generated by  $X$ , that is,  $X^+(y) = \frac{d}{dt}(\exp tX)y|_{t=0}$  for  $y \in \bar{N}$ .

**Lemma 3.1** ([7]). (1) For each  $X \in \bar{\mathfrak{n}}$  and each  $Y \in \bar{\mathfrak{n}}$ , we have

$$d \log(\exp Y) (X^+(\exp Y)) = \frac{\text{ad } Y}{e^{\text{ad } Y} - 1} (X).$$

(2) For each  $X \in \mathfrak{g}_0$  and each  $y \in \bar{N}$ , we have

$$\begin{aligned} \left. \frac{d}{dt} a(\exp(tX)y) \right|_{t=0} &= p_{\mathfrak{a}}(\text{Ad}(y^{-1})X) \\ \left. \frac{d}{dt} m(\exp(tX)y) \right|_{t=0} &= p_{\mathfrak{m}}(\text{Ad}(y^{-1})X) \\ \left. \frac{d}{dt} \bar{n}(\exp(tX)y) \right|_{t=0} &= (\text{Ad}(y)p_{\bar{\mathfrak{n}}}(\text{Ad}(y^{-1})X))^+(y). \end{aligned}$$

From this lemma, we immediately deduce the following proposition.

**Proposition 3.2.** For each  $(w, U) \in \mathfrak{g}$ ,  $\phi \in C_0(\bar{\mathfrak{n}}, E)$  and  $Y \in \bar{\mathfrak{n}}$ , we have

$$\begin{aligned} (d\pi(w, U)\phi)(Y) &= i\beta(\text{Ad}(\tilde{k}(\exp Y))\xi_1, w)\phi(Y) \\ &\quad + \rho(p_{\mathfrak{a}}(\text{Ad}(\exp(-Y))U))\phi(Y) + d\sigma(p_{\mathfrak{m}}(\text{Ad}(\exp(-Y))U))\phi(Y) \\ &\quad - d\phi(Y) \left( \frac{\text{ad } Y}{1 - e^{-\text{ad } Y}} p_{\bar{\mathfrak{n}}}(\text{Ad}(\exp(-Y))U) \right). \end{aligned}$$

#### 4. Dequantization

In this section, we first introduce the Berezin-Weyl calculus on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$ . Recall that the Berezin calculus on  $o(\xi_2)$  associates with each operator  $B$  on the finite-dimensional complex vector space  $E$  a complex-valued function  $s(B)$  on the orbit  $o(\xi_2)$  called the symbol of the operator  $B$  (see [2]). The following properties of the Berezin calculus are well-known, see [12], [3], [10].

**Proposition 4.1.** (1) *The map  $B \rightarrow s(B)$  is injective.*

(2) *For each operator  $B$  on  $E$ , we have  $s(B^*) = \overline{s(B)}$ .*

(3) *For each operator  $B$  on  $E$ , each  $m \in M$  and each  $\varphi \in o(\xi_2)$ , we have*

$$s(B)(\text{Ad}(m)\varphi) = s(\sigma(m)^{-1}B\sigma(m))(\varphi).$$

(4) *For  $X \in \mathfrak{m}$  and  $\varphi \in o(\xi_2)$ , we have  $s(d\sigma(X))(\varphi) = i\beta(\varphi, X)$ .*

In particular, we see that the map  $s^{-1}$  is an adapted Weyl transform on  $o(\xi_2)$  in the sense of Definition 1.1.

We say that a complex-valued smooth function  $f : (Y, Z, \varphi) \rightarrow f(Y, Z, \varphi)$  is a symbol on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$  if for each  $(Y, Z) \in \bar{\mathfrak{n}}^2$  the function  $\varphi \rightarrow f(Y, Z, \varphi)$  is the symbol in the Berezin calculus on  $o(\xi_2)$  of an operator on  $E$  denoted by  $\hat{f}(Y, Z)$ . A symbol  $f$  on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$  is called an  $S$ -symbol if the function  $\hat{f}$  belongs to the Schwartz space of rapidly decreasing smooth functions on  $\bar{\mathfrak{n}}^2$  with values in  $\text{End}(E)$ . The Weyl calculus for  $\text{End}(E)$ -valued functions is a slight refinement of the usual Weyl calculus for complex-valued functions [17], [15]. For each  $S$ -symbol on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$ , we define the operator  $\mathcal{W}(f)$  on the Hilbert space  $L^2(\bar{\mathfrak{n}}, E)$  by

$$(4.1) \quad (\mathcal{W}(f)\phi)(Y) = (2\pi)^{-n} \int_{\bar{\mathfrak{n}}^2} e^{i\langle T, Z \rangle} \hat{f}\left(Y + \frac{1}{2}T, Z\right) \phi(Y + T) dT dZ$$

for  $\phi \in C_0(\bar{\mathfrak{n}}, E)$ .

It is well-known that the Weyl calculus can be extended to much larger classes of symbols (see for instance [17]). Here we only consider a class of polynomial symbols. We say that a symbol  $f$  on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$  is a  $P$ -symbol if the function  $\hat{f}(Y, Z)$  is polynomial in the variable  $Z$ . Let  $f$  be the  $P$ -symbol defined by  $f(Y, Z, \varphi) = u(Y)Z^\alpha$  where  $u \in C^\infty(\bar{\mathfrak{n}})$  and  $Z^\alpha := z_1^{\alpha_1} z_2^{\alpha_2} \dots z_n^{\alpha_n}$  for each multi-index  $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_n)$ . Then, by [25], we have

$$(4.2) \quad (\mathcal{W}(f)\phi)(Y) = \left(i \frac{\partial}{\partial Z}\right)^\alpha \left(u\left(Y + \frac{1}{2}Z\right) \phi(Y + Z)\right) \Big|_{Z=0}.$$

In particular, if  $f(Y, Z, \varphi) = u(Y)$  then  $(\mathcal{W}(f)\phi)(Y) = u(Y)\phi(Y)$  and if  $f(Y, Z, \varphi) = u(Y)z_k$  then

$$(4.3) \quad (\mathcal{W}(f)\phi)(Y) = i \left(\frac{1}{2} \partial_k u(Y) \phi(Y) + u(Y) \partial_k \phi(Y)\right)$$

where  $\partial_k$  denotes partial derivative with respect to the variable  $y_k$ .

We need the following lemma. The trace of an endomorphism  $u$  of  $\bar{\mathfrak{n}}$  is denoted by  $\text{Tr}_{\bar{\mathfrak{n}}} u$ .

**Lemma 4.2.** *For  $U \in \mathfrak{k}$  let  $c_U : \bar{\mathfrak{n}} \rightarrow \bar{\mathfrak{n}}$  be the map defined by*

$$c_U(Y) = s(\text{ad } Y)p_{\bar{\mathfrak{n}}}(\text{Ad}(\exp(-Y))U)$$

where  $s$  is the function defined by  $s(z) = \frac{e^z}{1-e^{-z}}$  for  $z \neq 0$  and  $s(0) = 1$ . Then we have

$$\mathrm{Tr}_{\bar{\mathfrak{n}}} dc_U(Y) = -2\rho(p_{\mathfrak{a}}(\mathrm{Ad}(\exp(-Y))X)).$$

PROOF: This is a particular case of [7, Lemma 3.3]. □

Then we get the following proposition.

**Proposition 4.3.** *For each  $(w, U) \in \mathfrak{g}$ , the Berezin-Weyl symbol of the operator  $-id\pi(w, U)$  is the P-symbol  $f_{(w,U)}$  on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$  defined by*

$$f_{(w,U)}(Y, Z, \varphi) = \beta(\mathrm{Ad}(\tilde{k}(\exp Y))\xi_1, w) + \beta(p_{\mathfrak{m}}(\mathrm{Ad}(\exp(-Y))U), \varphi) + (c_U(Y), Z).$$

PROOF: Set  $c_U^k(Y) = (c_U(Y), E_k)$  for each  $k = 1, 2, \dots, n$ . By using (4) of Proposition 4.1 and Formula (4.3), we immediately see that the symbol of  $-id\pi(w, U)$  is

$$f_{(w,U)}(Y, Z, \varphi) = -i\rho(p_{\mathfrak{a}}(\mathrm{Ad}(\exp(-Y))U)) + \beta(\mathrm{Ad}(\tilde{k}(\exp Y))\xi_1, w) + \beta(p_{\mathfrak{m}}(\mathrm{Ad}(\exp(-Y))U), \varphi) + \sum_{k=1}^n c_U^k(Y)z_k - \frac{i}{2} \sum_{k=1}^n \partial_k c_U^k(Y).$$

But by Lemma 4.2, we have

$$-\frac{i}{2} \sum_{k=1}^n \partial_k c_U^k(Y) = -\frac{i}{2} \mathrm{Tr}_{\bar{\mathfrak{n}}}(dc_U(Y)) = i\rho(p_{\mathfrak{a}}(\mathrm{Ad}(\exp(-Y))U)).$$

The result follows. □

### 5. Adapted Weyl correspondence

In this section, we use the dequantization procedure of Section 4 in order to obtain an explicit diffeomorphism from  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$  onto the dense open subset  $\tilde{O}(\xi_1, \xi_2)$  of  $O(\xi_1, \xi_2)$  defined by

$$\tilde{O}(\xi_1, \xi_2) = \{(v, k) \cdot (\xi_1, \xi_2) : v \in V, k \in K \cap \bar{N}MAN\}$$

and then to construct an adapted Weyl correspondence on  $O(\xi_1, \xi_2)$ .

**Proposition 5.1.** *Let  $\Psi$  be the map from  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$  to  $\mathfrak{g}$  defined by*

$$\Psi(Y, Z, \varphi) = \left( \mathrm{Ad}(\tilde{k}(\exp Y))\xi_1, p_{\mathfrak{t}}^c \left( \mathrm{Ad}(\exp Y) \left( \varphi + p_{\mathfrak{n}} \left( \frac{\mathrm{ad} Y}{e^{\mathrm{ad} Y} - 1} \theta(Z) \right) \right) \right) \right).$$

Then, for each  $(w, U) \in \mathfrak{g}$ , we have

$$f_{(w,U)}(Y, Z, \phi) = \langle \Psi(Y, Z, \varphi), (w, U) \rangle.$$

PROOF: We use Proposition 4.3. Note that we have  $\beta(\mathfrak{a} + \mathfrak{m}, \mathfrak{n} + \bar{\mathfrak{n}}) = (0)$ ,  $\beta(\mathfrak{n}, \mathfrak{n}) = (0)$  and  $\beta(\bar{\mathfrak{n}}, \bar{\mathfrak{n}}) = (0)$ . Then for  $(Y, Z, \varphi) \in \bar{\mathfrak{n}}^2 \times o(\xi_2)$  and  $(w, U) \in \mathfrak{g}$ , we can write

$$\begin{aligned} (c_U(Y), Z) &= -\beta(c_U(Y), Z) \\ &= -\beta\left(\frac{\text{ad } Y}{1 - e^{-\text{ad } Y}} p_{\bar{\mathfrak{n}}}(\text{Ad}(\exp(-Y))U), \theta(Z)\right) \\ &= \beta\left(p_{\bar{\mathfrak{n}}}(\text{Ad}(\exp(-Y))U), \frac{\text{ad } Y}{e^{\text{ad } Y} - 1} \theta(Z)\right) \\ &= \beta\left(\text{Ad}(\exp(-Y))U, p_{\mathfrak{n}}\left(\frac{\text{ad } Y}{e^{\text{ad } Y} - 1} \theta(Z)\right)\right) \\ &= \beta\left(U, p_{\mathfrak{k}}^c\left(\text{Ad}(\exp Y) p_{\mathfrak{n}}\left(\frac{\text{ad } Y}{e^{\text{ad } Y} - 1} \theta(Z)\right)\right)\right). \end{aligned}$$

Similarly, we have

$$\beta(\varphi, p_{\mathfrak{m}}(\text{Ad}(\exp(-Y))U)) = \beta(\varphi, \text{Ad}(\exp(-Y))U) = \beta(\text{Ad}(\exp Y)\varphi, U).$$

The result then follows from Proposition 4.3. □

Let  $\omega$  and  $\omega_0$  be the Kirillov 2-forms on  $O(\xi_1, \xi_2)$  and  $o(\xi_2)$ , respectively. We endow  $\bar{\mathfrak{n}}^2$  with the symplectic form  $dY \wedge dZ := \sum_{k=1}^n dy_k \wedge dz_k$ .

**Proposition 5.2.** *The map  $\Psi$  is a symplectomorphism from the symplectic product  $(\bar{\mathfrak{n}}^2 \times o(\xi_2), (dY \wedge dZ) \otimes \omega_0)$  onto  $(\bar{O}(\xi_1, \xi_2), \omega|_{\bar{O}(\xi_1, \xi_2)})$ .*

PROOF: The proof is similar to that of Proposition 6.2 in [10]. □

Now, we obtain an adapted Weyl transform on  $O(\xi_1, \xi_2)$  by transferring to  $O(\xi_1, \xi_2)$  the Berezin-Weyl calculus on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$ . We say that a smooth function  $f$  on  $O(\xi_1, \xi_2)$  is a symbol on  $O(\xi_1, \xi_2)$  (respectively a  $P$ -symbol, an  $S$ -symbol) if  $f \circ \Psi$  is a symbol (respectively a  $P$ -symbol, an  $S$ -symbol) for the Berezin-Weyl calculus on  $\bar{\mathfrak{n}}^2 \times o(\xi_2)$ .

**Proposition 5.3.** *Let  $\mathcal{A}$  be the space of  $P$ -symbols on  $O(\xi_1, \xi_2)$  and let  $\mathcal{B}$  be the space of differential operators on  $\bar{\mathfrak{n}}$  with coefficients in  $C^\infty(\bar{\mathfrak{n}}, E)$ . Then the map  $W : \mathcal{A} \rightarrow \mathcal{B}$  that assigns to each  $f \in \mathcal{A}$  the operator  $\mathcal{W}(f \circ \Psi)$  on  $L^2(\bar{\mathfrak{n}}, E)$  is an adapted Weyl correspondence in the sense of Definition 1.1.*

PROOF: Properties (1), (2) and (3) of the definition of an adapted Weyl correspondence are clearly satisfied with  $\mathcal{D} = C_0(\bar{\mathfrak{n}}, E)$ . Property (4) follows from (2) of Proposition 4.1 and from the similar result for the usual Weyl calculus, see [17]. Finally, Property (5) is an immediate consequence of Proposition 4.1. □

Finally, let us consider the case when  $G_0$  is a complex Lie group. In this case, we have  $V = i\mathfrak{k}$  and  $M$  is the maximal torus  $\exp(ia)$  of  $K$  [19, p. 143 and p. 468].

Moreover,  $o(\xi_2)$  reduces to the point  $\xi_2$ ,  $\sigma$  is a character of  $M$  and  $E = \mathbb{C}$ . So, the map  $\mathcal{W}$  is just the usual Weyl calculus.

Note that the construction of [10] can also be applied in this case. In [10], we have defined a symplectomorphism  $\Psi_0$  from  $\mathfrak{n}^2$  onto  $\tilde{O}(\xi_1, \xi_2)$  and an adapted Weyl correspondence  $W_0$  on  $O(\xi_1, \xi_2)$ . We can easily verify that  $\Psi(Y, Z) = \Psi_0(\theta(Y), \theta(Z))$  for each  $(Y, Z) \in \bar{\mathfrak{n}} \times \bar{\mathfrak{n}}$  and that the spaces of symbols for  $W$  and for  $W_0$  are the same. Moreover, choosing the orthonormal basis for  $\bar{\mathfrak{n}}$  in Section 2 and for  $\mathfrak{n}$  in [10] in compatible ways, we have that  $W_0(f)(\phi \circ \theta) = (W(f)\phi) \circ \theta$  for each  $S$ -symbol  $f$  on  $O(\xi_1, \xi_2)$  and for each  $\phi \in C_0(\bar{\mathfrak{n}})$ .

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UNIVERSITÉ DE METZ, UFR-MIM, DÉPARTEMENT DE MATHÉMATIQUES, LMMAS,  
ISGMP-BÂT. A, ILE DU SAULCY 57045, METZ CEDEX 01, FRANCE

*E-mail:* cahen@univ-metz.fr

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