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Fuhai Zhu; Zhiqi Chen
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# NOVIKOV SUPERALGEBRAS WITH $A_{0}=A_{1} A_{1}$ 

Fuhai Zhu, Zhiqi Chen, Tianjin
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#### Abstract

Novikov superalgebras are related to quadratic conformal superalgebras which correspond to the Hamiltonian pairs and play a fundamental role in completely integrable systems. In this note we show that the Novikov superalgebras with $A_{0}=A_{1} A_{1}$ and $\operatorname{dim} A_{1}=2$ are of type $N$ and give a class of Novikov superalgebras of type $S$ with $A_{0}=$ $A_{1} A_{1}$.


Keywords: Novikov algebra, Novikov superalgebra, type $N$, type $S$
MSC 2010: 17A70, 17A30

## 1. Introduction

Novikov superalgebras are a super variant of Novikov algebras. They are closely related to popular algebraic objects such as conformal superalgebras [5], vertex operator superalgebras [8] and super Gel'fand-Dorfman bialgebras [7] which play an important role in the quantum field theory and the theory of completely integrable systems.

A Novikov superalgebra $A$ is a $\mathbb{Z}_{2}$-graded vector space $A=A_{0}+A_{1}$ with a bilinear product $(u, v) \mapsto u v$ for any $u \in A_{i}, v \in A_{j}, w \in A$ satisfying

$$
\begin{gather*}
(u v) w-u(v w)=(-1)^{i j}((v u) w-v(u w)),  \tag{1.1}\\
(w u) v=(-1)^{i j}(w v) u . \tag{1.2}
\end{gather*}
$$

The even part of a given Novikov superalgebra is what is said to be a Novikov algebra introduced in connection with the Poisson brackets of hydrodynamic type [1] and Hamiltonian operators in the formal variational calculus [2], [3], [4], [9], [10].

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Novikov superalgebras are classified into two types: $N$ and $S$. Let $A=A_{0}+A_{1}$ be a Novikov superalgebra with multiplication $(u, v) \mapsto u v$. If $A$ is also a Novikov algebra with respect to the same product and with superstructure forgotten, then $A$ is called a Novikov superalgebra of type $N$, otherwise $A$ is said to be of type $S$. It is proved in [6] that all Novikov superalgebras of dimensions up to 3 are of type $N$. Up to now, all examples of Novikov superalgebras have been of type $N$.

In this paper we show that the Novikov superalgebras with $A_{0}=A_{1} A_{1}$ and $\operatorname{dim} A_{1}=2$ are of type $N$. Furthermore, we provide a class of Novikov superalgebras of type $S$ with $A_{0}=A_{1} A_{1}$ and $\operatorname{dim} A=2 n+1$ for $n \geqslant 2$.

Throughout the paper we assume that the algebras are finite-dimensional over $\mathbb{C}$. Obvious proofs are omitted.

## 2. Novikov superalgebras with $A_{0}=A_{1} A_{1}$

Let $A=A_{0}+A_{1}$ be a Novikov superalgebra.
Lemma 2.1. The subspace $A_{1} A_{1}+A_{1}$ is an ideal of $A$.
Proof. For any $x \in A_{0}, y, z \in A_{1}$, we have

$$
\begin{gathered}
(y z) x=(y x) z \in A_{1} A_{1}, \\
x(y z)=y(x z)+(x y) z-(y x) z \in A_{1} A_{1} .
\end{gathered}
$$

It follows that $A_{1} A_{1}+A_{1}$ is an ideal of $A$.

Lemma 2.2. For any $x \in A_{1}, y \in A$, we have $(x x) y=x(x y)$ and $(y x) x=0$.
Now we consider the Novikov superalgebras $A=A_{0}+A_{1}$ with $A_{0}=A_{1} A_{1}$, $\operatorname{dim} A_{1}=2$ and $A_{0} \neq\{0\}$. Then one can easily see that $\operatorname{dim} A_{0} \leqslant 4$.

Let $f_{1}, f_{2}$ be a basis of $A_{1}$. Set

$$
\begin{equation*}
e_{1}=f_{1} f_{1}, \quad e_{2}=f_{2} f_{2}, \quad e_{3}=f_{1} f_{2}, \quad e_{4}=f_{2} f_{1} \tag{2.1}
\end{equation*}
$$

Then $A_{0}=\left\langle e_{1}, e_{2}, e_{3}, e_{4}\right\rangle$.
Assume that $x x=0$ for any $x \in A_{1}$. Then $\operatorname{dim} A_{0}=1$ and $f_{1} f_{2}=-f_{2} f_{1}$ is a basis of $A_{0}$. By [6], $A_{0} A_{1}=A_{1} A_{0}=A_{0} A_{0}=0$ and $A$ is of type $N$.

In the following, assume that $e_{1}=f_{1} f_{1} \neq 0$.

Lemma 2.3. $A_{0} A_{1}=0$.

Proof. By Lemma 2.2, one has $f_{1} e_{1}=e_{1} f_{1}=0, e_{4} f_{1}=0, f_{2} e_{2}=e_{2} f_{2}=0$ and $e_{3} f_{2}=0$. The following is to show that $e_{2} f_{1}=0$.

If $e_{2}$ is a multiple of $e_{1}$, the assertion is trivial. Assume that $e_{1}$ and $e_{2}$ are linearly independent. It is easy to see that

$$
\begin{equation*}
\left(e_{2} f_{1}\right) f_{1}=\left(e_{2} f_{1}\right) f_{2}=0 \tag{2.2}
\end{equation*}
$$

Furthermore,

$$
\begin{align*}
f_{2}\left(e_{2} f_{1}\right)=f_{2}\left(f_{2}\left(f_{2} f_{1}\right)\right) & =\left(f_{2} f_{2}\right)\left(f_{2} f_{1}\right)=\left(f_{2}\left(f_{2} f_{1}\right)\right) f_{2}  \tag{2.3}\\
& =-\left(\left(f_{2} f_{2}\right) f_{2}\right) f_{1}=0 .
\end{align*}
$$

Assume that $e_{2} f_{1}=a f_{1}+b f_{2}$, then by eqs. (2.2) and (2.3) we have that

$$
a e_{1}+b e_{4}=a e_{3}+b e_{2}=a e_{4}+b e_{2}=0
$$

It follows that $a=b=0$. Similarly, $e_{3} f_{1}=e_{4} f_{1}=0$ and $A_{0} f_{2}=0$.

Lemma 2.4. The subalgebra $A_{0}$ is skew-commutative.
Proof. For any $f_{i}, f_{j}, e_{k}$, we have

$$
\left(f_{i} f_{j}\right) e_{k}=\left(f_{i} e_{k}\right) f_{j}=\left(e_{k} f_{i}\right) f_{j}+f_{i}\left(e_{k} f_{j}\right)-e_{k}\left(f_{i} f_{j}\right)=-e_{k}\left(f_{i} f_{j}\right)
$$

since $e_{k} f_{i}=e_{k} f_{j}=0$ by Lemma 2.3.

## Lemma 2.5.

1) $f_{1} e_{2}=-f_{2} e_{3}, f_{1} e_{4}=-f_{2} e_{1}, f_{1} e_{1}=f_{1} e_{3}=0, f_{2} e_{2}=f_{2} e_{4}=0$.
2) $e_{1} e_{2}=e_{3} e_{4}, e_{1} e_{3}=e_{2} e_{4}=0$.

Proof. 1) By eq. (1.1), $\left(f_{1} f_{2}\right) f_{2}-f_{1}\left(f_{2} f_{2}\right)=-\left(f_{2} f_{1}\right) f_{2}+f_{2}\left(f_{1} f_{2}\right)$. So $f_{1} e_{2}=$ $-f_{2} e_{3}$. Similarly, $f_{1} e_{4}=-f_{2} e_{1}$. Also, one has $f_{1} e_{3}=f_{1}\left(f_{1} f_{2}\right)=\left(f_{1} f_{1}\right) f_{2}=0$. Similarly, $f_{2} e_{2}=f_{2} e_{4}=0$.
2) By eq. (1.2), $e_{1} e_{2}=\left(f_{1} f_{1}\right) e_{2}=\left(f_{1} e_{2}\right) f_{1}=-\left(f_{2} e_{3}\right) f_{1}=-\left(f_{2} f_{1}\right) e_{3}=-e_{4} e_{3}=$ $e_{3} e_{4}$. By 1) $e_{1} e_{3}=\left(f_{1} f_{1}\right) e_{3}=f_{1}\left(f_{1} e_{3}\right)=0$ and similarly $e_{2} e_{4}=0$ and $f_{1} e_{2} \neq 0$.

## Lemma 2.6.

1) The elements $f_{1} e_{2}$ and $f_{1} e_{4}$ commute with $f_{1}$ and $f_{2}$.
2) $A_{1} A_{0}=0$ and $A_{0} A_{0}=0$.

Proof. 1) For $x=e_{2}$ or $e_{4}, f_{1}\left(f_{1} x\right)=\left(f_{1} f_{1}\right) x=\left(f_{1} x\right) f_{1}$ and similarly $f_{2}$ commutes with $f_{2} y$ for $y=e_{1}$ or $e_{3}$. Then the assertion follows from Lemma 2.5.
2) If $A_{1} A_{0} \neq 0$, then assume that $f_{1} e_{2}=a_{1} f_{1}+a_{2} f_{2}$ and $f_{1} e_{4}=a_{3} f_{1}+a_{4} f_{2}$, where at least one $a_{i}$ is not zero. Then by 1$), a_{i} e_{3}=a_{i} e_{4}$ for $1 \leqslant i \leqslant 4$, so $e_{3}=e_{4}$. Consequently, $f_{1} e_{4}=f_{1} e_{3}=0$ and $f_{1} e_{2}=-f_{2} e_{3}=-f_{2} e_{4}=0$. So $A_{1} A_{0}=0$. Moreover, $A_{0} A_{0}=0$.

Theorem 2.7. Let $A=A_{0}+A_{1}$ be a Novikov superalgebra with $A_{0}=A_{1} A_{1}$ and $\operatorname{dim} A_{1}=2$. Then $A$ is of type $N$.

Example 2.8. Assume that $A=A_{0}+A_{1}$ is a vector space of dimension $2 n+1$, where $n \geqslant 2$. Let $e_{1}, \ldots, e_{n}$ be a basis of $A_{0}$ and $f_{1}, \ldots, f_{n}, f_{n+1}$ a basis of $A_{1}$. Set

$$
\begin{aligned}
f_{1} f_{i} & =e_{i}, \quad 1 \leqslant i \leqslant n, \\
e_{i} f_{n+1-i} & = \begin{cases}f_{n+1}, & 1 \leqslant i \leqslant\left[\frac{1}{2} n\right], \\
-f_{n+1}, & {\left[\frac{1}{2}(n+1)\right]+1 \leqslant i \leqslant n,}\end{cases} \\
f_{i} e_{n+1-i} & = \begin{cases}f_{n+1}, & 1 \leqslant i \leqslant\left[\frac{1}{2} n\right], \\
-f_{n+1}, & {\left[\frac{1}{2}(n+1)\right]+1 \leqslant i \leqslant n,}\end{cases}
\end{aligned}
$$

Then we have

$$
\begin{gathered}
\left(f_{1} f_{j}\right) f_{n+1-j}=-\left(f_{1} f_{n+1-j}\right) f_{j} \\
\left(f_{1} f_{1}\right) f_{n}=f_{1}\left(f_{1} f_{n}\right) \\
\left(f_{1} f_{j}\right) f_{n+1-j}-f_{1}\left(f_{j} f_{n+1-j}\right)=-\left(f_{j} f_{1}\right) f_{n+1-j}+f_{j}\left(f_{1} f_{n+1-j}\right), \quad j \geqslant 2
\end{gathered}
$$

It follows that $A$ is a Novikov superalgebra of type $S$ with $A_{0}=A_{1} A_{1}$.

## References

[1] A. A. Balinskii, S. P. Novikov: Poisson brackets of hydrodynamic type, Frobenius algebras and Lie algebras. Sov. Math. Dokl. 32 (1985), 228-231.
[2] I. M. Gel'fand, I. Ya. Dorfman: Hamiltonian operators and algebraic structures related to them. Funct. Anal. Appl. 13 (1980), 248-262.
[3] I. M. Gel'fand, I. Y. Dorfman: The Schouten brackets and Hamiltonian operators. Funct. Anal. Appl. 14 (1981), 223-226.
[4] I. M. Gel'fand, I. Y. Dorfman: Hamiltonian operators and infinite-dimensional Lie algebras. Funct. Anal. Appl. 15 (1982), 173-187.
[5] V. G. Kac: Vertex Algebras for Beginners. University Lecture Series, 10. American Mathematical Society (AMS), Providence, 1998.
[6] Y. F. Kang, Z. Q. Chen: Novikov superalgebras in low dimensions. J. Nonlinear Math. Phys 16 (2009), 251-257.
[7] X. P. Xu: Quadratic conformal superalgebras. J. Algebra 231 (2000), 1-38.
[8] X. P. Xu: Introduction to Vertex Operator Superalgebras and Their Modules. Kluwer, Dordercht, 1998.
[9] X. P. Xu: Hamiltonian operators and associative algebras with a derivation. Lett. Math. Phys. 33 (1995), 1-6.
[10] X. P. Xu: Hamiltonian superoperators. J. Phys A. Math. Gen. 28 (1995), 1681-1698.

Authors' addresses: F. Z h u, School of Mathematical Sciences and LPMC, Nankai University, Tianjin 300071, P.R.China; Z. Chen (corresponding author), School of Mathematical Sciences and LPMC, Nankai University, Tianjin 300071, P. R. China, e-mail: chenzhiqi@nankai.edu.cn.

