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## CALCULUS ON SYMPLECTIC MANIFOLDS

MICHAEL EASTWOOD AND JAN SLOVÁK

ABSTRACT. On a symplectic manifold, there is a natural elliptic complex replacing the de Rham complex. It can be coupled to a vector bundle with connection and, when the curvature of this connection is constrained to be a multiple of the symplectic form, we find a new complex. In particular, on complex projective space with its Fubini–Study form and connection, we can build a series of differential complexes akin to the Bernstein–Gelfand–Gelfand complexes from parabolic differential geometry.

### 1. INTRODUCTION

Throughout this article  $M$  will be a smooth manifold of dimension  $2n$  equipped with a symplectic form  $J_{ab}$ . Here, we are using Penrose’s abstract index notation [15] and non-degeneracy of this 2-form says that there is a skew contravariant 2-form  $J^{ab}$  such that  $J_{ab}J^{ac} = \delta_b^c$  where  $\delta_b^c$  is the canonical pairing between vectors and co-vectors.

Let  $\Lambda^k$  denote the bundle of  $k$ -forms on  $M$ . The homomorphism

$$\Lambda^k \rightarrow \Lambda^{k-2} \text{ given by } \phi_{abc\dots d} \mapsto J^{ab}\phi_{abc\dots d}$$

is surjective for  $2 \leq k \leq n$  with non-trivial kernel, corresponding to the irreducible representation

$$\begin{array}{ccccccccccc} \bullet & \bullet & \dots & \bullet & \bullet & \dots & \bullet & \bullet & \bullet & \bullet & \bullet \\ & & & \uparrow & & & & \swarrow & & & \\ & & & k^{\text{th}} \text{ node} & & & & & & & \end{array} \text{ of } \text{Sp}(2n, \mathbb{R}) \subset \text{GL}(2n, \mathbb{R}).$$

Denoting this bundle by  $\Lambda_{\perp}^k$ , there is a canonical splitting of the short exact sequence

$$0 \rightarrow \Lambda_{\perp}^k \xrightarrow{\pi} \Lambda^k \rightarrow \Lambda^{k-2} \rightarrow 0$$

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and an elliptic complex [2, 9, 11, 16, 18]

$$(1) \quad \begin{array}{ccccccccccccccc} 0 & \rightarrow & \Lambda^0 & \xrightarrow{d} & \Lambda^1 & \xrightarrow{d_\perp} & \Lambda^2_\perp & \xrightarrow{d_\perp} & \Lambda^3_\perp & \xrightarrow{d_\perp} & \dots & \xrightarrow{d_\perp} & \Lambda^n_\perp \\ & & & & & & & & & & & & \downarrow d_\perp^2 \\ & & 0 & \leftarrow & \Lambda^0 & \xleftarrow{d_\perp} & \Lambda^1 & \xleftarrow{d_\perp} & \Lambda^2_\perp & \xleftarrow{d_\perp} & \Lambda^3_\perp & \xleftarrow{d_\perp} & \dots & \xleftarrow{d_\perp} & \Lambda^n_\perp \end{array}$$

where

- $d: \Lambda^0 \rightarrow \Lambda^1$  is the exterior derivative,
- for  $1 \leq k < n$ , the operator  $d_\perp: \Lambda^k_\perp \rightarrow \Lambda^{k+1}_\perp$  is the composition

$$\Lambda^k_\perp \hookrightarrow \Lambda^k \xrightarrow{d} \Lambda^{k+1} \xrightarrow{\pi} \Lambda^{k+1}_\perp,$$

a first order operator,

- $d_\perp: \Lambda^{k+1}_\perp \rightarrow \Lambda^k_\perp$  are canonically defined first order operators, which may be seen as adjoint to  $d_\perp: \Lambda^k_\perp \rightarrow \Lambda^{k+1}_\perp$ ,
- $d_\perp^2: \Lambda^n_\perp \rightarrow \Lambda^n_\perp$  is the composition

$$\Lambda^n_\perp \xrightarrow{d_\perp} \Lambda^{n-1}_\perp \xrightarrow{d_\perp} \Lambda^n_\perp,$$

a second order operator.

More explicitly, formulæ for these operators may be given as follows. Firstly, it is convenient to choose a *symplectic connection*  $\nabla_a$ , namely a torsion-free connection such that  $\nabla_a J_{bc} = 0$ , equivalently  $\nabla_a J^{bc} = 0$ . As shown in [12], for example, such connections always exist and if  $\nabla_a$  is one such, then the general symplectic connection is

$$\hat{\nabla}_a \phi_b = \nabla_a \phi_b + J^{cd} \Xi_{abc} \phi_d \quad \text{where } \Xi_{abc} = \Xi_{(abc)}.$$

Then, for  $1 \leq k < n$ , the operator  $d_\perp: \Lambda^k_\perp \rightarrow \Lambda^{k+1}_\perp$  is given by

$$(2) \quad \phi_{def\dots g} \mapsto \nabla_{[c} \phi_{def\dots g]} - \frac{k}{2(n+1-k)} J^{ab} (\nabla_a \phi_{b[ef\dots g]}) J_{cd}$$

and  $d_\perp: \Lambda^{k+1}_\perp \rightarrow \Lambda^k_\perp$  is given by

$$(3) \quad \psi_{cdef\dots g} \mapsto J^{bc} \nabla_b \psi_{cdef\dots g}.$$

Now suppose  $E$  is a smooth vector bundle on  $M$  and  $\nabla: E \rightarrow \Lambda^1 \otimes E$  is a connection. Choosing any torsion-free connection on  $\Lambda^1$  induces a connection on  $\Lambda^1 \otimes E$  and, as is well-known, the composition

$$\Lambda^1 \otimes E \rightarrow \Lambda^1 \otimes \Lambda^1 \otimes E \rightarrow \Lambda^2 \otimes E$$

does not depend on this choice. (It is the second in a well-defined sequence of differential operators

$$(4) \quad E \xrightarrow{\nabla} \Lambda^1 \otimes E \xrightarrow{\nabla} \Lambda^2 \otimes E \xrightarrow{\nabla} \dots \xrightarrow{\nabla} \Lambda^{2n-1} \otimes E \xrightarrow{\nabla} \Lambda^{2n} \otimes E$$

known as the *coupled de Rham sequence*.) In particular, we may define a homomorphism  $\Theta: E \rightarrow E$  by

$$J^{ab} \nabla_a \nabla_b \Sigma = \frac{1}{2n} \Theta \Sigma \quad \text{for } \Sigma \in \Gamma(E).$$

It is part of the curvature of  $\nabla$  and if this is the only curvature, then

$$(5) \quad (\nabla_a \nabla_b - \nabla_b \nabla_a) \Sigma = 2J_{ab} \Theta \Sigma,$$

and we shall say that  $\nabla$  is *symplectically flat*. Looking back at (1), it is easy to see that there are coupled operators

$$E \xrightarrow{\nabla} \wedge^1 \otimes E \xrightarrow{\nabla_\perp} \wedge^2_\perp \otimes E \xrightarrow{\nabla_\perp} \dots \xrightarrow{\nabla_\perp} \wedge^{n-1}_\perp \otimes E \xrightarrow{\nabla_\perp} \wedge^n_\perp \otimes E,$$

explicit formulæ for which are just as in the uncoupled cases (2) and (3). To complete the coupled version of (1) let us use

$$(6) \quad \nabla^2_\perp - \frac{2}{n}\Theta : \wedge^n_\perp \otimes E \longrightarrow \wedge^n_\perp \otimes E$$

for the middle operator. It is evident that

$$E \xrightarrow{\nabla} \wedge^1 \otimes E \xrightarrow{\nabla_\perp} \wedge^2_\perp \otimes E$$

is a complex if and only if  $\nabla$  is symplectically flat. The reason for the curvature term in (6) is that this feature propagates as follows.

**Theorem 1.** *Suppose  $E \xrightarrow{\nabla} \wedge^1 \otimes E$  is a symplectically flat connection and define  $\Theta : E \rightarrow E$  by (5). Then the coupled version of (1)*

$$\begin{array}{ccccccc} 0 & \rightarrow & E & \xrightarrow{\nabla} & \wedge^1 \otimes E & \xrightarrow{\nabla_\perp} & \wedge^2_\perp \otimes E & \xrightarrow{\nabla_\perp} & \dots & \xrightarrow{\nabla_\perp} & \wedge^n_\perp \otimes E \\ & & & & & & & & & & \downarrow \nabla^2_\perp - \frac{2}{n}\Theta \\ 0 & \leftarrow & E & \xleftarrow{\nabla_\perp} & \wedge^1 \otimes E & \xleftarrow{\nabla_\perp} & \wedge^2_\perp \otimes E & \xleftarrow{\nabla_\perp} & \dots & \xleftarrow{\nabla_\perp} & \wedge^n_\perp \otimes E \end{array}$$

is a complex. It is locally exact except near the beginning where

$$\ker \nabla : E \rightarrow \wedge^1 \otimes E \quad \text{and} \quad \frac{\ker \nabla_\perp : \wedge^1 \otimes E \rightarrow \wedge^2_\perp \otimes E}{\text{im } \nabla : E \rightarrow \wedge^1 \otimes E}$$

may be identified with the kernel and cokernel, respectively, of  $\Theta$  as locally constant sheaves.

More precision and a proof of Theorem 1 will be provided in §2. Our next theorem yields some natural symplectically flat connections.

**Theorem 2.** *Suppose  $M$  is a  $2n$ -dimensional symplectic manifold with symplectic connection  $\nabla_a$ . Then there is a natural vector bundle  $\mathcal{T}$  on  $M$  of rank  $2n + 2$  equipped with a connection, which is symplectically flat if and only if the curvature  $R_{ab}{}^c{}_d$  of  $\nabla_a$  has the form*

$$(7) \quad R_{ab}{}^c{}_d = \delta_a{}^c P_{bd} - \delta_b{}^c P_{ad} + J_{ad} P_{be} J^{ce} - J_{bd} P_{ae} J^{ce} + 2J_{ab} P_{de} J^{ce},$$

for some symmetric tensor  $P_{ab}$ .

In particular, the Fubini–Study connection on complex projective space is symplectic for the standard Kähler form and its curvature is of the form (7) for  $P_{ab} = g_{ab}$ , the standard metric. More generally, if the symplectic connection  $\nabla_a$  arises from a Kähler metric, then we shall see that (7) holds precisely in the case of constant holomorphic sectional curvature.

After proving Theorems 1 and 2, the remainder of this article is concerned with the consequences of Theorem 1 for the vector bundle  $\mathcal{T}$  and those bundles, such as  $\odot^k \mathcal{T}$ , induced from it. In particular, these consequences pertain on complex projective space where we shall find a series of elliptic complexes closely following

the Bernstein-Gelfand-Gelfand complexes on the sphere  $S^{2n+1}$  as a homogeneous space for the Lie group  $\text{Sp}(2n + 2, \mathbb{R})$ .

This article is based on our earlier work [11] but here we focus on the simpler case where we are given a symplectic structure as background. This results in fewer technicalities and in this article we include more detail, especially in constructing the BGG-like complexes in §5. Further indications justifying the shape of our complexes can be found in [3, 4, 5, 6, 7].

2. THE RUMIN-SESHADRI COMPLEX

By the *Rumin-Seshadri complex*, we mean the differential complex (1) after [16]. However, the 4-dimensional case is due to R.T. Smith [17] and the general case is also independently due to Tseng and Yau [18]. In this section we shall derive the coupled version of this complex as in Theorem 1, our proof of which includes (1) as a special case. The following lemma is also the key step in [11].

**Lemma 1.** *Suppose  $E$  is a vector bundle on  $M$  with symplectically flat connection  $\nabla : E \rightarrow \Lambda^1 \otimes E$ . Define  $\Theta : E \rightarrow E$  by (5). Then  $\Theta$  has constant rank and the bundles  $\ker \Theta$  and  $\text{coker } \Theta$  acquire from  $\nabla$ , flat connections defining locally constant sheaves  $\underline{\ker \Theta}$  and  $\underline{\text{coker } \Theta}$ , respectively. There is an elliptic complex*

$$\begin{array}{ccccccccccc}
 E & \xrightarrow{\nabla} & \Lambda^1 \otimes E & \xrightarrow{\nabla} & \Lambda^2 \otimes E & \xrightarrow{\nabla} & \Lambda^3 \otimes E & \xrightarrow{\nabla} & \Lambda^4 \otimes E & & \dots \\
 & \searrow & \oplus & \swarrow \searrow & \oplus & \swarrow \searrow & \oplus & \swarrow \searrow & \oplus & & \\
 & & E & \longrightarrow & \Lambda^1 \otimes E & \longrightarrow & \Lambda^2 \otimes E & \longrightarrow & \Lambda^3 \otimes E & & 
 \end{array}$$

where the differentials are given by

$$\Sigma \mapsto \begin{bmatrix} \nabla \Sigma \\ \Theta \Sigma \end{bmatrix} \quad \begin{bmatrix} \phi \\ \eta \end{bmatrix} \mapsto \begin{bmatrix} \nabla \phi - J \otimes \eta \\ \nabla \eta - \Theta \phi \end{bmatrix} \quad \begin{bmatrix} \omega \\ \psi \end{bmatrix} \mapsto \begin{bmatrix} \nabla \omega + J \wedge \psi \\ \nabla \psi + \Theta \omega \end{bmatrix} \quad \dots$$

It is locally exact save for the zeroth and first cohomologies, which may be identified with  $\underline{\ker \Theta}$  and  $\underline{\text{coker } \Theta}$ , respectively.

**Proof.** From (5) the Bianchi identity for  $\nabla$  reads

$$0 = \nabla_{[a} (J_{bc]} \Theta) = J_{[ab} \nabla_{c]} \Theta$$

and non-degeneracy of  $J_{ab}$  implies that  $\nabla_a \Theta = 0$ . Consequently, the homomorphism  $\Theta$  has constant rank and the following diagram with exact rows commutes

$$\begin{array}{ccccccccc}
 0 & \rightarrow & \ker \Theta & \rightarrow & E & \xrightarrow{\Theta} & E & \rightarrow & \text{coker } \Theta & \rightarrow & 0 \\
 & & & & \downarrow \nabla & & \downarrow \nabla & & & & \\
 0 & \rightarrow & \Lambda^1 \otimes \ker \Theta & \rightarrow & \Lambda^1 \otimes E & \xrightarrow{\Theta} & \Lambda^1 \otimes E & \rightarrow & \Lambda^1 \otimes \text{coker } \Theta & \rightarrow & 0
 \end{array}$$

and yields the desired connections on  $\ker \Theta$  and  $\text{coker } \Theta$ , which are easily seen to be flat. Ellipticity of the given complex is readily verified and, by definition, the kernel of its first differential is  $\underline{\ker \Theta}$ . To identify the higher local cohomology of this complex the key observation is that locally we may choose a 1-form  $\tau$  such that  $d\tau = J$  and, having done this, the connection

$$\Gamma(E) \ni \Sigma \xrightarrow{\nabla} \nabla \Sigma - \tau \otimes \Theta \Sigma \in \Gamma(\Lambda^1 \otimes E)$$

is flat. The rest of the proof is diagram chasing, using exactness of

$$E \xrightarrow{\tilde{\nabla}} \wedge^1 \otimes E \xrightarrow{\tilde{\nabla}} \wedge^2 \otimes E \xrightarrow{\tilde{\nabla}} \wedge^3 \otimes E \xrightarrow{\tilde{\nabla}} \wedge^4 \otimes E \xrightarrow{\tilde{\nabla}} \dots$$

If needed, the details are in [11]. □

**Proof of Theorem 1.** In [11], the corresponding result [11, Theorem 4] is proved by invoking a spectral sequence. Here, we shall, instead, prove two typical cases ‘by hand,’ leaving the rest of the proof to the reader.

For our first case, let us suppose  $n \geq 3$  and prove local exactness of

$$\wedge^1 \otimes E \xrightarrow{\nabla_{\perp}} \wedge^2_{\perp} \otimes E \xrightarrow{\nabla_{\perp}} \wedge^3_{\perp} \otimes E.$$

Thus, we are required to show that if  $\omega_{ab}$  has values in  $E$  and

$$\omega_{ab} = \omega_{[ab]} \quad J^{ab}\omega_{ab} = 0 \quad \nabla_{[c}\omega_{de]} = \frac{1}{n-1}J^{ab}(\nabla_a\omega_{b[c}J_{de]}),$$

then locally there is  $\phi_a \in \Gamma(\wedge^1 \otimes E)$  such that

$$\omega_{cd} = \nabla_{[c}\phi_{d]} - \frac{1}{2n}J^{ab}(\nabla_a\phi_b)J_{cd}.$$

If we set  $\psi_c \equiv -\frac{1}{n-1}J^{ab}\nabla_a\omega_{bc}$ , then  $\nabla_{[c}\omega_{de]} + J_{[cd}\psi_{e]} = 0$  so

$$0 = \nabla_{[b}\nabla_c\omega_{de]} + J_{[bc}\nabla_d\psi_{e]} = J_{[bc}\Theta\omega_{de]} + J_{[bc}\nabla_d\psi_{e]}$$

and since  $J \wedge_{\perp} : \wedge^2 \rightarrow \wedge^4$  is injective it follows that

$$\nabla_{[c}\psi_{d]} + \Theta\omega_{cd} = 0.$$

In other words, we have shown that

$$\begin{aligned} \nabla\omega + J \wedge \psi &= 0 \\ \nabla\psi + \Theta\omega &= 0 \end{aligned}$$

and Lemma 1 locally yields  $\phi_a \in \Gamma(\wedge^1 \otimes E)$  and  $\eta \in \Gamma(E)$  such that

$$\begin{aligned} \nabla_{[a}\phi_{b]} - J_{ab}\eta &= \omega_{ab}, \\ \nabla_a\eta - \Theta\phi_a &= \psi_a. \end{aligned}$$

In particular,

$$J^{ab}\nabla_a\phi_b - 2n\eta = J^{ab}(\nabla_a\phi_b - J_{ab}\eta) = J^{ab}\omega_{ab} = 0$$

and, therefore,

$$\nabla_{[c}\phi_{d]} - \frac{1}{2n}J^{ab}(\nabla_a\phi_b)J_{cd} = \nabla_{[c}\phi_{d]} - \eta J_{cd} = \omega_{cd},$$

as required.

Our second case is more involved. It is to show that

$$(8) \quad \wedge^k_{\perp} \otimes E \xrightarrow{\nabla^2_{\perp} - \frac{2}{n}\Theta} \wedge^k_{\perp} \otimes E \xrightarrow{\nabla_{\perp}} \wedge^{k-1}_{\perp} \otimes E$$

is locally exact. As regards  $\nabla_{\perp} : \wedge^k_{\perp} \otimes E \rightarrow \wedge^{k-1}_{\perp} \otimes E$ , notice that

$$J^{bc}\nabla_b\psi_{cdef\dots g} = \frac{n+1}{2}J^{bc}\nabla_{[b}\psi_{cdef\dots g]}$$

and that if  $\phi_{def\dots g} \in \Gamma(\wedge^k \otimes E)$ , then

$$(9) \quad J^{bc}J_{[bc}\phi_{def\dots g]} = \frac{4(n-k)}{(k+1)(k+2)}\phi_{def\dots g} + \frac{k(k-1)}{(k+1)(k+2)}J_{[de}\phi_{f\dots g]bc}J^{bc}$$

so if  $\phi_{def\dots g} \in \Gamma(\wedge_{\perp}^{n-1} \otimes E)$ , then

$$J^{bc} J_{[bc} \phi_{def\dots g]} = \frac{4}{n(n+1)} \phi_{def\dots g}.$$

Therefore,  $\nabla_{\perp} \psi \in \Gamma(\wedge_{\perp}^{n-1} \otimes E)$  is characterised by

$$(10) \quad J \wedge \nabla_{\perp} \psi = \frac{2}{n} \nabla \psi$$

as an equation in  $\wedge^{n+1} \otimes E$ . In particular, in  $\wedge^{n+2} \otimes E$  we find

$$J \wedge \nabla \nabla_{\perp} \psi = \nabla (J \wedge \nabla_{\perp} \psi) = \frac{2}{n} \nabla^2 \psi = J \wedge \Theta \psi = 0$$

whence  $\nabla \nabla_{\perp} \psi$  already lies in  $\wedge^n \otimes E$  and there is no need to remove the trace as in (2) to form  $\nabla_{\perp}^2 \psi$ . Therefore, invoking (10) once again, the composition

$$\wedge_{\perp}^n \otimes E \xrightarrow{\nabla_{\perp}} \wedge_{\perp}^{n-1} \otimes E \xrightarrow{\nabla_{\perp}} \wedge_{\perp}^n \otimes E \xrightarrow{\nabla_{\perp}} \wedge_{\perp}^{n-1} \otimes E$$

is characterised by

$$J \wedge \nabla_{\perp}^3 \psi = \frac{2}{n} \nabla \nabla_{\perp}^2 \psi = \frac{2}{n} \nabla^2 \nabla_{\perp} \psi = \frac{2}{n} J \wedge \Theta \nabla_{\perp} \psi = \frac{2}{n} J \wedge \nabla_{\perp} \Theta \psi$$

and, since  $J \wedge \_ : \wedge^{n-1} \rightarrow \wedge^{n+1}$  is an isomorphism, we conclude that  $\nabla_{\perp}^3 \psi = \frac{2}{n} \nabla_{\perp} \Theta \psi$ , equivalently that (8) is a complex.

Before proceeding, let us remark on another consequence of (9), namely that for  $\nu_{cdef\dots g} \in \Gamma(\wedge^n \otimes E)$ ,

$$(11) \quad J_{[ab} \nu_{cdef\dots g]} = 0 \iff J^{cd} \nu_{cdef\dots g} = 0.$$

Now to establish local exactness, suppose  $\nu \in \Gamma(\wedge_{\perp}^n \otimes E)$  satisfies  $\nabla_{\perp} \nu = 0$ . Equivalently, according to (10) and (11)

$$\nu \in \Gamma(\wedge^n \otimes E) \quad \text{satisfies} \quad \nabla \nu = 0 \quad \text{and} \quad J \wedge \nu = 0.$$

Lemma 1 implies that locally there are

$$\begin{aligned} \phi \in \Gamma(\wedge^n \otimes E) \\ \eta \in \Gamma(\wedge^{n-1} \otimes E) \end{aligned} \quad \text{such that} \quad \begin{aligned} \nabla \phi - J \wedge \eta &= 0 \\ \nabla \eta - \Theta \phi &= \nu. \end{aligned}$$

Since

$$0 \rightarrow \wedge^{n-2} \xrightarrow{J \wedge \_} \wedge^n \rightarrow \wedge_{\perp}^n \rightarrow 0$$

is exact, we can write  $\phi$  uniquely as

$$\phi = \psi + J \wedge \tau,$$

where  $\psi \in \Gamma(\wedge_{\perp}^n \otimes E)$  and  $\tau \in \Gamma(\wedge^{n-2} \otimes E)$ . We conclude that

$$\begin{aligned} \nabla \psi - J \wedge \hat{\eta} &= 0 \\ \nabla \hat{\eta} - \Theta \psi &= \nu, \end{aligned} \quad (\text{where } \hat{\eta} = \eta - \nabla \tau).$$

However, as discussed above, these equations say exactly that

$$\nabla_{\perp}^2 \psi - \frac{2}{n} \Theta \psi = \nu,$$

and exactness is shown. □

3. TRACTOR BUNDLES

For the rest of the article we suppose that we are given, not only a manifold  $M$  with symplectic form  $J_{ab}$ , but also a torsion-free connection  $\nabla_a$  on the tangent bundle (and hence on all other tensor bundles) such that  $\nabla_a J_{bc} = 0$ . This is sometimes called a *Fedosov structure* [12] on  $M$ . The curvature  $R_{ab}{}^c{}_d$  of  $\nabla_a$ , characterised by

$$(\nabla_a \nabla_b - \nabla_b \nabla_a)X^c = R_{ab}{}^c{}_d X^d,$$

satisfies

$$R_{ab}{}^c{}_d = R_{[ab]}{}^c{}_d \quad R_{[ab}{}^c{}_d] = 0 \quad R_{ab}{}^c{}_d J_{ce} = R_{ab}{}^e{}_d J_{cd}$$

and enjoys the following decomposition into irreducible parts

$$R_{ab}{}^c{}_d = V_{ab}{}^c{}_d + \delta_a{}^c P_{bd} - \delta_b{}^c P_{ad} + J_{ad} P_{be} J^{ce} - J_{bd} P_{ae} J^{ce} + 2J_{ab} P_{de} J^{ce},$$

for some symmetric  $P_{ab}$ , where  $V_{ab}{}^a{}_d = 0$  (reflecting the branching

$$\square\square = \square\square_{\perp} \oplus \square\square$$

of representations under  $GL(2n, \mathbb{R}) \supset Sp(2n, \mathbb{R})$ ). Notice that

$$(12) \quad P_{bd} = \frac{1}{2(n+1)} R_{ab}{}^a{}_d = \frac{1}{4(n+1)} J^{ae} R_{ae}{}^c{}_b J_{cd}.$$

We define the *standard tractor bundle* to be the rank  $2n + 2$  vector bundle  $\mathcal{T} \equiv \wedge^0 \oplus \wedge^1 \oplus \wedge^0$  with its *tractor connection*

$$\nabla_a \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix} = \begin{bmatrix} \nabla_a \sigma - \mu_a \\ \nabla_a \mu_b + J_{ab} \rho + P_{ab} \sigma \\ \nabla_a \rho - P_{ab} J^{bc} \mu_c + S_a \sigma \end{bmatrix}, \text{ where } S_a \equiv \frac{1}{2n+1} J^{bc} \nabla_c P_{ab}.$$

Readers familiar with conformal differential geometry may recognise the form of this connection as following the tractor connection in that setting [1]. If needs be, we shall write *symplectic tractor connection* to distinguish the connection just defined from any alternatives. We shall need the following curvature identities.

**Lemma 2.** *Let  $Y_{abc} \equiv \frac{1}{2n+1} \nabla_c V_{ab}{}^c{}_d$ . Then*

$$(13) \quad Y_{abc} = 2\nabla_{[a} P_{b]c} - 2J_{c[a} S_{b]} + 2J_{ab} S_c$$

and

$$(14) \quad \begin{aligned} J^{ad} \nabla_a Y_{bcd} &= J^{ad} V_{bc}{}^e{}_a P_{ed} + 4n(J^{ad} P_{ba} P_{cd} - \nabla_{[b} S_{c]}) \\ &+ 2J_{bc} J^{ad} (\nabla_a S_d - J^{ef} P_{ae} P_{df}). \end{aligned}$$

**Proof.** Writing the Bianchi identity  $\nabla_{[a} R_{bc]d}{}^e = 0$  in terms of  $V_{ab}{}^c{}_d$  and  $P_{ab}$  yields

$$\nabla_{[a} V_{bc]d}{}^e = -2\delta_{[b}{}^d \nabla_a P_{c]e} + 2J^{df} J_{e[b} \nabla_a P_{c]f} - 2J^{df} J_{[bc} \nabla_a] P_{ef}.$$

and contracting over  $a^d$  gives

$$\begin{aligned} \frac{1}{3} \nabla_a V_{bc}{}^a{}_e &= \frac{4(n-1)}{3} \nabla_{[b} P_{c]e} + \frac{2}{3} [\nabla_{[b} P_{c]e} - (2n+1) J_{e[b} S_{c]}] \\ &+ \frac{2}{3} [(2n+1) J_{bc} S_e + 2\nabla_{[b} P_{c]e}], \end{aligned}$$

which is easily rearranged as (13). For (14), firstly notice that

$$J^{ad}R_{ab}{}^e{}_d = J^{ed}R_{ab}{}^a{}_d = 2(n+1)J^{ed}P_{bd}$$

and the Bianchi symmetry may be written as  $R_{a[b}{}^e{}_c] = -\frac{1}{2}R_{bc}{}^e{}_a$ . Thus,

$$\begin{aligned} J^{ad}\nabla_a\nabla_bP_{cd} &= \nabla_bJ^{ad}\nabla_aP_{cd} - J^{ad}R_{ab}{}^e{}_cP_{ed} - J^{ad}R_{ab}{}^e{}_dP_{ce} \\ &= -(2n+1)\nabla_bS_c - J^{ad}R_{ab}{}^e{}_cP_{ed} + 2(n+1)J^{de}P_{bd}P_{ce} \end{aligned}$$

and so

$$J^{ad}\nabla_a\nabla_{[b}P_{c]d} = -(2n+1)\nabla_{[b}S_{c]} + \frac{1}{2}J^{ad}R_{bc}{}^e{}_aP_{ed} + 2(n+1)J^{de}P_{bd}P_{ce}.$$

From (13) we see that

$$J^{ad}\nabla_aY_{bcd} = 2J^{ad}\nabla_a\nabla_{[b}P_{c]d} + 2\nabla_{[b}S_{c]} + 2J_{bc}J^{ad}\nabla_aS_d.$$

Therefore,

$$J^{ad}\nabla_aY_{bcd} = J^{ad}R_{bc}{}^e{}_aP_{ed} - 4n\nabla_{[b}S_{c]} + 4(n+1)J^{de}P_{bd}P_{ce} + 2J_{bc}J^{ad}\nabla_aS_d.$$

Finally,

$$J^{ad}R_{bc}{}^e{}_aP_{ed} = J^{ad}V_{bc}{}^e{}_aP_{ed} - 4J^{ad}P_{ba}P_{cd} - 2J_{bc}J^{ad}J^{ef}P_{ae}P_{df},$$

so

$$\begin{aligned} J^{ad}\nabla_aY_{bcd} &= J^{ad}V_{bc}{}^e{}_aP_{ed} + 4nJ^{ad}P_{ba}P_{cd} - 2J_{bc}J^{ad}J^{ef}P_{ae}P_{df} \\ &\quad - 4n\nabla_{[b}S_{c]} + 2J_{bc}J^{ad}\nabla_aS_d, \end{aligned}$$

which may be rearranged as (14). □

**Proposition 1.** *The tractor connection  $\mathcal{T} \rightarrow \Lambda^1 \otimes \mathcal{T}$  preserves the non-degenerate skew form*

$$\left\langle \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix}, \begin{bmatrix} \tilde{\sigma} \\ \tilde{\mu}_c \\ \tilde{\rho} \end{bmatrix} \right\rangle \equiv \sigma\tilde{\rho} + J^{bc}\mu_b\tilde{\mu}_c - \rho\tilde{\sigma}$$

and its curvature is given by

$$\begin{aligned} (\nabla_a\nabla_a - \nabla_b\nabla_a) \begin{bmatrix} \sigma \\ \mu_d \\ \rho \end{bmatrix} &= \begin{bmatrix} 0 \\ -V_{ab}{}^c{}_d\mu_c + Y_{abd}\sigma \\ -Y_{abc}J^{cd}\mu_d + \frac{1}{2n}(J^{cd}V_{ab}{}^e{}_cP_{de} - J^{cd}\nabla_cY_{abd})\sigma \end{bmatrix} \\ &\quad + 2J_{ab} \begin{bmatrix} \rho \\ J^{ce}P_{cd}\mu_e - S_d\sigma \\ S_cJ^{cd}\mu_d + \frac{1}{2n}J^{cd}(\nabla_cS_d - J^{ef}P_{ce}P_{df})\sigma \end{bmatrix}. \end{aligned}$$

**Proof.** We expand

$$\left\langle \nabla_a \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix}, \begin{bmatrix} \tilde{\sigma} \\ \tilde{\mu}_c \\ \tilde{\rho} \end{bmatrix} \right\rangle + \left\langle \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix}, \nabla_a \begin{bmatrix} \tilde{\sigma} \\ \tilde{\mu}_c \\ \tilde{\rho} \end{bmatrix} \right\rangle$$

to obtain

$$\begin{aligned} & (\nabla_a \sigma - \mu_a) \tilde{\rho} + \sigma (\nabla \tilde{\rho} - P_{ab} J^{bc} \tilde{\mu}_c + S_a \tilde{\sigma}) \\ & \quad + J^{bc} (\nabla_a \mu_b + J_{ab} \rho + P_{ab} \sigma) \tilde{\mu}_c + J^{bc} \mu_b (\nabla_a \tilde{\mu}_c + J_{ac} \tilde{\rho} + P_{ac} \tilde{\sigma}) \\ & \quad - (\nabla_a \rho - P_{ab} J^{bc} \mu_c + S_a \sigma) \tilde{\sigma} - \rho (\nabla_a \tilde{\sigma} - \tilde{\mu}_a) \end{aligned}$$

in which all terms cancel save for

$$(\nabla_a \sigma) \tilde{\rho} + \sigma \nabla \tilde{\rho} + J^{bc} (\nabla_a \mu_b) \tilde{\mu}_c + J^{bc} \mu_b \nabla_a \tilde{\mu}_c - (\nabla_a \rho) \tilde{\sigma} - \rho \nabla_a \tilde{\sigma},$$

which reduces to

$$\nabla_a (\sigma \tilde{\rho} + J^{bc} \mu_b \tilde{\mu}_c - \rho \tilde{\sigma}),$$

as required. For the curvature, we readily compute

$$\nabla_{[a} \nabla_{b]} \begin{bmatrix} \sigma \\ \mu_d \\ \rho \end{bmatrix} = \begin{bmatrix} \nabla_{[a} \nabla_{b]} \sigma - J_{ba} \rho \\ \nabla_{[a} \nabla_{b]} \mu_d + J_{d[a} P_{b]c} J^{ce} \mu_e - P_{d[a} \mu_{b]} + T_{abd} \sigma \\ \nabla_{[a} \nabla_{b]} \rho - T_{abc} J^{cd} \mu_d + (\nabla_{[a} S_{b]} - J^{cd} P_{ac} P_{bd}) \sigma \end{bmatrix},$$

where  $T_{abc} \equiv \nabla_{[a} P_{b]c} - J_{c[a} S_{b]}$ . Lemma 2, however, states that

$$T_{abc} = \frac{1}{2} Y_{abc} - J_{ab} S_c$$

and

$$\begin{aligned} 4n(\nabla_{[a} S_{b]} - J^{cd} P_{ac} P_{bd}) &= J^{cd} V_{ab}{}^e{}_c P_{de} - J^{cd} \nabla_c Y_{abd} \\ &\quad + 2J_{ab} J^{cd} (\nabla_c S_d - J^{ef} P_{ce} P_{df}). \end{aligned}$$

Therefore,

$$\begin{aligned} \nabla_{[a} \nabla_{b]} \begin{bmatrix} \sigma \\ \mu_d \\ \rho \end{bmatrix} &= \begin{bmatrix} 0 \\ \nabla_{[a} \nabla_{b]} \mu_d + J_{d[a} P_{b]c} J^{ce} \mu_e - P_{d[a} \mu_{b]} + \frac{1}{2} Y_{abd} \sigma \\ -\frac{1}{2} Y_{abc} J^{cd} \mu_d + \frac{1}{4n} (J^{cd} V_{ab}{}^e{}_c P_{de} - J^{cd} \nabla_c Y_{abd}) \sigma \end{bmatrix} \\ &\quad + J_{ab} \begin{bmatrix} \rho \\ -S_d \sigma \\ S_c J^{cd} \mu_d + \frac{1}{2n} J^{cd} (\nabla_c S_d - J^{ef} P_{ce} P_{df}) \sigma \end{bmatrix}. \end{aligned}$$

Finally,

$$R_{ab}{}^c{}_d \mu_c = V_{ab}{}^c{}_d \mu_c - 2P_{d[a} \mu_{b]} + 2J_{d[a} P_{b]c} J^{ce} \mu_e + 2J_{ab} P_{de} J^{ce} \mu_c,$$

so

$$\nabla_{[a} \nabla_{b]} \mu_d + J_{d[a} P_{b]c} J^{ce} \mu_e - P_{d[a} \mu_{b]} = -\frac{1}{2} V_{ab}{}^c{}_d \mu_c - J_{ab} P_{de} J^{ce} \mu_c$$

whence

$$\begin{aligned} \nabla_{[a} \nabla_{b]} \begin{bmatrix} \sigma \\ \mu_d \\ \rho \end{bmatrix} &= \begin{bmatrix} 0 \\ -\frac{1}{2} V_{ab}{}^c{}_d \mu_c + \frac{1}{2} Y_{abd} \sigma \\ -\frac{1}{2} Y_{abc} J^{cd} \mu_d + \frac{1}{4n} (J^{cd} V_{ab}{}^e{}_c P_{de} - J^{cd} \nabla_c Y_{abd}) \sigma \end{bmatrix} \\ &\quad + J_{ab} \begin{bmatrix} \rho \\ J^{ce} P_{cd} \mu_e - S_d \sigma \\ S_c J^{cd} \mu_d + \frac{1}{2n} J^{cd} (\nabla_c S_d - J^{ef} P_{ce} P_{df}) \sigma \end{bmatrix}, \end{aligned}$$

as required.  $\square$

**Corollary 1.** *The tractor connection is symplectically flat if and only if the curvature tensor  $V_{ab}{}^c{}_d$  vanishes.*

4. KÄHLER GEOMETRY

Kähler manifolds provide a familiar source of symplectic manifolds equipped with a compatible torsion-free connection as in §3. In this case, the connection  $\nabla_a$  is the Levi-Civita connection of a metric  $g_{ab}$  and  $J_a{}^b \equiv J_{ac}g^{bc}$  is an almost complex structure on  $M$  whose integrability is equivalent to the vanishing of  $\nabla_a J_{bc}$ . In Kähler geometry, the Riemann curvature tensor decomposes into three irreducible parts:

$$(15) \quad \begin{aligned} R_{ab}{}^c{}_d &= U_{ab}{}^c{}_d \\ &+ \delta_a{}^c \Xi_{bd} - \delta_b{}^c \Xi_{ad} - g_{ad} \Xi_b{}^c + g_{bd} \Xi_a{}^c \\ &+ J_a{}^c \Sigma_{bd} - J_b{}^c \Sigma_{ad} - J_{ad} \Sigma_b{}^c + J_{bd} \Sigma_a{}^c + 2J_{ab} \Sigma_c{}^d + 2J^c{}_d \Sigma_{ab} \\ &+ \Lambda(\delta_a{}^c g_{bd} - \delta_b{}^c g_{ad} + J_a{}^c J_{bd} - J_b{}^c J_{ad} + 2J_{ab} J^c{}_d), \end{aligned}$$

where indices have been raised using  $g^{ab}$  and

- $U_{ab}{}^c{}_d$  is totally trace-free with respect to  $g^{ab}$ ,  $J_a{}^b$ , and  $J^{ab}$ ,
- $\Xi_{ab}$  is trace-free symmetric whilst  $\Sigma_{ab} \equiv J_a{}^c \Xi_{bc}$  is skew.

Computing the Ricci curvature from this decomposition, we find

$$R_{bd} \equiv R_{ab}{}^a{}_d = 2(n + 2)\Xi_{bd} + 2(n + 1)\Lambda g_{bd}$$

and therefore from (12) conclude that

$$P_{ab} = \frac{n+2}{n+1}\Xi_{ab} + \Lambda g_{ab}.$$

Hence

$$\begin{aligned} J_c{}^a R_{ab}{}^c{}_d &= J_c{}^a V_{ab}{}^c{}_d - J_{bd} P_a{}^a - 2J_b{}^a P_{da} \\ &= J_c{}^a V_{ab}{}^c{}_d - 2\frac{n+2}{n+1}\Sigma_{bd} - 2(n+1)\Lambda J_{bd}. \end{aligned}$$

On the other hand, from (15) we find

$$J_c{}^a R_{ab}{}^c{}_d = -2(n+2)\Sigma_{bd} - 2(n+1)\Lambda J_{bd}$$

and, comparing these two expressions gives

$$J_c{}^a V_{ab}{}^c{}_d - 2\frac{n+2}{n+1}\Sigma_{bd} = -2(n+2)\Sigma_{bd}$$

and we have established the following.

**Proposition 2.** *Concerning the symplectic curvature decomposition on a Kähler manifold,*

$$J_c{}^a V_{ab}{}^c{}_d = -2\frac{n(n+2)}{n+1}\Sigma_{bd}.$$

**Corollary 2.** *The symplectic tractor connection on a Kähler manifold is symplectically flat if and only if the metric has constant holomorphic sectional curvature.*

**Proof.** According to Corollary 1, we have to interpret the constraint  $V_{ab}{}^c{}_d = 0$  in the Kähler case. From (15) it is already clear that  $U_{ab}{}^c{}_d = 0$  and Proposition 2 implies that also  $\Sigma_{ab} = 0$  so (15) reduces to

$$R_{ab}{}^c{}_d = \Lambda(\delta_a{}^c g_{bd} - \delta_b{}^c g_{ad} + J_a{}^c J_{bd} - J_b{}^c J_{ad} + 2J_{ab} J^c{}_d),$$

which is exactly the constancy of holomorphic sectional curvature. □

### 5. BGG-LIKE COMPLEXES ON $\mathbb{C}\mathbb{P}_n$

Fix a real vector space  $\mathfrak{g}_{-1}$  of dimension  $2n$ , let  $\mathfrak{g}_1$  denotes its dual, and fix a non-degenerate 2-form  $J_{ab} \in \Lambda^2 \mathfrak{g}_1$ . The  $(2n + 1)$ -dimensional Heisenberg Lie algebra may be realised as

$$\mathfrak{h} = \mathbb{R} \oplus \mathfrak{g}_{-1},$$

where the first summand is the 1-dimensional centre of  $\mathfrak{h}$  and the Lie bracket on  $\mathfrak{g}_{-1}$  is given by

$$[X, Y] = 2J_{ab} X^a Y^b \in \mathbb{R} \hookrightarrow \mathfrak{h}.$$

We should admit right away that the reason for this seemingly arcane notation is that we shall soon have occasion to write

$$(16) \quad \mathfrak{sp}(2n + 2, \mathbb{R}) = \begin{array}{ccccccc} \mathfrak{g}_{-2} & \oplus & \mathfrak{g}_{-1} & \oplus & \mathfrak{g}_0 & \oplus & \mathfrak{g}_1 & \oplus & \mathfrak{g}_2 \\ & & \parallel & & \parallel & & \parallel & & \\ & & \mathbb{R} & & \mathfrak{sp}(2n, \mathbb{R}) \oplus \mathbb{R} & & \mathbb{R} & & \end{array}$$

(a  $|2|$ -graded Lie algebra as in [8, §4.2.6]) and, in particular, regard  $\mathfrak{h} = \mathbb{R} \oplus \mathfrak{g}_{-1} = \mathfrak{g}_{-2} \oplus \mathfrak{g}_{-1}$  as a Lie subalgebra of  $\mathfrak{sp}(2n + 2, \mathbb{R})$ . Be that as it may, let us suppose that  $\mathbb{V}$  is a finite-dimensional representation of  $\mathfrak{h}$ . The Lie algebra cohomology  $H^r(\mathfrak{h}, \mathbb{V})$  may be realised as the cohomology of the Chevalley-Eilenberg complex

$$(17) \quad 0 \rightarrow \mathbb{V} \rightarrow \mathfrak{h}^* \otimes \mathbb{V} \rightarrow \dots \rightarrow \wedge^r \mathfrak{h}^* \otimes \mathbb{V} \rightarrow \wedge^{r+1} \mathfrak{h}^* \otimes \mathbb{V} \rightarrow \dots$$

as, for example, in [13, Chapter IV]. We shall require, however, the following alternative realisation.

**Lemma 3.** *There is a complex*

$$(18) \quad \begin{array}{ccccccccccc} 0 & \rightarrow & \mathbb{V} & \xrightarrow{\partial} & \mathfrak{g}_1 \otimes \mathbb{V} & \xrightarrow{\partial_\perp} & \wedge^2_\perp \mathfrak{g}_1 \otimes \mathbb{V} & \xrightarrow{\partial_\perp} & \dots & \xrightarrow{\partial_\perp} & \wedge^n_\perp \mathfrak{g}_1 \otimes \mathbb{V} \\ & & & & & & & & & & \downarrow \\ 0 & \leftarrow & \mathbb{V} & \xleftarrow{\partial_\perp} & \mathfrak{g}_1 \otimes \mathbb{V} & \xleftarrow{\partial_\perp} & \wedge^2_\perp \mathfrak{g}_1 \otimes \mathbb{V} & \xleftarrow{\partial_\perp} & \dots & \xleftarrow{\partial_\perp} & \wedge^n_\perp \mathfrak{g}_1 \otimes \mathbb{V} \end{array}$$

whose cohomology realises  $H^r(\mathfrak{h}, \mathbb{V})$ . Here, we are writing

$$\wedge^r_\perp \mathfrak{g}_1 \equiv \{\omega_{abc\dots d} \in \wedge^r \mathfrak{g}_1 \mid J^{ab} \omega_{abc\dots d} = 0\},$$

where  $J^{ab} \in \Lambda^2 \mathfrak{g}_{-1}$  is the inverse of  $J_{ab} \in \Lambda^2 \mathfrak{g}_1$  (let's say normalised so that  $J_{ab} J^{ac} = \delta_b{}^c$ ).

**Proof.** Notice that any representation  $\rho : \mathfrak{h} \rightarrow \text{End}(\mathbb{V})$  is determined by its restriction to  $\mathfrak{g}_{-1} \subset \mathfrak{h}$ . Indeed, writing  $\partial_a : \mathfrak{g}_{-1} \rightarrow \text{End}(\mathbb{V})$  for this restriction, to say that  $\rho$  is a representation of  $\mathfrak{h}$  is to say that

$$(19) \quad \left. \begin{array}{l} (\partial_a \partial_b - \partial_b \partial_a)v = 2J_{ab} \theta v \\ (\partial_a \theta - \theta \partial_a)v = 0 \end{array} \right\} \quad \forall v \in \mathbb{V},$$

where  $\theta \in \text{End}(\mathbb{V})$  is  $\rho(1)$  for  $1 \in \mathbb{R} \subset \mathfrak{h}$ .

The splitting  $\mathfrak{h}^* = \mathfrak{g}_1 \oplus \mathbb{R}$  allows us to write (17) as

$$(20) \quad \begin{array}{ccccccc} \mathbb{V} & \longrightarrow & \mathfrak{h}^* \otimes \mathbb{V} & \longrightarrow & \wedge^2 \mathfrak{h}^* \otimes \mathbb{V} & \longrightarrow & \wedge^3 \mathfrak{h}^* \otimes \mathbb{V} \longrightarrow \dots \\ \parallel & & \parallel & & \parallel & & \parallel \\ \mathbb{V} & \longrightarrow & \mathfrak{g}_1 \otimes \mathbb{V} & \longrightarrow & \wedge^2 \mathfrak{g}_1 \otimes \mathbb{V} & \longrightarrow & \wedge^3 \mathfrak{g}_1 \otimes \mathbb{V} \longrightarrow \dots, \\ & \searrow & \oplus & \swarrow \searrow & \oplus & \swarrow \searrow & \oplus & \swarrow \searrow \\ & & \mathbb{V} & \longrightarrow & \mathfrak{g}_1 \otimes \mathbb{V} & \longrightarrow & \wedge^2 \mathfrak{g}_1 \otimes \mathbb{V} & \longrightarrow \dots \end{array}$$

where the differentials are given by

$$v \mapsto \begin{bmatrix} \partial_a v \\ \theta v \end{bmatrix} \quad \begin{bmatrix} \phi_a \\ \eta \end{bmatrix} \mapsto \begin{bmatrix} \partial_{[a} \phi_{b]} - J_{ab} \eta \\ \partial_a \eta - \theta \phi_a \end{bmatrix} \quad \begin{bmatrix} \omega_{ab} \\ \psi_a \end{bmatrix} \mapsto \begin{bmatrix} \partial_{[a} \omega_{bc]} + J_{[ab} \psi_{c]} \\ \partial_{[a} \psi_{b]} + \theta \omega_{ab} \end{bmatrix}$$

et cetera. In particular, notice that the homomorphisms

$$(21) \quad \wedge^{r-1} \mathfrak{g}_1 \ni \psi \mapsto \pm J \wedge \psi \in \wedge^{r+1} \mathfrak{g}_1$$

are

- independent of the representation on  $\mathbb{V}$ ,
- injective for  $1 \leq r < n$ ,
- an isomorphism for  $r = n$ ,
- surjective for  $n < r \leq 2n - 1$ .

Note that  $\wedge_{\perp}^{r+1} \mathfrak{g}_1$  is complementary to the image of (21) for  $1 \leq r < n$ . Also note the isomorphisms

$$\wedge^{2n+1-r} \mathfrak{g}_1 \xrightarrow{J \wedge J \wedge \dots \wedge J} \wedge^{r-1} \mathfrak{g}_1, \quad \text{for } n < r \leq 2n + 1,$$

under which the kernel of (21) may be identified with

$$\wedge_{\perp}^{2n+1-r} \mathfrak{g}_1, \quad \text{for } n < r \leq 2n - 1.$$

Diagram chasing in (20) (or the spectral sequence of a filtered complex) finishes the proof. □

**Remark.** Evidently, the equations (19) are algebraic versions of

$$\left. \begin{aligned} (\nabla_a \nabla_b - \nabla_b \nabla_a) \Sigma &= 2J_{ab} \Theta \Sigma \\ (\nabla_a \Theta - \Theta \nabla_a) \Sigma &= 0 \end{aligned} \right\} \quad \forall \Sigma \in \Gamma(E),$$

which hold for a symplectically flat connection  $\nabla_a$  on smooth vector bundle  $E$  on  $M$ . Also (20) is the evident algebraic counterpart to the differential complex of Lemma 1. It follows that explicit formulæ for the operators  $\partial_{\perp}$  in the complex (18) follow the differential versions (2) and (3) with  $\wedge_{\perp}^n \mathfrak{g} \otimes \mathbb{V} \rightarrow \wedge_{\perp}^n \mathfrak{g} \otimes \mathbb{V}$  being given by  $\partial_{\perp}^2 - \frac{2}{n} \theta$ .

Let us now consider the tractor connection on  $\mathbb{C}\mathbb{P}_n$ . According to Theorem 2, the remarks following its statement, and the discussions in §3, this is the connection on  $\mathcal{T} = \wedge^0 \oplus \wedge^1 \oplus \wedge^0$  given by

$$\nabla_a \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix} = \begin{bmatrix} \nabla_a \sigma - \mu_a \\ \nabla_a \mu_b + J_{ab} \rho + g_{ab} \sigma \\ \nabla_a \rho - J_a{}^b \mu_b \end{bmatrix} = \begin{bmatrix} \nabla_a \sigma \\ \nabla_a \mu_b + g_{ab} \sigma \\ \nabla_a \rho - J_a{}^b \mu_b \end{bmatrix} + \begin{bmatrix} -\mu_a \\ J_{ab} \rho \\ 0 \end{bmatrix}.$$

The induced operator  $\nabla: \Lambda^1 \otimes \mathcal{T} \rightarrow \Lambda^2 \otimes \mathcal{T}$  is

$$\begin{bmatrix} \sigma_b \\ \mu_{bc} \\ \rho_b \end{bmatrix} \mapsto \begin{bmatrix} \nabla_{[a}\sigma_{b]} \\ \nabla_{[a}\mu_{b]}{}_c + g_{c[a}\sigma_{b]} \\ \nabla_{[a}\rho_{b]} - J_{[a}{}^c\mu_{b]}{}_c \end{bmatrix} + \begin{bmatrix} \mu_{[ab]} \\ -J_{c[a}\rho_{b]} \\ 0 \end{bmatrix}$$

but Corollary 2 says the tractor connection on  $\mathbb{C}\mathbb{P}_n$  is symplectically flat so we should contemplate  $\nabla_{\perp}: \Lambda^1 \otimes \mathcal{T} \rightarrow \Lambda^2_{\perp} \otimes \mathcal{T}$  from Theorem 1, viz.

$$\begin{bmatrix} \sigma_b \\ \mu_{bc} \\ \rho_b \end{bmatrix} \mapsto \begin{bmatrix} \nabla_{[a}\sigma_{b]} - \frac{1}{2n}J^{cd}\nabla_c\sigma_d J_{ab} \\ \dots \\ \dots \end{bmatrix} + \begin{bmatrix} \mu_{[ab]} - \frac{1}{2n}J^{cd}\mu_{cd}J_{ab} \\ -J_{c[a}\rho_{b]} - \frac{1}{2n}\rho_c J_{ab} \\ 0 \end{bmatrix}.$$

From these formulæ, let us focus attention on the homomorphisms

$$(22) \quad \begin{array}{ccccccc} 0 & \rightarrow & \mathcal{T} & \rightarrow & \Lambda^1 \otimes \mathcal{T} & \rightarrow & \Lambda^2_{\perp} \otimes \mathcal{T} & \rightarrow & \dots \\ & & \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix} & \mapsto & \begin{bmatrix} -\mu_a \\ J_{ab}\rho \\ 0 \end{bmatrix} & & & & \\ & & & & \begin{bmatrix} \sigma_b \\ \mu_{bc} \\ \rho_b \end{bmatrix} & \mapsto & \begin{bmatrix} \mu_{[ab]} - \frac{1}{2n}J^{cd}\mu_{cd}J_{ab} \\ -J_{c[a}\rho_{b]} - \frac{1}{2n}\rho_c J_{ab} \\ 0 \end{bmatrix} & & \end{array}$$

It is evident that this is a complex and that its cohomology so far is

$$\Lambda^0 \text{ in degree } 0 \quad \text{and} \quad \odot^2\Lambda^1 \text{ in degree } 1.$$

On the other hand, one may check that the defining representation of the Lie algebra  $\mathfrak{sp}(2n+2, \mathbb{R})$  on  $\mathbb{R}^{2n+2} = \mathbb{R} \oplus \mathbb{R}^{2n} \oplus \mathbb{R}$  restricts via (16) to a representation of the Heisenberg Lie algebra  $\mathfrak{h} = \mathbb{R} \oplus \mathfrak{g}_{-1}$ , given explicitly by

$$\mathbb{R}^{2n+2} \xrightarrow{\theta} \mathbb{R}^{2n+2} \quad \text{and} \quad \mathbb{R}^{2n+2} \xrightarrow{\partial_a} \mathfrak{g}_1 \otimes \mathbb{R}^{2n+2}$$

$$\begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix} \mapsto \begin{bmatrix} \rho \\ 0 \\ 0 \end{bmatrix} \quad \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix} \mapsto \begin{bmatrix} -\mu_a \\ J_{ab}\rho \\ 0 \end{bmatrix}$$

(noticing that equations (19) hold, as they must). We may also find  $\theta$  as part of the curvature of the tractor connection on  $\mathbb{C}\mathbb{P}_n$ . Specifically, the formula from Proposition 1 reduces to

$$(23) \quad (\nabla_a\nabla_a - \nabla_b\nabla_a) \begin{bmatrix} \sigma \\ \mu_d \\ \rho \end{bmatrix} = 2J_{ab} \begin{bmatrix} \rho \\ J_d{}^e\mu_e \\ -\sigma \end{bmatrix}$$

and we find  $\theta$  as the top component of  $\Theta: \mathcal{T} \rightarrow \mathcal{T}$  where  $\Theta$  is defined by (5). If we now consider the entire complex from Theorem 1, with filtration induced by

$$\Lambda^0 \subset \Lambda^1 \oplus \Lambda^0 \subset \Lambda^0 \oplus \Lambda^1 \oplus \Lambda^0 = \mathcal{T}$$

$$\begin{bmatrix} 0 \\ 0 \\ \rho \end{bmatrix} \quad \begin{bmatrix} 0 \\ \mu_b \\ \rho \end{bmatrix} \quad \begin{bmatrix} \sigma \\ \mu_b \\ \rho \end{bmatrix}$$

of  $\mathcal{T}$ , then the associated spectral sequence (or corresponding diagram chasing) yields (22) continuing as in (18) including the middle operator  $\nabla_{\perp}^2 - \frac{2}{n}\theta: \wedge_{\perp}^n \rightarrow \wedge_{\perp}^n$ . The same reasoning pertains for any Fedosov structure with  $V_{ab}{}^c{}_d = 0$  as in Corollary 1. Evidently, this sequence of vector bundle homomorphisms is induced by the complex (18) and, together with Lemma 3, the spectral sequence of a filtered complex (or the appropriate diagram chasing) immediately yields the following.

**Theorem 3.** *Suppose  $\nabla_a$  is a torsion-free connection on a symplectic manifold  $(M, J_{ab})$ , such that  $\nabla_a J_{bc} = 0$  and so that the corresponding curvature tensor  $V_{ab}{}^c{}_d$  vanishes. Fix a finite-dimensional representation  $\mathbb{E}$  of  $\mathrm{Sp}(2n + 2, \mathbb{R})$  and let  $E$  denote the associated ‘tractor bundle’ induced from the standard tractor bundle and the representation  $\mathbb{E}$  (so that the standard representation of  $\mathrm{Sp}(2n + 2, \mathbb{R})$  on  $\mathbb{R}^{2n+2}$  yields the standard tractor bundle). In accordance with Corollary 1, the induced ‘tractor connection’  $\nabla: E \rightarrow \wedge^1 \otimes E$  is symplectically flat and we may define  $\Theta: E \rightarrow E$  by (5). Having done this, there are complexes of differential operators*

$$\begin{array}{ccccccc}
 0 \rightarrow & H^0(\mathfrak{h}, E) & \rightarrow & H^1(\mathfrak{h}, E) & \rightarrow & H^2(\mathfrak{h}, E) & \rightarrow \dots \rightarrow H^n(\mathfrak{h}, E) \\
 & & & & & & \downarrow \\
 0 \leftarrow & H^{2n+1}(\mathfrak{h}, E) & \leftarrow & H^{2n}(\mathfrak{h}, E) & \leftarrow & H^{2n-1}(\mathfrak{h}, E) & \leftarrow \dots \leftarrow H^{n+1}(\mathfrak{h}, E)
 \end{array}$$

where  $H^r(\mathfrak{h}, E)$  denotes the tensor bundle on  $M$  that is induced by the cohomology  $H^r(\mathfrak{h}, \mathbb{E})$  as a representation of  $\mathrm{Sp}(2n, \mathbb{R})$ . This complex is locally exact except near the beginning where

$$\ker: H^0(\mathfrak{h}, E) \rightarrow H^1(\mathfrak{h}, E) \quad \text{and} \quad \frac{\ker: H^1(\mathfrak{h}, E) \rightarrow H^2(\mathfrak{h}, E)}{\mathrm{im}: H^0(\mathfrak{h}, E) \rightarrow H^1(\mathfrak{h}, E)}$$

may be identified with the locally constant sheaves  $\ker \Theta$  and  $\mathrm{coker} \Theta$ , respectively. In particular, for  $\mathbb{C}\mathbb{P}_n$  with its Fubini–Study connection, these sheaves vanish and the complex is locally exact everywhere.

**Proof.** It remains only to observe that for the Fubini–Study connection we see from (23) that  $\Theta: \mathcal{T} \rightarrow \mathcal{T}$  is an isomorphism. □

The main point about Theorem 3, however, is that if the representation  $\mathbb{E}$  of  $\mathrm{Sp}(2n + 2, \mathbb{R})$  is irreducible, then the representations  $H^r(\mathfrak{h}, \mathbb{E})$  of  $\mathrm{Sp}(2n, \mathbb{R})$  are also irreducible and are computed by a theorem due to Kostant [14]. Specifically, if we denote the irreducible representations of  $\mathrm{Sp}(2n + 2, \mathbb{R})$  and  $\mathrm{Sp}(2n, \mathbb{R})$  by writing the highest weight as a linear combination of fundamental weights and recording the coefficients over the corresponding nodes of the Dynkin diagrams for  $C_{n+1}$  and  $C_n$ , as is often done, then Kostant’s Theorem says that

$$\begin{array}{l}
 H^0(\mathfrak{h}, \overset{a}{\bullet} \overset{b}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet}) = \overset{b}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet} \\
 H^1(\mathfrak{h}, \overset{a}{\bullet} \overset{b}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet}) = \overset{a+b+1}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet} \\
 H^2(\mathfrak{h}, \overset{a}{\bullet} \overset{b}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet}) = \overset{a}{\bullet} \overset{b+c+1}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet} \\
 H^3(\mathfrak{h}, \overset{a}{\bullet} \overset{b}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet}) = \overset{a}{\bullet} \overset{b}{\bullet} \overset{c+d+1}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet} \\
 \vdots \\
 H^n(\mathfrak{h}, \overset{a}{\bullet} \overset{b}{\bullet} \overset{c}{\bullet} \overset{d}{\bullet} \dots \overset{e}{\bullet} \overset{f}{\bullet}) = \overset{a}{\bullet} \overset{b}{\bullet} \overset{c}{\bullet} \dots \overset{e+f+1}{\bullet}
 \end{array}$$

and for  $r \geq n + 1$ , there are isomorphisms  $H^r(\mathfrak{h}, \mathbb{E}) = H^{2n+1-r}(\mathfrak{h}, \mathbb{E})$ . Using the same notation for the bundles  $H^r(\mathfrak{h}, E)$ , the complexes of Theorem 3 become

$$\begin{array}{c}
 \begin{array}{c}
 \bullet \xrightarrow{b} \bullet \xrightarrow{c} \bullet \xrightarrow{d} \cdots \bullet \xrightarrow{e} \bullet \xrightarrow{f} \bullet \\
 \xrightarrow{\nabla^{a+1}} \bullet \xrightarrow{a+b+1} \bullet \xrightarrow{c} \bullet \xrightarrow{d} \cdots \bullet \xrightarrow{e} \bullet \xrightarrow{f} \bullet \\
 \xrightarrow{\nabla^{b+1}} \bullet \xrightarrow{a} \bullet \xrightarrow{b+c+1} \bullet \xrightarrow{d} \cdots \bullet \xrightarrow{e} \bullet \xrightarrow{f} \bullet \\
 \xrightarrow{\nabla^{c+1}} \bullet \xrightarrow{a} \bullet \xrightarrow{b} \bullet \xrightarrow{c+d+1} \cdots \bullet \xrightarrow{e} \bullet \xrightarrow{f} \bullet \\
 \vdots \\
 \xrightarrow{\nabla^{e+1}} \bullet \xrightarrow{a} \bullet \xrightarrow{b} \bullet \xrightarrow{c} \cdots \bullet \xrightarrow{e+f+1} \bullet \\
 \xrightarrow{\nabla^{2f+2}} \bullet \xrightarrow{a} \bullet \xrightarrow{b} \bullet \xrightarrow{c} \cdots \bullet \xrightarrow{e+f+1} \bullet \\
 \xrightarrow{\nabla^{e+1}} \cdots \\
 \vdots \\
 \xrightarrow{\nabla^{a+1}} \bullet \xrightarrow{b} \bullet \xrightarrow{c} \bullet \xrightarrow{d} \cdots \bullet \xrightarrow{e} \bullet \xrightarrow{f} \bullet,
 \end{array}
 \end{array}$$

for arbitrary non-negative integers  $a, b, c, d, \dots, e, f$ . When all these integers are zero, this is the Rumin–Seshadri complex. Just the first three terms in this complex, in the special case when only  $a$  is non-zero, are already essential in [10]. For example, if  $a = 1$ , then the first two differential operators are

$$\sigma \mapsto \nabla_a \nabla_b \sigma + P_{ab} \sigma \quad \text{and} \quad \phi_{bc} \mapsto (\nabla_a \phi_{bc} - \nabla_b \phi_{ac})_{\perp}$$

where  $\phi_{bc}$  is symmetric and  $(\ )_{\perp}$  means to take the trace-free part with respect to  $J_{ab}$ . From the curvature decomposition and Bianchi identity we find that their composition is

$$\sigma \mapsto V_{ab}{}^d{}_c \nabla_d \sigma + Y_{abc} \sigma,$$

which vanishes in case  $V_{ab}{}^c{}_d = 0$ . In case  $\Theta$  is invertible, as for the Fubini–Study connection, we conclude that this sequence of differential operators is locally exact.

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