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NATURAL 2-FORMS ON THE TANGENT BUNDLE OF A RIEMANNIAN MANIFOLD

JOSEF JANYŠKA

ABSTRACT. 1-order natural differential operators from metrics to 2-forms on the tangent bundle are classified. Some natural transformations from TT to T^*T for Riemannian manifolds are described.

Introduction

It is very well known that on the cotangent bundle $q_M : T^*M \rightarrow M$ of a manifold M there is the canonical symplectic 2-form given by the exterior differential of the Liouville 1-form. Similar canonical construction on the tangent bundle $p_M : TM \rightarrow M$ is not possible. But we can construct the canonical symplectic form on the tangent bundle of a Riemannian manifold. Namely, let (M, g) be a Riemannian manifold and let $h(u) = \frac{1}{2}g(u, u)$, $u \in TM$, be the induced function on TM . The canonical symplectic 2-form on TM is given by

$$\Omega(g) = dd_v h,$$

where d_v denotes the vertical differential, [G].

From the point of view of natural geometry, [N], [KMS], [KJ], Ω is a natural 1-order differential operator, over the identity of T , from the natural bundle functor of Riemannian metrics to the natural bundle functor of exterior 2-forms on the tangent bundle. 2-forms on the tangent bundle of a Riemannian manifold which arise as the results of natural operators from metrics will be called *natural 2-forms* on TM . The aim of this paper is to give the full classification of natural 2-forms of order 1 on TM . We deduce that the family of natural 2-forms on TM depends on some smooth functions of one variable.

Kowalski and Sekizawa, [KS], gave the full classification of natural symmetric (0,2)-tensor fields of order 1 on TM which, together with our results, gives the complete classification of natural (0,2)-tensor fields on TM .

This paper is in final form and no version of it will be submitted for publication elsewhere.

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Kolář and Radziszewski, [KR], classified all natural transformations $TT^*M \rightarrow T^*TM$. They pointed out that there is no natural equivalence $TTM \rightarrow T^*TM$. It corresponds to the fact that there is no natural symplectic form on TM . But in the case of Riemannian manifolds metrics admit wide possibility to construct natural transformations of TTM to T^*TM . In Section 3 we use the natural transformations by Kolář and Radziszewski, [KR], and natural (0,2)-tensor fields on TM described in Section 2 to show some families of natural transformations $TTM \rightarrow T^*TM$ for Riemannian manifolds.

All manifolds and mappings are assumed to be infinitely differentiable.

1. The canonical example

Let M be a manifold with a Riemannian metric g and (x^i) be local coordinates on M . Then

$$g_x = g_{ij}(x)dx^i \odot dx^j, \quad g_{ij}(x) = g_{ji}(x), \quad \det(g_{ij}(x)) \neq 0.$$

We consider the induced function h on TM , $h(u) = \frac{1}{2}\|u\|^2 = \frac{1}{2}g_x(u, u)$, $u \in T_xM$. The vertical differential of h is a 1-form on TM with the coordinate expression

$$d_v h(u) = \frac{\partial h(u)}{\partial u^i} dx^i = g_{im}(x)u^m dx^i,$$

where (x^i, u^i) are the induced fibred coordinates on TM . The canonical symplectic 2-form on TM is then defined by

$$\Omega_u(g) = dd_v h(u)$$

with coordinate expression

$$\Omega_u(g) = dd_v h(u) = \partial_i g_{mj}(x)u^m dx^i \wedge dx^j - g_{ij}(x)dx^i \wedge du^j.$$

In what follows we shall write $g_{jk,l}$ instead of $\partial_l g_{jk}(x)$. We shall also use the matrix notation

$$(1.1) \quad \Omega_u(g) = \begin{bmatrix} (g_{mj,i} - g_{mi,j})u^m & -g_{ij} \\ g_{ij} & 0 \end{bmatrix}.$$

Now, we shall give another description of the canonical symplectic form, which will be more convenient for our purposes. Let Γ be the Levi-Civita connection on M , i.e. its Christoffel symbols are given by

$$(1.2) \quad \Gamma_{jk}^i = \frac{g^{im}}{2}(g_{mj,k} + g_{mk,j} - g_{jk,m}).$$

Then for any $u \in TM$ the tangent space $T_u TM$ splits with respect to Γ into the horizontal and the vertical subspaces, i.e.

$$T_u TM = H_u \oplus V_u.$$

The connection Γ defines the isomorphism between the vector spaces $T_x M$ and H_u , $p_M(u) = x$. This isomorphism is called the *horizontal lift* and for $\xi_x \in T_x M$ the horizontal lift will be denoted $\xi_u^H \in H_u$.

The *vertical lift* of a vector $\xi_x \in T_x M$ is a vector $\xi_u^V \in V_u$ such that $\xi_u^V(df) = \xi_x f$ for all $f \in C^\infty M$. Here df is considered as a function on TM , i.e. $df(u) = uf$. The vertical lift defines an isomorphism between $T_x M$ and V_u . Obviously, each vector $\zeta \in T_u TM$ can be written in the form $\zeta_u = \xi_u^H + \eta_u^V$, where $\xi, \eta \in T_x M$ are uniquely determined vectors.

Now we can define a 2-form on TM as follows

$$(1.3) \quad \begin{aligned} \Omega_u(g)(\xi^H, \eta^H) &= 0, & \Omega_u(g)(\xi^H, \eta^V) &= -g_x(\xi, \eta), \\ \Omega_u(g)(\xi^V, \eta^H) &= g_x(\eta, \xi), & \Omega_u(g)(\xi^V, \eta^V) &= 0. \end{aligned}$$

for all $\xi, \eta \in T_x M$. The matrix expression of (1.3) is

$$(1.4) \quad \Omega_u(g) = \begin{bmatrix} (g_{mj}\Gamma_{ai}^m - g_{mi}\Gamma_{aj}^m)u^a & -g_{ij} \\ g_{ij} & 0 \end{bmatrix}.$$

From (1.2) we can easily see that the matrix expressions (1.1) and (1.4) coincide and hence (1.3) defines the canonical symplectic form on TM .

Remark 1.1. From (1.3) we see that the canonical symplectic form $\Omega(g)$ is defined by the construction which is similar to the construction of the horizontal lift, [KS], of a metric to a symmetric (0,2)-tensor field on TM . It is why the construction (1.3) is called the *horizontal lift* of a metric to a 2-form on TM .

2. Natural 2-forms on TM

Let $S_+^2 T^* \subset T^* \otimes T^*$ be the natural bundle functor of Riemannian metrics. The canonical symplectic form described in Section 1 is a natural 1-order operator from $S_+^2 T^* \oplus T$ to $\wedge^2 T^*(T)$ over the identity of T . We shall classify all such 1-order (with respect to metric) operators. It is very well known that such operators are in a bijective correspondence with G_n^2 -equivariant mappings from the standard fibre of the bundle functor $J^1(S_+^2) \oplus T$ to the standard fibre of $\wedge^2 T^*(T)$. To determine these equivariant mappings we use the infinitesimal method, [KS], [KJ].

Let us denote $Q = \odot^2 \mathbf{R}^{n*} \times (\odot^2 \mathbf{R}^{n*} \otimes \mathbf{R}^{n*}) \times \mathbf{R}^n$ the standard fibre of $J^1(S_+^2) \oplus T$ and $(g_{ij}, g_{ij,k}, u^i)$ the canonical coordinates on Q . The action of G_n^2 on Q is given by

$$\begin{aligned} \bar{g}_{ij} &= \tilde{a}_i^p \tilde{a}_j^q g_{pq} \\ \bar{g}_{ij,k} &= \tilde{a}_i^p \tilde{a}_j^q \tilde{a}_k^r g_{pq,r} + (\tilde{a}_{ik}^p \tilde{a}_j^q + \tilde{a}_i^p \tilde{a}_{jk}^q) g_{pq} \\ \bar{u}^i &= a_p^i u^p, \end{aligned}$$

where (a_j^i, a_{jk}^i) are the canonical coordinates on G_n^2 and tilde denotes the inverse element.

The fundamental vector fields on Q relative to this action are

$$(2.1) \quad \xi_p^q(Q) = u^q \frac{\partial}{\partial u^p} - 2g_{ap} \frac{\partial}{\partial g_{aq}} - (\delta_a^q g_{pb,c} + \delta_b^q g_{ap,c} + \delta_c^q g_{ab,p}) \frac{\partial}{\partial g_{ab,c}},$$

$$(2.2) \quad \xi_p^{qr}(Q) = -g_{ap} \left(\frac{\partial}{\partial g_{aq,r}} + \frac{\partial}{\partial g_{ar,q}} \right).$$

Let us denote $S = \mathbf{R}^n \times \wedge^2 \mathbf{R}^{2n*}$ the standard fibre of $\wedge^2 T^*(T)$ with the canonical coordinates $(u^i, \begin{pmatrix} u_{ij}^1 & u_{ij}^2 \\ -u_{ji}^2 & u_{ij}^4 \end{pmatrix})$, $u_{ij}^1 = -u_{ji}^1$, $u_{ij}^4 = -u_{ji}^4$. The action of G_n^2 on S is given by

$$\begin{aligned}\bar{u}^i &= a_p^i u^p, \\ \bar{u}_{ij}^1 &= \tilde{a}_i^p \tilde{a}_j^q u_{pq}^1 + (\tilde{a}_i^p \tilde{a}_{jm}^q - \tilde{a}_j^p \tilde{a}_{im}^q) a_r^m u^r u_{pq}^2 + \tilde{a}_{im}^p \tilde{a}_{jk}^q a_r^m a_s^k u^r u^s u_{pq}^4, \\ \bar{u}_{ij}^2 &= \tilde{a}_i^p \tilde{a}_j^q u_{pq}^2 + \tilde{a}_{im}^p \tilde{a}_j^q a_r^m u^r u_{pq}^4, \\ \bar{u}_{ij}^4 &= \tilde{a}_i^p \tilde{a}_j^q u_{pq}^4.\end{aligned}$$

The fundamental vector fields on S relative to this action are

$$(2.3) \quad \xi_p^q(S) = u^q \frac{\partial}{\partial u^p} - 2u_{ap}^1 \frac{\partial}{\partial u_{aq}^1} - u_{pa}^2 \frac{\partial}{\partial u_{qa}^2} - u_{ap}^2 \frac{\partial}{\partial u_{aq}^2} - 2u_{ap}^4 \frac{\partial}{\partial u_{aq}^4},$$

$$(2.4) \quad 2\xi_p^{qr}(S) = (\delta_a^q u^r u_{bp}^2 + \delta_a^r u^q u_{bp}^2 - \delta_b^q u^r u_{ap}^2 - \delta_b^r u^q u_{ap}^2) \frac{\partial}{\partial u_{ab}^1} - (\delta_a^q u^r u_{pb}^4 + \delta_a^r u^q u_{pb}^4) \frac{\partial}{\partial u_{ab}^2}.$$

A mapping $F: Q \rightarrow S$ is G_n^2 -equivariant iff the corresponding fundamental vector fields are F -related. If F has the coordinate expression

$$u^i = u^i, \quad u_{ij}^\alpha = F_{ij}^\alpha(g_{ab}, g_{ab,c}, u^a), \quad \alpha = 1, 2, 4,$$

then F_{ij}^α have to satisfy the following system of partial differential equations

$$(2.5) \quad 2g_{ap} \frac{\partial F_{ij}^\alpha}{\partial g_{aq}} + (\delta_a^q g_{pb,c} + \delta_b^q g_{ap,c} + \delta_c^q g_{ab,p}) \frac{\partial F_{ij}^\alpha}{\partial g_{ab,c}} - u^q \frac{\partial F_{ij}^\alpha}{\partial u^p} = F_{ip}^\alpha \delta_j^q + F_{pj}^\alpha \delta_i^q,$$

$$(2.6) \quad g_{pa} \left(\frac{\partial F_{ij}^4}{\partial g_{aq,r}} + \frac{\partial F_{ij}^4}{\partial g_{ar,q}} \right) = 0,$$

$$(2.7) \quad g_{pa} \left(\frac{\partial F_{ij}^2}{\partial g_{aq,r}} + \frac{\partial F_{ij}^2}{\partial g_{ar,q}} \right) = F_{pj}^4 u^r \delta_i^q + F_{pj}^4 u^q \delta_i^r,$$

$$(2.8) \quad 2g_{pa} \left(\frac{\partial F_{ij}^1}{\partial g_{aq,r}} + \frac{\partial F_{ij}^1}{\partial g_{ar,q}} \right) = F_{ip}^2 u^r \delta_j^q + F_{ip}^2 u^q \delta_j^r - F_{jp}^2 u^r \delta_i^q - F_{jp}^2 u^q \delta_i^r.$$

Theorem 2.1. All G_n^2 -equivariant mappings $F : Q \rightarrow S$ are given by the formulas

$$(2.9) \quad \begin{aligned} F_{ij}^1 &= u^b u^c \Gamma_{bi}^r \Gamma_{cj}^s \alpha_{rs} + u^b \Gamma_{bj}^r \beta_{ri} - u^b \Gamma_{bi}^r \beta_{rj} + \gamma_{ij}, \\ F_{ij}^2 &= u^b \Gamma_{bi}^r \alpha_{rj} + \beta_{ij}, \\ F_{ij}^4 &= \alpha_{ij}, \end{aligned}$$

where Γ_{jk}^i are the formal Christoffel symbols and $\alpha_{ij}, \beta_{ij}, \gamma_{ij}$ are functions on Q which are solutions of the following system of differential equations

$$(2.10) \quad \frac{\partial \zeta_{ij}}{\partial g_{pq,r}} = 0,$$

$$(2.11) \quad 2g_{ap} \frac{\partial \zeta_{ij}}{\partial g_{aq}} - u^q \frac{\partial \zeta_{ij}}{\partial u^p} = \zeta_{ip} \delta_j^q + \zeta_{pj} \delta_i^q.$$

Moreover, $\alpha_{ij} = -\alpha_{ji}, \gamma_{ij} = -\gamma_{ji}$.

Proof. We have to show that all solutions of (2.5) - (2.8) are of the form (2.9). We contract both sides of (2.6)-(2.8) by g^{pq} and use the cyclic permutation of the indices p, q, r . We get

$$(2.12) \quad \frac{\partial F_{ij}^4}{\partial g_{pq,r}} = 0,$$

$$(2.13) \quad 4 \frac{\partial F_{ij}^2}{\partial g_{pq,r}} = F_{mj}^4 g^{ma} u^b (\delta_{abi}^{pgr} + \delta_{aib}^{pgr} + \delta_{iab}^{pgr} - \delta_{iba}^{pgr} - \delta_{bia}^{pgr} + \delta_{bai}^{pgr}),$$

$$(2.14) \quad \begin{aligned} 4 \frac{\partial F_{ij}^1}{\partial g_{pq,r}} &= F_{jm}^2 g^{ma} u^b (\delta_{iba}^{pgr} + \delta_{bia}^{pgr} - \delta_{abi}^{pgr} - \delta_{aib}^{pgr} - \delta_{iab}^{pgr} - \delta_{bai}^{pgr}) + \\ &+ F_{im}^2 g^{ma} u^b (\delta_{abj}^{pgr} + \delta_{ajb}^{pgr} + \delta_{jab}^{pgr} - \delta_{jba}^{pgr} - \delta_{bja}^{pgr} + \delta_{baj}^{pgr}), \end{aligned}$$

which can be rewrite, by using (1.2), in the form

$$(2.15) \quad \frac{\partial F_{ij}^2}{\partial g_{pq,r}} = F_{mj}^4 u^s \frac{\partial \Gamma_{si}^m}{\partial g_{pq,r}},$$

$$(2.16) \quad \frac{\partial F_{ij}^1}{\partial g_{pq,r}} = F_{im}^2 u^s \frac{\partial \Gamma_{sj}^m}{\partial g_{pq,r}} - F_{jm}^2 u^s \frac{\partial \Gamma_{si}^m}{\partial g_{pq,r}}.$$

Putting $F_{ij}^4 = \alpha_{ij}$ and substituting it into (2.15) we get after the integration

$$(2.17) \quad F_{ij}^2 = u^b \Gamma_{bi}^r \alpha_{rj} + \beta_{ij},$$

where $\frac{\partial \beta_{ij}}{\partial g_{pq,r}} = 0$ and substituting (2.17) into (2.16) we get after the integration

$$(2.18) \quad F_{ij}^1 = u^b u^c \Gamma_{bi}^r \Gamma_{cj}^s \alpha_{rs} + u^b \Gamma_{bj}^r \beta_{ri} - u^b \Gamma_{bi}^r \beta_{rj} + \gamma_{ij},$$

where $\frac{\partial \gamma_{ij}}{\partial g_{pq,r}} = 0$. It is easy to see that $\alpha_{ij} = -\alpha_{ji}, \gamma_{ij} = -\gamma_{ji}$ and α_{ij}, β_{ij} and γ_{ij} satisfy (2.11). $\square$

Remark 2.1. From (2.10) and (2.11) it follows that $\alpha_{ij}, \beta_{ij}, \gamma_{ij}$ are the components of (0,2)-tensor fields on M which are given as 0-order natural differential operators from $S_+^2 \oplus T$ to $\otimes^2 T^*$. Such natural tensor fields are called *natural F -metrics*, [KS].

Now we can easily prove

Theorem 2.2. All natural 2-forms $\Omega(g)$ of order 1 on TM are of the form

$$(2.19) \quad \begin{aligned} \Omega_u(g)(\xi^H, \eta^H) &= \gamma_x(\xi, \eta), & \Omega_u(g)(\xi^H, \eta^V) &= \beta_x(\xi, \eta), \\ \Omega_u(g)(\xi^V, \eta^H) &= -\beta_x(\eta, \xi), & \Omega_u(g)(\xi^V, \eta^V) &= \alpha_x(\xi, \eta). \end{aligned}$$

where α, β, γ are natural F -metrics and moreover γ and α are skew-symmetric.

Proof. It is easy to see that the coordinate expression of (2.19) coincides with (2.9). $\square$

We recall here the classifying theorem for natural 1-order symmetric (0,2)-tensors on TM by Kowalski and Sekizawa, [KS], (see also [KMS]).

Theorem 2.3. All natural 1-order symmetric (0,2)-tensor fields $G(g)$ on TM are of the form

$$(2.20) \quad \begin{aligned} G_u(g)(\xi^H, \eta^H) &= \gamma_x(\xi, \eta), & G_u(g)(\xi^H, \eta^V) &= \beta_x(\xi, \eta), \\ G_u(g)(\xi^V, \eta^H) &= \beta_x(\eta, \xi), & G_u(g)(\xi^V, \eta^V) &= \alpha_x(\xi, \eta). \end{aligned}$$

where α, β, γ are natural F -metrics and moreover γ and α are symmetric. $\square$

Hence the problem of classifying natural (0,2)-tensor fields of order 1 on TM is reduced to the problem of classifying natural F -metrics. This problem was completely solved by Kowalski and Sekizawa, [KS].

Theorem 2.4. Let (M, g) be an oriented Riemannian manifold of dimension n . Then all natural F -metrics on M derived from g are given as follows:

i) For $n = 1$, all natural F -metrics are of the form

$$(2.21) \quad \zeta_u(\xi, \eta) = \mu(\|u\|^2)g(\xi, \eta),$$

where μ is an arbitrary function of $\|u\|^2 = g(u, u)$.

ii) For $n = 2$, all symmetric natural F -metrics are of the form

$$(2.22) \quad \begin{aligned} \zeta_u(\xi, \eta) &= \mu(\|u\|^2)g(\xi, \eta) + \nu(\|u\|^2)g(\xi, u)g(\eta, u) + \\ &+ \kappa(\|u\|^2)[g(\xi, u)g(\eta, Ju) + g(\eta, u)g(\xi, Ju)], \end{aligned}$$

and all skew-symmetric natural F -metrics are of the form

$$(2.23) \quad \zeta_u(\xi, \eta) = \lambda(\|u\|^2)[g(\xi, u)g(\eta, Ju) - g(\eta, u)g(\xi, Ju)],$$

where $\mu, \nu, \kappa, \lambda$ are arbitrary functions of $\|u\|^2$ and J is one of the two canonical almost complex structures on (M, g) .

iii) For $n = 3$, all symmetric natural F -metrics are of the form

$$(2.24) \quad \zeta_u(\xi, \eta) = \mu(\|u\|^2)g(\xi, \eta) + \nu(\|u\|^2)g(\xi, u)g(\eta, u),$$

and all skew-symmetric natural F -metrics are of the form

$$(2.25) \quad \zeta_u(\xi, \eta) = \lambda(\|u\|^2)g(\xi \times \eta, u),$$

where μ, ν, λ are arbitrary functions of $\|u\|^2$ and $\xi \times \eta$ is the usual vector product of ξ and η .

iv) For $n > 3$, all natural F -metrics are symmetric and are of the form (2.24). $\square$

Remark 2.2. M was supposed to be oriented in Theorem 2.4. In the case of non-oriented Riemannian manifold all natural F -metrics are of the form (2.21), for $n = 1$, and (2.24), for $n \geq 2$.

Remark 2.3. If we combine Theorem 2.2 and Theorem 2.4 we get the following: If $n = 1$, the family of all natural 2-forms depends on one arbitrary function of one variable, for $n = 2$ it depends on six arbitrary functions of one variable, for $n = 3$ on five functions and for $n > 3$ on two arbitrary functions of one variable.

Remark 2.4. In the general case (for oriented Riemannian manifolds if $n = 1$ or $n > 3$, for non-oriented Riemannian manifolds in all dimensions) all natural 2-forms are horizontal lifts of a natural F -metric to a 2-form on TM . The canonical symplectic form is then given for $\zeta = -g$.

Remark 2.5. The restriction of order of our operators is necessary. If we consider higher order operations we get many further natural $(0,2)$ -tensor fields on TM . For instance let (M, ∇) be a manifold with a linear connection ∇ . Let R be the Ricci tensor of ∇ . Then

$$\begin{aligned} \Omega_u(g)(\xi^H, \eta^H) &= 0, & \Omega_u(g)(\xi^H, \eta^V) &= R_x(\xi, \eta), \\ \Omega_u(g)(\xi^V, \eta^H) &= -R_x(\eta, \xi), & \Omega_u(g)(\xi^V, \eta^V) &= 0. \end{aligned}$$

is a 2-form on TM which is naturally induced from ∇ and is of order 1 with respect to ∇ . Hence, if M is a Riemannian manifold and ∇ is the Levi-Civita connection, we get 2-order (with respect to a metric) natural 2-form on TM .

3. Some natural transformations $TT \rightarrow T^*T$

Kolář and Radziszewski, [KR], classified all natural transformations of the bundle functor TT^* to T^*T . They pointed out that there is no natural equivalence of TT to T^*T . It is a consequence of different geometrical properties of these bundle functors and it corresponds to the fact that there is no canonical natural symplectic form on TM . In Section 1 we have constructed the canonical natural (with respect to a metric) symplectic form Ω on the tangent bundle of a Riemannian manifold, which gives a natural transformation $S_\Omega : TTM \rightarrow T^*TM$. This transformation is in fact a natural differential operator

$$S : S^2_+ T^* \oplus TT \rightarrow T^*T$$

of order 1 in metrics. In this section we shall give some natural transformations $TT \rightarrow T^*T$ for Riemannian manifolds, which are of order 1 with respect to metrics.

First we recall the main result by Kolář and Radziszewski, [KR]. We give two canonical natural transformations $TT^* \rightarrow T^*T$. The first is the transformation $s_M : TT^*M \rightarrow T^*TM$ by Modugno and Stefani, [MS], which can be described geometrically as follows. Every $A \in TT^*M$ is a vector tangent to a curve $\gamma(t) : \mathbf{R} \rightarrow T^*M$ at $t = 0$. If $B \in T_{T_{q_M(A)}TM}$, then iB is tangent to the curve $\delta(t) : \mathbf{R} \rightarrow TM$ over the curve $q_M(\gamma(t))$ on M . $i : TTM \rightarrow T^*TM$ is the canonical involution. Hence we can evaluate $\langle \gamma(t), \delta(t) \rangle$ for every t and the derivative $\frac{d}{dt}|_0 \langle \gamma(t), \delta(t) \rangle =: \sigma(A, B)$ depends on A and B only. This determines a linear map $T_{T_{q_M(A)}TM} \rightarrow \mathbf{R}$, $B \mapsto \sigma(A, B)$, i.e. an element $s_M(A) \in T^*TM$.

The second construction is the following. We have the injection $\kappa_M : T^*M \rightarrow T^*TM$ given by the pullback with respect to the projection $p_M : TM \rightarrow M$. I.e. $\kappa_M(A)(B) = \langle A, T p_M(B) \rangle$, $A \in T_x^*M$, $B \in T_x TM$. Then $\kappa_M \circ p_{T^*M} : TT^*M \rightarrow T^*TM$ is a natural transformation $TT^*M \rightarrow T^*TM$ such that the diagram

$$\begin{array}{ccccc}
 & & TT^*M & & \\
 & \swarrow p_{T^*M} & \downarrow & \searrow T_{q_M} & \\
 T^*M & \xrightarrow{\kappa_M} & T^*TM & \xrightarrow{q_{TM}} & TM \\
 & \searrow q_M & & \swarrow p_M & \\
 & & M & &
 \end{array}$$

commutes.

Finally we denote $Y \mapsto (k)_1 Y$ and $Y \mapsto (k)_2 Y$, $k \in \mathbf{R}$, the scalar multiplications in TT^*M with respect to two vector bundle structures $p_{T^*M} : TT^*M \rightarrow T^*M$ and $T_{q_M} : TT^*M \rightarrow TM$, respectively. In the notation $X \in TT^*M$, $p = p_{T^*M}(X) \in T^*M$, $\xi = T_{q_M}(X) \in TM$ we have, [KR], [KMS],

Theorem 3.1. *All natural transformations of TT^*M to T^*TM are of the form*

$$(3.1) \quad (F(\langle p, \xi \rangle))_1 (G(\langle p, \xi \rangle))_2 s_M(X) + \kappa_M(H(\langle p, \xi \rangle)p),$$

where $F(t), G(t), H(t)$ are three arbitrary smooth functions of one variable. $\square$

Let us express (3.1) in coordinates. If (x^i) are local coordinates on M , then we have the induced fibred coordinates (x^i, u^i) on TM , (x^i, u^i, ξ_i, U_i) on T^*TM and (x^i, p_i, ξ^i, π_i) on TT^*M . Then the coordinate expression of (3.1) is

$$\begin{aligned}
 u^i &= F(p_m \xi^m) \xi^i, \\
 \xi_i &= F(p_m \xi^m) G(p_m \xi^m) \pi_i + H(p_m \xi^m) p_i, \\
 U_i &= G(p_m \xi^m) p_i.
 \end{aligned}
 \quad (3.2)$$

The canonical transformation s_M is then given by $F = 1, G = 1, H = 0$, i.e.

$$u^i = \xi^i, \quad \xi_i = \pi_i, \quad U_i = p_i$$

and $\kappa_M \circ p_{T^*M}$ is given by $F = 1, G = 0, H = 1$, i.e.

$$u^i = \xi^i, \quad \xi_i = p_i, \quad U_i = 0.$$

Now we are in the position to describe some natural transformations of TTM to T^*TM for a Riemannian manifolds by using Theorem 3.1. Let us suppose that we have a $(0,2)$ -tensor field ζ on M which defines the mapping $S_\zeta : TM \rightarrow T^*M$, over M , by

$$(3.3) \quad \langle S_\zeta(u), \xi \rangle_x = \zeta_x(\xi, u), \quad \xi, u \in T_x M.$$

Then we can define two families of natural transformations (natural also with respect to ζ) of TTM to T^*TM . The first family is given by the commutativity of the diagram

$$\begin{array}{ccc} TTM & \xrightarrow{TS_\zeta} & TT^*M \\ & \searrow Z_M & \swarrow \Sigma_M \\ & T^*TM & \end{array}$$

where Σ_M are the natural transformations from Theorem 3.1. Transformations $Z_M : TTM \rightarrow T^*TM$ are over the natural transformation of TM given by the scalar multiplication in TM such that the diagram

$$\begin{array}{ccc} TTM & \xrightarrow{Z_M} & T^*TM \\ T p_M \downarrow & & \downarrow q_{TM} \\ TM & \xrightarrow{F(\zeta_{km}\xi^k u^m)} & TM \end{array}$$

commutes, where F is the function from Theorem 3.1. The coordinate expression of the family Z_M is

$$\begin{aligned} u^i &= F(\zeta_{km}\xi^k u^m)\xi^i, \\ \xi_i &= F(\zeta_{km}\xi^k u^m)G(\zeta_{km}\xi^k u^m)(\zeta_{im,k}u^m\xi^k + \zeta_{im}\Xi^m) + H(\zeta_{km}\xi^k u^m)\zeta_{im}u^m, \\ U_i &= G(\zeta_{km}\xi^k u^m)\zeta_{im}u^m, \end{aligned}$$

where (x^i, u^i, ξ^i, Ξ^i) are the induced fibred coordinates on TTM .

The second family of natural transformations of TTM to T^*TM is given by the following commutative diagram

$$\begin{array}{ccc} TTM & \xrightarrow{TS_\zeta} & TT^*M \\ i_M \uparrow & & \downarrow \Sigma_M \\ TTM & \xrightarrow{\tilde{Z}_M} & T^*TM \end{array}$$

where i_M is the canonical involution of TTM . The family of natural transformations $\tilde{Z}_M$ is over the scalar multiplication in TM via the commutative diagram

$$\begin{array}{ccc} TTM & \xrightarrow{\tilde{Z}_M} & T^*TM \\ p_{TM} \downarrow & & \downarrow q_{TM} \\ TM & \xrightarrow{F(\zeta_{km}u^k\xi^m)} & TM \end{array}$$

The coordinate expression of the family $\tilde{Z}_M$ is

$$\begin{aligned} u^i &= F(\zeta_{km}u^k\xi^m)u^i, \\ \xi_i &= F(\zeta_{km}u^k\xi^m)G(\zeta_{km}u^k\xi^m)(\zeta_{im,k}u^k\xi^m + \zeta_{im}\Xi^m) + H(\zeta_{km}u^k\xi^m)\zeta_{im}\xi^m, \\ U_i &= G(\zeta_{km}u^k\xi^m)\zeta_{im}\xi^m. \end{aligned}$$

Now if ζ is a natural F -metric from Theorem 2.4 we get two families of natural transformations of TTM to T^*TM for Riemannian manifolds. These families depend on functions of one variable via the natural F -metric and via Σ_M , where as arguments appear $g(u, u)$, $g(\xi, \xi)$ and $g(u, \xi)$.

The third possibility how to construct natural transformations of TTM to T^*TM is to use a natural lift of a metric to $(0,2)$ -tensor fields on TM . Namely, if $\Omega(g)$ is a natural lift described in Theorem 2.2 or 2.3, then

$$S_{\Omega(g)} : TTM \rightarrow T^*TM$$

defined by (3.3) is a natural transformation. All these transformations are over the identity of TM and some of them are contained in the family $\tilde{Z}_M$.

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