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# Properties of Operators Related to the Boltzmann Collision Term 

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Some recent results on boundedness and compactness of certain operators closely related to the Boltzmann collision term are reviewed.

Je dán přehled některých nových výsledků o omezenosti a kompaktnosti jistých operátorủ úzce souvisejícich s Boltzmannovým srážkovým členem.

Сведены некоторые новейшие результаты об ограниченности и компактности определенных операторов находящихся в тесной связи со столкновительным членом Больцмана.

## Introduction

The Boltzmann collision term $Q$, characterizing the effect of collisions between two sorts of particles with velocity distribution functions $f$ and $g$ upon the change of $f$, has the form

$$
\begin{align*}
Q[f, g](v)= & \int_{\left.\left.R^{3} \times<0, \pi / 2\right) \times<0,2 \pi\right)}\left[f\left(v^{\prime}\right) g\left(v_{*}^{\prime}\right)-f(v) g\left(v_{*}\right)\right] \times  \tag{1}\\
& \times B(\vartheta,\|V\|) \mathrm{d} v_{*} \mathrm{~d} \vartheta \mathrm{~d} \varepsilon\left(v \in \mathrm{IR}^{3}\right),
\end{align*}
$$

provided only binary elastic collisions are admitted. Here,

$$
\begin{gathered}
V=v-v_{*}, \\
v^{\prime}=v-\frac{2}{1+x}(V \cdot n) V, \\
v_{*}^{\prime}=v_{*}+\frac{2 x}{1+x}(V . n) V, \\
n=\left(\begin{array}{c}
\sin \vartheta \cos \varepsilon \\
\sin \vartheta \sin \varepsilon \\
\cos \vartheta
\end{array}\right),
\end{gathered}
$$

and $x$ is the ratio of the masses of the colliding particles;

$$
B:\left\langle 0, \frac{1}{2} \pi\right) \times(0, \infty) \rightarrow\langle 0, \infty)
$$

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is a measurable function closely related to the differential cross section [1]. In case that the collisions are governed by a repulsive force inversely proportional to some power, say, $k$-th $(k>2)$, of the distance between the (two) colliding particles, the function $B$ is factorized as follows [1]:

$$
\begin{equation*}
B, \vartheta, u)=u^{\gamma} b(\vartheta)\left(0 \leqq \vartheta<\frac{1}{2} \pi, u>0\right) \tag{2}
\end{equation*}
$$

where

$$
\gamma=\frac{k-5}{k-1}
$$

(hence, $-3<\gamma<1$ ). If an angular cut-off [1], [2] is introduced, the function $b$ becomes integrable over $\left\langle 0, \frac{1}{2} \pi\right.$ ). Note that the well-known hard sphere model of interactions gives the function $B$ of the form (2) with $\gamma=1$.

Throughout this note we shall deal with a more general version of (1), namely

$$
\begin{gathered}
\left.Q[f, g](v)=\int_{R^{3} \times\langle 0, \pi / 2 \times<0,2 \pi)}\left[f\left(v^{\prime}\right) g\left(v_{*}^{\prime}\right)-f^{\prime} v\right) g\left(v_{*}\right)\right] \times \\
\times B(\vartheta,\|V\|) \exp \left(-\lambda\left\|v_{*}\right\|^{2}\right) \mathrm{d} v_{*} \mathrm{~d} \vartheta \mathrm{~d} \varepsilon \quad\left(v \in \mathbb{R}^{3}\right)
\end{gathered}
$$

where $\lambda \geqq 0$ is a parameter. A term of this type appears in an equation arising from the Boltzmann kinetic equation by means of a cetrain transformation [3]. Define

$$
\left.v[g](v)=2 \pi \int_{R^{3} \times\langle 0, \pi / 2)} g^{\prime} v_{*}\right) B(\vartheta,\|V\|) \exp \left(-\lambda\left\|v_{*}\right\|^{2}\right) \mathrm{d} v_{*} \mathrm{~d} \vartheta \quad\left(v \in \mathbb{R}^{3}\right)
$$

Given functions $f$ and $g$ such that $Q[f, g]$ exists and $v[g]$ is finite on $\mathbb{R}^{\mathbf{3}}$, we can decompose $Q$ as follows:

$$
Q[f, g]=Q^{+}[f, g]-f . v[g]
$$

where

$$
\begin{gathered}
Q^{+}[f, g](v)=\int_{R^{3} \times\langle 0, \pi / 2) \times\langle 0,2 \pi)} f\left(v^{\prime}\right) g\left(v_{*}^{\prime}\right) B(\vartheta,\|V\|) \times \\
\times \exp \left(-\lambda\left\|v_{*}\right\|^{2}\right) \mathrm{d} v_{*} \mathrm{~d} \vartheta \mathrm{~d} \varepsilon \quad\left(v \in \mathbb{R}^{3}\right) .
\end{gathered}
$$

Letting $f$ and $g$ vary within some function space, we may regard $Q^{+}$as a bilinear operator. If one of the functions $f$ and $g$ is fixed, we consider the following linear operators:

$$
Q^{(1)}=Q^{+}[f, \cdot], \quad Q^{(2)}=Q^{+}[\cdot, g], \quad Q^{(3)}=f \cdot v[\cdot]
$$

The purpose of this note is to review criteria on boundedness of the operators $Q^{+}$and $Q^{(1)}, Q^{(2)}, Q^{(3)}$, as well as criteria on compactness of the last three operators in quotes. The function $B$ will correspond to an interaction model from a specified class involving, inter alia, the models of inverse-power repulsive interparticle forces with angular cut-off and the hard sphere model. The subject matter presented was treated in full detail in recent author's papers [3], [4], which extend former results due to Carleman [5], Grad [2], [6], Dorfman [7], and Molinet [8].

## Notation and agreements

We shall work in spaces of the $L_{p}$-type. The space $L_{p}\left(\mathbb{R}^{3}\right)$ will be denoted just $L_{p}$ and the norm $\|\cdot\|_{p}$. For $\sigma \geqq 1$ we consider the space

$$
L_{p(\sigma)}=\left\{h ;\left(1+\|v\|^{2}\right)^{\sigma / 2} h(v) \in L_{p}\right\}
$$

endowed with the norm

$$
\|h\|_{p(\sigma)}=\left\|\left(1+\|v\|^{2}\right)^{\sigma / 2} h(v)\right\|_{p} .
$$

For easy reference recall that $p$ is conjugated with $q, 1 \leqq q \leqq \infty$, if
(3) $\frac{1}{p}+\frac{1}{q}=1$ for $1<q<\infty, p=1$ for $q=\infty, p=\infty$ for $q=1$.

Given $\gamma$, define $\mathscr{P}_{\gamma}$ to be the set of measurable functions $B(\vartheta, u)$ satisfying

$$
0 \leqq B(\vartheta, u) \leqq \text { const } \cdot u^{\gamma} \sin \vartheta \cos ^{\gamma} \vartheta \quad\left(0 \leqq \vartheta<\frac{1}{2} \pi, u>0\right),
$$

the indicated multiplicative constant depending on $B$.
Given $\alpha$, define $\mathscr{M}_{\alpha}$ to be the set of measurable functions $h(v)$ satisfying

$$
\left|h^{\prime}(v)\right| \leqq \text { const } \cdot \exp \left(-\alpha\|v\|^{2}\right) \quad\left(v \in \mathbb{R}^{3}\right)
$$

the indicated multiplicative constant depending on $h$.
In any discussion involving $x$ it will be tacitly assumed that $x$ is an arbitrary positive parametre unless its value is explicitly specified.

## Results concerning the bilinear operator $Q^{+}[\cdot, \cdot]$

Theorem (see [4], 2.1). Let

$$
\lambda>0 ;
$$

let $B(\vartheta, u)$ be a (nonnegative measurable) function such that

$$
\left.B(\vartheta, u) \leqq u^{\nu} b^{\prime} \vartheta\right) \quad\left(0 \leqq \vartheta<\frac{1}{2} \pi, u>0\right)
$$

with

$$
-3<\gamma \leqq 0
$$

and let

$$
\left.b \in L_{q}^{\prime} 0, \frac{1}{2} \pi\right)
$$

for some $q$ satisfying

$$
\begin{array}{ll}
1 \leqq q<\frac{3}{|\gamma|} & \text { if } \quad \gamma \neq 0 \\
1 \leqq q \leqq \infty & \text { if } \tag{4}
\end{array} \quad \gamma=0
$$

Defining $p$ by (3), $Q^{+}[\cdot, \cdot]$ is a bounded bilinear operator from $L_{p} \times L_{p}$ into $L_{p}$.

Corollary (see [4], 2.3). Let

$$
\lambda>0
$$

and

$$
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 0 .
$$

Let $q$ satisfy $1 \leqq q<1 /|\gamma|$ if $\gamma \neq 0$, or (4). Then, defining $p$ by (3), $Q^{+}[\cdot, \cdot]$ is a bounded bilinear operator from $L_{p} \times L_{p}$ into $L_{p}$.

Theorem (see [4], 3.5). Let

$$
\lambda>0
$$

and

$$
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 0
$$

Let $q$ and $\sigma$ satisfy

$$
\begin{gathered}
1 \leqq q \leqq 2+|\gamma| \\
1 \leqq \sigma \leqq \frac{|\gamma|}{q}+\min \left\{1, \frac{2}{q}\right\}
\end{gathered}
$$

Then, defining $p$ by (3), $Q^{+}[\cdot, \cdot]$ is a bounded bilinear operator from $L_{p(\sigma)} \times$ $\times L_{p(\sigma)}$ into $L_{p(\sigma)}$.

Theorem (see [4], 3.6). Let

$$
\begin{aligned}
& \lambda>0, \\
& x=1,
\end{aligned}
$$

and

$$
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 1
$$

Then for every $1 \leqq p \leqq \infty$ and every $1 \leqq \sigma(<\infty), Q^{+}[\cdot, \cdot]$ is a bounded bilinear operator from $L_{p(\sigma)} \times L_{p(\sigma)}$ into $L_{p(\sigma)}$.

## Results concerning the linear operators $Q^{(i)}(i=1,2,3)$

Some criteria on boundedness of $Q^{(1)}$ and $Q^{(2)}$ arise immediately from those of $Q^{+}[\cdot, \cdot]$. If the (fixed) functions $f$ and $g$ are majorized by a Maxwellian, more conclusive results can be established, as seen below.

Theorem (see [3], 5.7). Let

$$
\begin{gathered}
\lambda=0, \\
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 0,
\end{gathered}
$$

and

$$
f, g \in \mathscr{M}_{\alpha}, \quad \alpha>0
$$

Then each of $Q^{(i)}(i=1,2,3)$ is a bounded linear operator from $L_{1}$ into itself.

Theorem (see [3], 5.7*)). Let

$$
\begin{gathered}
\lambda=0, \\
x \neq 1, \\
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 0,
\end{gathered}
$$

and

$$
g \in \mathscr{M}_{\alpha}, \quad \alpha>0
$$

Then for every $1 \leqq p \leqq \infty, Q^{(2)}$ is a bounded linear operator from $L_{p}$ into itself.
Theorem (see [3], 6.9). Let

$$
\begin{gathered}
\lambda>0, \\
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 1,
\end{gathered}
$$

and

$$
f, g \in \mathscr{M}_{\alpha}, \quad \alpha>0
$$

Then for every $1 \leqq p \leqq \infty$, each of $Q^{(i)}(i=1,2,3)$ is a bounded linear operator from $L_{p}$ into itself.

Theorem (see [3], 6.16). Let

$$
\begin{gathered}
\lambda>0, \\
B \in \mathscr{B}_{\gamma}, \quad-1<\gamma \leqq 1,
\end{gathered}
$$

and

$$
f, g \in \mathscr{M}_{\alpha}, \quad \alpha>0
$$

Then each of $Q^{(i)}(i=1,2,3)$ is a compact operator from $L_{2}$ into itself.

## Concluding remarks

The results reviewed above may be found interesting for mathematicians and physicists engaged in the kinetic theory. The results originate from author's thesis, which provides mathematical background for investigating transport processes in ionized gases in an external electromagnetic field. A problem of this sort is considered, e.g., in connection with determining the mobility of charged particles in operation media of electrotechnical devices under high-voltage stresses. The thesis was worked out at the Mathematical Institute of the Charles University, Prague.

## References

[1] Cercignani C.: Theory and application of the Boltzmann equation. Scott. Acad. Press, Edinburgh-London 1975.

[^0][2] Grad H.: Asymptotic theory of the Boltzmann equation, II. Rarefied gas dynamics, vol. 1 (J. A. Laurmann, ed.). Acad. Press, New York-London 1963, 26-59.
[3] Chvála F.: On the Boltzmann collision term and related operators. Acta Techn. ČSAV 30 (1985), 327-364.
[4] Chvála F.: On boundedness of a bilinear operator derived from the Boltzmann collision term. To appear in Acta Techn. ČSAV.
[5] Carleman T.: Problèmes mathématiques dans la théorie cinétique des gaz. Almqvist Wiksells, Uppsala 1957.
[6] Grad H.: Asymptotic equivalence of the Navier-Stokes and nonlinear Boltzmann equations. Proc. Symp. Appl. Math., vol. 17 (R. Finn, ed.), Amer. Math. Soc., Providence 1965, 154-183.
[7] Dorfman R.: Note on the linearized Boltzmann integral equation for rigid sphere molecules. Proc. Nat. Acad. Sci. USA 50 (1963), 804-806.
[8] Molinet F. A.: Existence, uniqueness and properties of the solutions of the Boltzmann kinetic equation for a weakly ionized gas. J. Math. Phys. 18 (1977), 984-996.


[^0]:    ${ }^{*}$ ) The clause "provided $x \neq 1$ " should be attached to the last sentence of Theorem 5.7 in [3].

