

Marcella Palese

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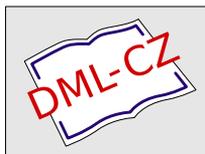
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## Erratum to: Variations by generalized symmetries of local Noether strong currents equivalent to global canonical Noether currents

Marcella Palese

**Abstract.** We correct misprints in a formula in the last sentence at the end of page 129; the first paragraph of subsection 4.1; misprints at the end of page 132 and in Proposition 1 at page 133 of the paper ‘Variations by generalized symmetries of local Noether strong currents equivalent to global canonical Noether currents’, Communications in Mathematics 24 (2016), 125–135. DOI: 10.1515/cm-2016-0009

### 1 Errata corrige

- The formulae in the last sentence at the end of page 129 should read as follows.

$$\begin{aligned}\mathcal{E}_{n+1}(\eta) &= 0 & \mathfrak{d}\eta &= 0 \\ \delta_n(\eta) \neq 0 &\implies \mathfrak{d}\lambda_i \neq 0,\end{aligned}$$

- The first paragraph of subsection 4.1 should read as follows.

Suppose that  $\mathfrak{d}\mathcal{L}_{j^r\Xi}\lambda_i = 0$ . We are here particularly concerned with the case when  $\mathcal{L}_{j^r\Xi}\lambda_i \neq 0$  i.e.  $j^r\Xi$  is a generalized symmetry, that means  $\mathcal{L}_{j^r\Xi}\lambda_i =_{\text{loc.}} d_H\beta_i$ . In fact, since  $\mathcal{E}_n(\mathcal{L}_{j^r\Xi}\lambda_i) = 0$  then  $\mathcal{L}_{j^r\Xi}\lambda_i$  defines a cohomology class and we have  $\delta_{n-1}(\mathcal{L}_{j^r\Xi}\lambda_i) \neq 0$  which implies  $\mathfrak{d}\beta_i \neq 0$ . Notice that if  $\mathfrak{d}\mathcal{L}_{j^r\Xi}\lambda_i = 0$  then  $0 = \mathfrak{d}d_H\beta_i = d_H\mathfrak{d}\beta_i$ .

- the last line at page 132 should read:

[...] always satisfied for vertical generalized symmetries of global dynamical forms.

- Proposition 1 at page 133 should read as follows.

**Proposition 1.** *Let  $j^r\Xi$  be a vertical generalized symmetry. The coboundary of the strong Noether currents is locally exact, i.e.  $d_H(\mathfrak{d}(\nu_i + \epsilon_i)) = 0$ .*

*Proof.* To get the assertion it is enough to prove that  $\mathfrak{d}\mathcal{L}_{j^r\Xi}\lambda_i = 0$ . To prove this we note that we have, identically,

$$\delta_n \mathcal{E}_n(\mathfrak{d}\mathcal{L}_{j^r\Xi}\lambda_i) = 0.$$

Indeed, by linearity

$$\delta_n \mathcal{E}_n(\mathfrak{d}\mathcal{L}_{j^r\Xi}\lambda_i) = \delta_n \mathfrak{d}\mathcal{E}_n(\mathcal{L}_{j^r\Xi}\lambda_i) = \delta_n \mathfrak{d}\mathcal{L}_{j^r\Xi}\mathcal{E}_n(\lambda_i),$$

and being  $j^r\Xi$  a generalized symmetry  $\mathcal{L}_{j^r\Xi}\mathcal{E}_n(\lambda_i) = 0$ , thus we get immediately the result, because this implies  $\mathfrak{d}(\mathfrak{d}\mathcal{L}_{j^r\Xi}(\lambda_i) - i(d_H\bar{\beta}_i)) = 0$ , i.e.  $d_H\bar{\beta}_i = \mathfrak{d}d_H\bar{\psi}_i$ ; for vertical generalized symmetries we have  $\mathfrak{d}\bar{\psi}_i = d_H\hat{\psi}_{ij}$  which gives us the results.

Note that the result holds true for any generalized symmetry (not necessarily vertical) of  $\eta_{\lambda_i}$  generating Noether conserved currents for Lagrangians of the type  $\lambda_i - d_H\mu_i$ , with  $\mu_i$  satisfying  $d_H\beta_i = d_H\mathcal{L}_{j^r\Xi}\mu_i$  and  $\mathcal{L}_{j^r\Xi}\mu_i = d_H\psi_i + \mathfrak{d}\bar{\psi}_i$ . A necessary condition for any generalized symmetry to generate Noether strong currents with closed coboundary is  $\mathcal{L}_{j^r\Xi}\mathcal{L}_{j^r\Xi}\mathfrak{d}\lambda_i = 0$ .  $\square$

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