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## ABOUT DUALITY AND KILLING TENSORS

DUMITRU BALEANU

**ABSTRACT.** In this paper the isometries of the dual space were investigated. The dual structural equations of a Killing tensor of order two were found. The general results are applied to the case of the flat space.

### 1. INTRODUCTION

Killing tensors are indispensable tools in the quest for exact solutions in many branches of general relativity as well as classical mechanics [1]. Killing tensors are important for solving the equations of motion in particular space-times. The notable example here is the Kerr metric which admits a second rank Killing tensor [1]. The Killing tensors give rise to new exact solutions in perfect fluid Bianchi and Katowski-Sachs cosmologies as well in inflationary models with a scalar field sources [2]. Recently the Killing tensors of third rank in  $(1+1)$  dimensional geometry were investigated and classified [3]. In a geometrical setting, symmetries are connected with isometries associated with Killing vectors, and more generally, Killing tensors on the configuration space of the system. An example is the motion of a point particle in a space with isometries [4], which is a physicist's way of studying the geodesic structure of a manifold. Any symmetrical tensor  $K_{\alpha\beta}$  satisfying the condition

$$K_{(\alpha\beta;\gamma)} = 0, \quad (1)$$

is called a Killing tensor. Here the parenthesis denotes a full symmetrization with all indices and coma denotes a covariant derivative.  $K_{\alpha\beta}$  will be called redundant if it is equal to some linear combination with constants coefficients of the metric tensor  $g_{\alpha\beta}$  and of the form  $S_{(\alpha}B_{\beta)}$  where  $A_\alpha$  and  $B_\beta$  are Killing vectors. For any Killing vector  $K_\alpha$  we have [5]

$$K_{\beta;\alpha} = \omega_{\alpha\beta} = -\omega_{\beta\alpha}, \quad (2)$$

$$\omega_{\alpha\beta;\gamma} = R_{\alpha\beta\gamma\delta}K^\delta. \quad (3)$$

The equations (2) and (3) may be regarded as a system of linear homogeneous first-order equations in the components  $K_\beta$  and  $\omega_{\alpha\beta}$ . The equations analogous to the above ones for a Killing vector were derived for  $K_{\alpha\beta}$  in [5].

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The paper is in the final form and no version of it will be submitted elsewhere.

Recently Holten has presented a theorem concerning the reciprocal relation between two local geometries described by metrics which are Killing tensors with respect to one another [6].

In this paper the geometric duality was presented and the structural equations of  $K_{\alpha\beta}$  were analyzed.

The plan of the paper is as follows.

In Sec.2 the geometric duality is presented. In Sec.3 the structural equations of  $K_{\alpha\beta}$  are investigated. Our comments and concluding remarks are presented in Sec.4.

## 2. GEOMETRIC DUALITY

Let us consider that the space with a metric  $g_{\alpha\beta}$  admits a Killing tensor field  $K_{\alpha\beta}$ .

As it is well known the equation of motion of a particle on a geodesic is derived from the action [7]

$$S = \int d\tau \left( \frac{1}{2} g_{\alpha\beta} \dot{x}^\alpha \dot{x}^\beta \right). \quad (4)$$

The Hamiltonian is constructed in the following form  $H = \frac{1}{2} g_{\alpha\beta} p^\alpha p^\beta$  and the Poisson brackets are

$$\{x_\alpha, p^\beta\} = \delta_\alpha^\beta. \quad (5)$$

The equation of motion for a phase space function  $F(x, p)$  can be computed from the Poisson brackets with the Hamiltonian

$$\dot{F} = \{F, H\}, \quad (6)$$

where  $\dot{F} = \frac{dF}{d\tau}$ . From the covariant component  $K_{\alpha\beta}$  of the Killing tensor we can construct a constant of motion  $K$

$$K = \frac{1}{2} K_{\alpha\beta} p^\alpha p^\beta. \quad (7)$$

We can easily verify that

$$\{H, K\} = 0. \quad (8)$$

The formal similarity between the constants of motion  $H$  and  $K$ , and the symmetrical nature of the condition implying the existence of the Killing tensor amount to a reciprocal relation between two different models: the model with Hamiltonian  $H$  and constant of motion  $K$ , and a model with constant of motion  $H$  and Hamiltonian  $K$ . The relation between the two models has a geometrical interpretation: it implies that if  $K_{\alpha\beta}$  are the contravariant components of a Killing tensor with respect to the metric  $g_{\alpha\beta}$ , then  $g_{\alpha\beta}$  must represent a Killing tensor with respect to the metric defined by  $K_{\alpha\beta}$ . When  $K_{\alpha\beta}$  has an inverse we interpret it as the metric of another space and we can define the associated Riemann-Christoffel connection  $\hat{\Gamma}_{\alpha\beta}^\lambda$  as usual through the metric postulate  $\hat{D}_\lambda K_{\alpha\beta} = 0$ . Here  $\hat{D}$  represents the covariant derivative with respect to  $K_{\alpha\beta}$ . The relation between connections  $\hat{\Gamma}_{\alpha\beta}^\mu$  and  $\Gamma_{\alpha\beta}^\mu$  is [8]

$$\hat{\Gamma}_{\alpha\beta}^\mu = \Gamma_{\alpha\beta}^\mu - K^{\mu\delta} D_\delta K_{\alpha\beta}. \quad (9)$$

As it is well known for a given metric  $g_{\alpha\beta}$  the conformal transformation is defined as  $\hat{g}_{\alpha\beta} = e^{2U(x)}g_{\alpha\beta}$  and the relation between the corresponding connections is

$$\hat{\Gamma}_{\alpha\beta}^{\lambda} = \Gamma_{\alpha\beta}^{\lambda} + 2\delta_{(\alpha}^{\lambda}U_{\beta)}' - g_{\alpha\beta}U'^{\lambda}, \quad (10)$$

where  $U'^{\lambda} = \frac{dU^{\lambda}}{dx}$ . After some calculations we conclude that the dual transformation (9) is not a conformal transformation.

For this reason is interesting to investigate the case when the manifold and its dual have the same isometries. Let us denote by  $\chi_{\alpha}$  a Killing vector corresponding to  $g_{\alpha\beta}$  and by  $\hat{\chi}_{\alpha}$  a Killing vector corresponding to  $K_{\alpha\beta}$ . Then the following Proposition holds.

**Proposition.** *The manifold and its dual have the same Killing vectors iff*

$$(D_{\delta}K_{\alpha\beta})\hat{\chi}^{\delta} = 0. \quad (11)$$

**Proof.** Let us consider  $\chi_{\sigma}$  a vector satisfies

$$D_{\alpha}\chi_{\beta} + D_{\beta}\chi_{\alpha} = 0. \quad (12)$$

Using (9) the corresponding dual Killing vector's equations are

$$D_{\alpha}\hat{\chi}_{\beta} + D_{\beta}\hat{\chi}_{\alpha} + 2K^{\delta\sigma}(D_{\delta}K_{\alpha\beta})\hat{\chi}_{\sigma} = 0. \quad (13)$$

If we assume that  $\hat{\chi}_{\alpha} = \chi_{\alpha}$ , then using (12) and (13) we obtain

$$(D_{\delta}K_{\alpha\beta})\hat{\chi}^{\delta} = 0. \quad (14)$$

Conversely if we suppose that (11) holds, then from (13) we can deduce immediately that  $\chi_{\alpha} = \hat{\chi}_{\alpha}$ .  $\square$

As an example we mention here the separable coordinates in  $(1+1)$  dimensions because in this case  $D_{\lambda}K_{\alpha\beta} = 0$ .

### 3. THE STRUCTURAL EQUATIONS

The following two vectors play roles analogous to that of a bivector  $\omega_{\alpha\beta}$

$$L_{\alpha\beta\gamma} = K_{\beta\gamma;\alpha} - K_{\alpha\gamma;\beta}, \quad (15)$$

$$M_{\alpha\beta\gamma\delta} = \frac{1}{2}(L_{\alpha\beta[\gamma;\delta]} + L_{\gamma\delta[\alpha;\beta]}). \quad (16)$$

The properties of the tensors  $L_{\alpha\beta\gamma}$  and  $M_{\alpha\beta\gamma\delta}$  were derived in [5].  $M_{\alpha\beta\gamma\delta}$  has the same symmetries as the Riemannian tensor and the covariant derivatives of  $K_{\alpha\beta}$  and  $L_{\alpha\beta\gamma}$  satisfy relations reminiscent of those satisfied by Killing vectors.

From (15) and (16) we found that  $M_{\alpha\beta\gamma\delta} = \frac{1}{2}(K_{\beta\gamma;(\alpha\delta)} + K_{\alpha\delta;(\beta\gamma)} - K_{\alpha\gamma;(\beta\delta)} - K_{\beta\delta;(\alpha\gamma)})$  and  $M_{\alpha\beta\gamma\delta} + M_{\gamma\alpha\beta\delta} + M_{\beta\gamma\alpha\delta} = 0$  [5].

Let us define a tensor  $H_{\alpha\beta}^{\mu}$  as

$$H_{\alpha\beta}^{\mu} = \hat{\Gamma}_{\alpha\beta}^{\mu} - \Gamma_{\alpha\beta}^{\mu}. \quad (17)$$

Taking into account (9) we found

$$H_{\alpha\beta}^{\delta}K_{\delta\gamma} = -D_{\gamma}K_{\alpha\beta}. \quad (18)$$

Using (15) and (18) we obtain

$$L_{\alpha\beta\gamma} = H_{\alpha\gamma}^{\delta}K_{\delta\beta} - H_{\beta\gamma}^{\delta}K_{\delta\alpha}. \quad (19)$$

From (19) we conclude that  $L_{\alpha\beta\gamma}$  looks like an angular momentum. This result is in agreement with those presented in [5]. Taking into account (16) and (18) the expression of  $M_{\alpha\beta\gamma\delta}$  becomes

$$\begin{aligned} M_{\alpha\beta\gamma\delta} &= \frac{K_{\sigma\alpha}}{2} [-D_\delta H_{\beta\gamma}^\sigma + D_\gamma H_{\beta\delta}^\sigma + H_{\beta\gamma}^\sigma H_{\delta\beta}^\theta - H_{\gamma\beta}^\theta H_{\delta\delta}^\sigma] \\ &+ \frac{K_{\sigma\beta}}{2} [D_\delta H_{\alpha\gamma}^\sigma - D_\gamma H_{\alpha\delta}^\sigma + H_{\delta\alpha}^\theta H_{\beta\gamma}^\sigma - H_{\gamma\alpha}^\theta H_{\delta\delta}^\sigma] \\ &+ \frac{K_{\sigma\gamma}}{2} [-D_\beta H_{\delta\alpha}^\sigma + D_\alpha H_{\delta\beta}^\sigma - H_{\beta\delta}^\theta H_{\alpha\alpha}^\sigma + H_{\alpha\delta}^\theta H_{\beta\beta}^\sigma] \\ &+ \frac{K_{\sigma\delta}}{2} [-D_\alpha H_{\gamma\beta}^\sigma + D_\beta H_{\gamma\alpha}^\sigma + H_{\beta\gamma}^\theta H_{\alpha\alpha}^\sigma - H_{\alpha\gamma}^\theta H_{\beta\beta}^\sigma]. \end{aligned} \quad (20)$$

We will investigate now the dual structural equations. Using (17) and (18) we found the dual expressions of  $L_{\alpha\beta\gamma\delta}$  and  $M_{\alpha\beta\gamma\delta}$  as

$$\begin{aligned} \hat{L}_{\alpha\beta\gamma} &= \hat{D}_\alpha g_{\beta\gamma} - \hat{D}_\beta g_{\alpha\gamma} = -H_{\alpha\gamma}^\delta g_{\delta\beta} + H_{\beta\gamma}^\delta g_{\delta\alpha}, \\ \hat{M}_{\alpha\beta\gamma\delta} &= \frac{1}{2} (\hat{L}_{\alpha\beta[\gamma;\delta]} + \hat{L}_{\gamma\delta[\alpha;\beta]}) = -g_{\sigma\alpha} R_{\beta\gamma\delta}^{\prime\sigma} - g_{\sigma\beta} R_{\alpha\delta\gamma}^{\prime\sigma} - g_{\sigma\gamma} R_{\delta\alpha\beta}^{\prime\sigma} - g_{\sigma\delta} R_{\gamma\beta\alpha}^{\prime\sigma} \\ &- g_{\sigma\chi} (H_{\beta\delta}^\chi \hat{G}_{\alpha\gamma}^\sigma - H_{\alpha\delta}^\chi \hat{G}_{\beta\gamma}^\sigma - H_{\beta\gamma}^\chi \hat{G}_{\alpha\delta}^\sigma + H_{\alpha\gamma}^\chi \hat{G}_{\beta\delta}^\sigma) \\ &- H_{\beta\delta}^\sigma g_{\alpha\gamma,\sigma} - H_{\alpha\gamma}^\sigma g_{\beta\delta,\sigma} + H_{\beta\gamma}^\sigma g_{\alpha\delta,\sigma} + H_{\alpha\delta}^\sigma g_{\beta\gamma,\sigma}. \end{aligned} \quad (21)$$

Here  $\hat{G}_{\alpha\delta}^\sigma = H_{\alpha\delta}^\sigma + \hat{\Gamma}_{\alpha\delta}^\sigma$ , the semicolon denotes the dual covariant derivative and  $R_{\nu\rho\sigma}^{\prime\beta}$  has the following form

$$R_{\nu\rho\sigma}^{\prime\beta} = H_{\nu\sigma,\rho}^\beta - H_{\nu\rho,\sigma}^\beta + H_{\nu\sigma}^\alpha H_{\alpha\rho}^\beta - H_{\nu\rho}^\alpha H_{\alpha\sigma}^\beta. \quad (22)$$

From (22) we conclude that  $R_{\nu\rho\sigma}^{\prime\beta}$  looks like as the curvature tensor  $R_{\nu\rho\sigma}^\beta$  [9].

Taking into account (17) we found a new identity for  $K_{\alpha\beta}$

$$K_\beta^\sigma D_\sigma K_{\nu\lambda} + K_\lambda^\sigma D_\sigma K_{\beta\nu} + K_\nu^\sigma D_\sigma K_{\beta\lambda} = 0. \quad (23)$$

By duality we get from (23) another identity

$$g_\beta^\sigma \hat{D}_\sigma g_{\nu\lambda} + g_\lambda^\sigma \hat{D}_\sigma g_{\beta\nu} + g_\nu^\sigma \hat{D}_\sigma g_{\beta\lambda} = 0. \quad (24)$$

In the flat space case the general solution of eq. (1) has the form

$$K_{\beta\gamma} = s_{\beta\gamma} + \frac{2}{3} B_{\alpha(\beta\gamma)} x^\alpha + \frac{1}{3} A_{\alpha\beta\gamma\delta} x^\alpha x^\delta. \quad (25)$$

Here  $s_{\beta\gamma}$ ,  $B_{\alpha\beta\gamma}$  and  $A_{\alpha\beta\gamma\delta}$  are constant tensors having the same symmetries as  $K_{\beta\gamma}$ ,  $L_{\alpha\beta\gamma}$  and  $M_{\alpha\beta\gamma\delta}$  respectively. From (19) and (25) we get

$$\begin{aligned} L_{\alpha\beta\gamma} &= H_{\alpha\gamma}^\delta (s_{\delta\beta} + \frac{2}{3} B_{\sigma(\delta\beta)} x^\sigma + \frac{1}{3} A_{\sigma\delta\beta\lambda} x^\sigma x^\lambda) \\ &- H_{\beta\gamma}^\delta (s_{\delta\alpha} + \frac{2}{3} B_{\sigma(\delta\alpha)} x^\sigma + \frac{1}{3} A_{\sigma\delta\alpha\lambda} x^\sigma x^\lambda). \end{aligned} \quad (26)$$

Using (20) and (25) the expression of  $M_{\alpha\beta\gamma\delta}$  becomes

$$M_{\alpha\beta\gamma\delta} = \frac{1}{2} (K_{\sigma\alpha} R_{\beta\gamma\delta}^{\prime\sigma} + K_{\sigma\beta} R_{\alpha\delta\gamma}^{\prime\sigma} + K_{\sigma\gamma} R_{\delta\alpha\beta}^{\prime\sigma} + K_{\sigma\delta} R_{\gamma\beta\alpha}^{\prime\sigma}). \quad (27)$$

Let us suppose now that  $K_{\alpha\beta}$  given by (25) is non-degenerate [10]. In this case we found that

$$\hat{L}_{\alpha\beta\gamma} = -H_{\beta\alpha\gamma} + H_{\alpha\beta\gamma}, \quad (28)$$

and

$$\begin{aligned} \hat{M}_{\alpha\beta\gamma\delta} = & -R'_{\alpha\beta\gamma\delta} - R'_{\beta\alpha\delta\gamma} - R'_{\gamma\delta\alpha\beta} - R'_{\delta\gamma\beta\alpha} \\ & - H_{\beta\delta}^{\sigma} \hat{G}_{\alpha\gamma}^{\sigma} + H_{\alpha\delta}^{\sigma} \hat{G}_{\beta\gamma}^{\sigma} + H_{\beta\gamma}^{\sigma} \hat{G}_{\alpha\delta}^{\sigma} - H_{\alpha\gamma}^{\sigma} \hat{G}_{\beta\delta}^{\sigma}. \end{aligned} \quad (29)$$

#### 4. CONCLUSIONS

The geometric duality between local geometries described by  $g_{\alpha\beta}$  and by  $K_{\alpha\beta}$  was presented. We found the relation between connections corresponding to  $g_{\alpha\beta}$  and  $K_{\alpha\beta}$  respectively and we have shown that the dual transformation is not a conformal transformation. The manifold and its dual have the same isometries if  $D_{\lambda}K_{\alpha\beta} = 0$ . We have shown that  $L_{\alpha\beta\gamma}$  looks like an angular momentum. The dual structural equations were analyzed and the expressions of  $\hat{L}_{\alpha\beta\gamma}$  and  $\hat{M}_{\alpha\beta\gamma\delta}$  were calculated. In the flat space case the general forms of  $(L_{\alpha\beta\gamma}, \hat{L}_{\alpha\beta\gamma})$  and  $(M_{\alpha\beta\gamma\delta}, \hat{M}_{\alpha\beta\gamma\delta})$  were found.

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