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# Two Models of Non-Euclidean Spaces Generated by Associative Algebras 

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#### Abstract

We present a nontrivial example how to generate non-Euclidean geometries from associative unital algebras. We consider bundles of the sphere of the degenerate non-Euclidean space and its two models. The first (conformal) model is obtained by the mapping $S$ onto a plane passing through the origin. It is analogous to the stereographic mapping. The second model (projective) is constructed by the Norden normalization method, where we project the sphere onto a plane of normalization defining the metric and Christoffel symbols which allow us to find geodesic curves.


## 1. Introduction

A lot of models of non-Euclidean spaces were studied in the past, especially spaces of a constant curvature, projective spaces and the conformal planes (e.g. [10], [11], [12], [19]). There exists a lot of studies on how these models can be generated by algebras. It is well known that algebras define some structures in bundle manifolds of different types (e.g. [5], [9], [13]). In the literature, we can find many applications of this approach on the cases of non-Eucledian spaces (e.g. [4], [6], [16], [17], [20]).

We would like to present non-standard models within this framework. In the preliminaries we describe how an associative algebra generates a vector space and we also discuss some of its properties. In the next section we define a sphere and the map $S$ in this vector space and we use it to construct a conformal model. In the last section we remind the reader of some facts about the Norden normalization method [7] and we use it for the construction of a projective model.

[^0]Foundations of the theory of finite-dimensional associative algebras were made by E. Cartan (1898), Wedderburn (1908) and F. E. Molin (1983), who described the structure of any algebra over an arbitrary base field [2]. E. Study and E. Cartan in [15] classified all 3 and 4-dimensional unital associative irreducible ${ }^{1}$ algebras up to an isomorphism. This classification can be also found in [18]. In this paper we consider only one type of 3-dimensional algebra $\mathfrak{A}$.

## 2. Preliminaries

Let $\mathfrak{A}$ be an unital associative 3 -dimensional algebra and $\left\{\mathbf{1}, e_{1}, e_{2}\right\}$ be its basis with the identity element $\mathbf{1}$. The multiplication rules are:

$$
\begin{equation*}
\left(e_{1}\right)^{2}=1, \quad\left(e_{2}\right)^{2}=0, \quad e_{1} e_{2}=-e_{2} e_{1}=e_{2} \tag{1}
\end{equation*}
$$

The algebra $\mathfrak{A}$ is the set of upper triangular matrices

$$
\begin{gather*}
\left(\begin{array}{cc}
x_{0} & x_{2} \\
0 & x_{1}
\end{array}\right)=x_{0} \cdot \mathbf{1}+x_{1} \cdot e_{1}+x_{2} \cdot e_{2}, \text { where } \\
\mathbf{1}=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), \quad e_{1}=\left(\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right), \quad e_{2}=\left(\begin{array}{ll}
0 & 1 \\
0 & 0
\end{array}\right) \tag{2}
\end{gather*}
$$

are basic elements [2].
The algebra $\mathfrak{A}$ admits the following conjugation

$$
x=x_{0}+x_{1} e_{1}+x_{2} e_{2} \quad \rightarrow \quad \bar{x}=x_{0}-x_{1} e_{1}-x_{2} e_{2}
$$

with the property $\overline{x y}=\bar{y} \bar{x}$.
We consider the bilinear form $(x, y)$ which takes real values and determines a degenerate scalar product:

$$
\begin{equation*}
(x, y)=\frac{1}{2}(x \bar{y}+y \bar{x})=x_{0} y_{0}-x_{1} y_{1} . \tag{3}
\end{equation*}
$$

It defines the structure of a degenerate pseudo-Euclidean vector space of rank 2 on $\mathfrak{A}$. (It is also possible to call this space "semi-pseudo-Euclidean", but later we will call it just "pseudo-Euclidean".) The set of invertible elements $G=\{x \in \mathfrak{U} \mid$ $\left.\mid\left(x_{0}\right)^{2}-\left(x_{1}\right)^{2} \neq 0\right\}$ is a non-Abelian Lie group with the same multiplication rule ([1]). Its underlying manifold is obtained from $\mathbb{R}^{3}$ by removing two transversal planes, hence it consists of 4 connected components.

The distance is defined as usual, $d(x, y)^{2}=(x-y, x-y)$. The geodesic curves $x(t)$ are then

$$
x_{0}=a_{0} t+b_{0} \quad x_{1}=a_{1} t+b_{1} \quad x_{2}=f(t)
$$

where $f(t)$ is an arbitrary function of $t$ and $a_{0}, a_{1}, b_{0}, b_{1}$ are the numerical coefficients.
In the basis (2) we can find two subalgebras: $R\left(e_{1}\right)$ with basis $\left\{\mathbf{1}, e_{1}\right\}$, it is an algebra of double numbers, and a subalgebra $R\left(e_{2}\right)$ with basis $\left\{\mathbf{1}, e_{2}\right\}$, it is an algebra of dual

[^1]numbers. The set of their invertible elements $H_{1}=\left\{x_{0}+x_{1} e_{1} \in R\left(e_{1}\right) \mid x_{0}^{2}-x_{1}^{2} \neq 0\right\}$ and $H_{2}=\left\{x_{0}+x_{2} e_{2} \in R\left(e_{2}\right) \mid x_{0} \neq 0\right\}$ are Lie subgroups of the Lie group $G$.

The space of right cosets $H_{1} x$ defines a trivial bundle $\left(G, \pi, M=G / H_{1}\right)$ over the real line $\mathbb{R}$ with the structure group $H_{1}$, where $\pi$ is a canonical projection ([3]). The fiber is a plane without two transversal lines and the structure group is $H_{1}$. The manifold of the group $G$ is diffeomorphic to direct sum $\mathbb{R} \times H_{1}$. The coordinate view of the canonical projection $\pi$ is:

$$
\begin{equation*}
\pi(x)=\frac{x_{2}}{x_{0}-x_{1}} \tag{4}
\end{equation*}
$$

The equation of fibers is:

$$
\begin{equation*}
u\left(x_{0}-x_{1}\right)-x_{2}=0, u \in \mathbb{R} \tag{5}
\end{equation*}
$$

Let us investigate the isometry group of the pseudo-Euclidean space $G$. We can easily find that it has no dilations and inversions while there is a vertical translation $x \rightarrow x+a, a \in G$. Furthermore, the isometry group includes rotations, resp. antirotations,

$$
x^{\prime}=a x \quad \text { or } \quad x^{\prime}=x a
$$

with $|a|^{2}=1$, resp. $|a|^{2}=-1$. These elements can be represented as:

$$
a=\cosh \varphi \pm \sinh \varphi e_{1}+u \sinh \varphi e_{2}
$$

$$
\text { resp. } \quad a=\sinh \varphi \pm \cosh \varphi e_{1}+u \cosh \varphi e_{2},
$$

where $u \in \mathbb{R}$. The anti-rotations map the elements with the positive norms into the elements with the negative norms and visa versa.

The bilinear form (3) in the algebra $\mathfrak{A}$ takes the real values, therefore it is possible to present it as: $(x, y)=\frac{1}{2}(x \bar{y}+y \bar{x})=\frac{1}{2}(\bar{x} y+\bar{y} x)$. Consequently, in the case of rotations the hyperbolic cosine of an angle between $x$ and $x^{\prime}$ is equal to

$$
\begin{gather*}
\cosh \left(x, x^{\prime}\right)=\frac{(x, a x)}{|x||a x|}=\frac{1 / 2(x \overline{a x}+a x \bar{x})}{|x|^{2}}= \\
\frac{1 / 2(x \bar{x} \bar{a}+a x \bar{x})}{|x|^{2}}=\frac{1}{2}(\bar{a}+a)=\cosh \varphi, \tag{6}
\end{gather*}
$$

and the same for the right multiplication. Similarly we get $\sinh \varphi$ for anti-rotations. Note that the angle $\varphi$ does not depend on $x$.

Isometries

$$
\begin{equation*}
x^{\prime}=a x b \tag{7}
\end{equation*}
$$

where $|a|^{2}= \pm 1,|b|^{2}= \pm 1$, are compositions of rotations and/or anti-rotations $x^{\prime}=a x$ and $x^{\prime}=x b$. We see that (7) defines proper rotations and anti-rotations.

Similarly,

$$
\begin{equation*}
x^{\prime}=a \bar{x} b \tag{8}
\end{equation*}
$$

are compositions of the reflection $x^{\prime}=\bar{x}$ and transformations (7). These are improper rotations and anti-rotations.

Proposition. Any proper or improper rotation/anti-rotation of the pseudoEuclidean space $G$ can be represented by (7) or (8).

Proof. Rotations and anti-rotations (7), (8) are compositions of odd and even numbers of reflections of planes passing through the origin. There corresponds an orthonormal vector $n$ to each plane. If vectors $x_{1}$ and $n$ are collinear, then $\bar{x}_{1} n=\bar{n} x_{1}$ and $x_{1}^{\prime}=-n \bar{x}_{1} n=-n \bar{n} x_{1}=-x_{1}$. If vectors $x_{2}$ and $n$ are orthogonal, then $\bar{x}_{2} n+\bar{n} x_{2}=0$ and $x_{2}^{\prime}=-n \bar{x}_{2} n=n \bar{n} x_{2}=x_{2}$. On the other hand, any vector $x$ can be represented by a sum of vectors $x_{1}$ and $x_{2}$. It means, that a reflection of the plane is: $x^{\prime}=-n \bar{x} n$. Therefore, the composition of even, resp. odd number of reflections of planes are isometries (7), resp. (8).

Corollary. Only translations, rotations and anti-rotations are isometries of $G$.
They all can be written in a known form (for further discussion see e.g. [19])

$$
\left\{\begin{array}{l}
x_{0}^{\prime}=x_{0} \cosh \varphi+x_{1} \sinh \varphi+a_{0}  \tag{9}\\
x_{1}^{\prime}=x_{1} \cosh \varphi+x_{0} \sinh \varphi+a_{1} \\
x_{2}^{\prime}=u_{0} x_{0}+u_{1} x_{1}+u_{2} x_{2}+a_{2}
\end{array}\right.
$$

where $a=a_{i} e_{i} \in G$ and $u_{i} \in \mathbb{R}$.
Let us introduce adapted coordinates $(u, \lambda, \varphi)$ of the bundle in semi-Euclidean space, here $u$ is a basic coordinate, $\lambda, \varphi$ are fiber coordinates. If $|x|^{2}>0$, we denote $\lambda= \pm \sqrt{x_{0}^{2}-x_{1}^{2}} \neq 0$, the sign of $\lambda$ is equal to the sign of $x_{0}$. The adapted coordinates of the bundle in this case are:

$$
\begin{equation*}
x_{0}=\lambda \cosh \varphi, \quad x_{1}=\lambda \sinh \varphi, \quad x_{2}=u \lambda \exp \varphi, \tag{10}
\end{equation*}
$$

where $\lambda \in \mathbb{R}_{0}, \quad u, \varphi \in \mathbb{R}$.
If $|x|^{2}<0$, then we write $\lambda= \pm \sqrt{x_{1}^{2}-x_{0}^{2}}$, the sign of $\lambda$ is equal to the sign of $x_{1}$ :

$$
\begin{equation*}
x_{0}=\lambda \sinh \varphi, \quad x_{1}=\lambda \cosh \varphi, \quad x_{2}=u \lambda \exp \varphi . \tag{11}
\end{equation*}
$$

The structure group acts as follows:

$$
\begin{equation*}
u^{\prime}=u, \quad \lambda^{\prime}=\lambda \rho, \quad \varphi^{\prime}=\varphi+\psi, \tag{12}
\end{equation*}
$$

where the element $a(0, \rho, \psi)$ of the structure group acts on the element $x(u, \lambda, \varphi) \in G$. This group consists of 4 connected components.

## 3. Conformalmodelof a sphere

Definition. We call semi-Euclidean sphere with an unit radius the set of all elements of algebra $\mathfrak{A}$ whose square is equal to one,

$$
S^{2}(1)=\left\{x \in \mathfrak{A} \mid x_{0}^{2}-x_{1}^{2}=1\right\} .
$$

Analogously, the set of elements with an imaginary unit module $|x|^{2}=-1$ we call semi-Euclidean sphere with an imaginary unit radius $S^{2}(-1)$.

One of these spheres can be obtained from another by rotation. The isometries (9) are now constrained by additional relation $x_{0}^{2}-x_{1}^{2}=1$, therefore, only rotations and vertical translations remain, $a_{0}=a_{1}=0$.

We consider the subbundle of the bundle ( $G, \pi, M=G / H_{1}$ ) of semi-Euclidean sphere $S^{2}(1)$, i.e. the bundle $\pi: S^{2}(1) \rightarrow M$. The fibers of the new bundle are intersections of $S^{2}(1)$ and planes (5). The restriction of the group of double numbers $H_{1}$ to $S^{2}(1)$ is a Lie subgroup $S_{1}$ of double numbers with an unit module

$$
S_{1}=\left\{a_{0}+a_{1} e_{1} \in H_{1} \mid a_{0}^{2}-a_{1}^{2}=1\right\} .
$$

This group consists of two connected components. The bundle $\left(S^{2}(1), \pi, M\right)$ is a trivial bundle of the group $S^{2}(1)$ by the Lie subgroup $S_{1}$ to right cosets.

We define coordinates adapted to the bundle on semi-Euclidean sphere $S^{2}(1)$. If $x \in S^{2}(1)$ then from (10) we get $\lambda=\varepsilon, \varepsilon= \pm 1$. The parametric equation of semiEuclidean sphere in the adapted coordinates $(u, \varphi)$ is:

$$
\begin{equation*}
\mathbf{r}(u, \varphi)=\varepsilon(\cosh \varphi, \sinh \varphi, u \exp \varphi) \tag{13}
\end{equation*}
$$

where $u$ is a basis coordinate, $\varphi$ is a fiber coordinate. Different values of $\varepsilon$ correspond to different connected components of semi-Euclidean sphere $S^{2}(1)$.

Let us define the action of the structure group $S_{1}$ on semi-Euclidean sphere. From (12) and using the adapted coordinates of elements $a\left(0, \varepsilon_{1}, \psi\right), x(u, \varepsilon, \varphi) \in S^{2}(1)$ we get:

$$
u^{\prime}=u, \quad \varepsilon^{\prime}=\varepsilon \varepsilon_{1}, \quad \varphi^{\prime}=\varphi+\psi .
$$

This group also consists of two connected components.
The metric tensor for semi-Euclidean sphere has the matrix representation:

$$
\left(g_{i j}\right)=\left(\begin{array}{cc}
0 & 0 \\
0 & -1
\end{array}\right) .
$$

The linear element of the metric is:

$$
\begin{equation*}
d s_{1}^{2}=-d \varphi^{2} \tag{14}
\end{equation*}
$$

Now, we want to define the conformal model of the bundle $\left(S^{2}(1), \pi, \mathbb{R}\right)$. For that we need to introduce the conformal map of the sphere to a disconnected plane $f: S^{2}(1) \rightarrow Q \in \mathbb{R}^{2} . Q$ is located at $x_{0}=0$. We know that the sphere consists of two disconnected components, one with $x_{0}>0$, and other with $x_{0}<0$. We choose a pole at the first one, $N(1,0,0)$. All points of $S^{2}(1)$ except the line passing through the pole $N$ are stereographically projected to $Q$ such that the first component of the sphere with $x_{0}>0$ is mapped on $\left\{\left(0, x_{1}, x_{2}\right) \mid x_{1} \in(-\infty,-1) \cup(1, \infty), x_{2} \in \mathbb{R}\right\}$ while the second component with $x_{0}<0$ is mapped on the strip $\left\{\left(0, x_{1}, x_{2}\right) \mid x_{1} \in(-1,1), x_{2} \in \mathbb{R}\right\}$. We denote $x, y$ coordinates on $Q$ such that the $x$ axis lies along $x_{1}$ while the $y$ axis along $x_{2}$. Then

$$
\begin{equation*}
x=\frac{x_{1}}{1-x_{0}}, \quad y=\frac{x_{2}}{1-x_{0}} . \tag{15}
\end{equation*}
$$

An inverse map $f^{-1}: Q \rightarrow S^{2}(1)$, where $x \neq \pm 1$, is:

$$
\begin{equation*}
x_{0}=-\frac{1+x^{2}}{1-x^{2}}, \quad x_{1}=\frac{2 x}{1-x^{2}}, \quad x_{2}=\frac{2 y}{1-x^{2}} . \tag{16}
\end{equation*}
$$

If we substitute formulas (15) into (13) then we obtain the relations between coordinates $x, y$ and adapted coordinates $u, \varphi$ which are on semi-Euclidean sphere:

$$
f: \quad x=\frac{\sinh \varphi}{\varepsilon-\cosh \varphi}, \quad y=\frac{u \exp \varphi}{\varepsilon-\cosh \varphi} .
$$

Then the inverse map is:

$$
\begin{equation*}
\varphi=\ln \left(\varepsilon \frac{x-1}{x+1}\right), \quad u=-\frac{2 y}{(1-x)^{2}} . \tag{17}
\end{equation*}
$$

Note that the lines $x= \pm 1$ are not included in the mapping and $Q$ consists of three disconnected components. Also, the line $x_{0}=1, x_{1}=0$ has no image in this mapping. We add it by hand, identifying the image of this line with the points $\{(x, y) \mid x=$ $= \pm \infty, y \in \mathbb{R}\}$ on $Q$. Then two disconnected parts $\{(x, y) \mid x \in(-\infty,-1), y \in \mathbb{R}\}$ and $\{(x, y) \mid x \in(1, \infty), y \in \mathbb{R}\}$ are connected and we call this plane $C^{2}$.

In particular, after enlarging $Q$ into $C^{2}$ by the infinitely distant point and ideal line crossing this point, then the stereographic map $f$ becomes diffeomorphism $S$. Note that the infinitely distant point is the image of point $N$. The ideal line is the image of the straight line belonging to $S^{2}(1)$ and crossing the pole: $x_{0}=1, x_{1}=0$.

Let us now consider the commutative diagram:


The map $p=\pi \circ S^{-1}: C^{2} \rightarrow R$ is defined by this diagram. We find the coordinate form of this map:

$$
u=-\frac{2 y}{(1-x)^{2}}
$$

The map $p: C^{2} \rightarrow \mathbb{R}$ defines the trivial bundle with the base $\mathbb{R}$ and the structure group $S_{1}$.

Theorem. Let $S$ is the map : $S^{2}(1) \rightarrow C^{2}$ as described before. Then $S$ is $a$ conformal map.

Proof. The metric on $G$ induces the metric on $C^{2}$. In the coordinates $x, y$ it has the form:

$$
\begin{equation*}
d \vec{s}^{2}=-d x^{2} \tag{18}
\end{equation*}
$$

Let us find the metric of semi-Euclidean sphere from the metric on $C^{2}$. From (17) we get $d \varphi=\frac{2}{x^{2}-1} d x$ and using (14) and (18) we find:

$$
d s_{1}^{2}=\frac{4}{\left(x^{2}-1\right)^{2}} d \vec{s}^{2} .
$$

Hence, the linear element of semi-Euclidean sphere differs from the linear element of $C^{2}$ by a conformal factor and therefore, the map $S$ is conformal.

We find the equation of fibers on $C^{2}$. The 1-parametric fibers family of the bundle $\left(S^{2}(1), \pi, \mathbb{R}\right)$ in the adaptive coordinates (13) is: $u=c, c \in \mathbb{R}$. From (17) we get the image of this family under the map $S$ :

$$
\begin{equation*}
y=-c / 2 \cdot(x-1)^{2} . \tag{19}
\end{equation*}
$$

The $C^{2}$ plane is also fibred by this 1-parametric family of parabolas.

## 4. Theprojective conformal model

Now we construct the projective semi-conformal model of the sphere $S^{2}(1)$ and the principal bundle on it. We use a normalization method of A.P.Norden [7], [8]. A. P. Shirokov in his work [14] constructed conformal models of Non-Euclidean spaces with this method.

Definition. A hypersurface $X_{n-1}$ as an absolute in a projective space $P_{n}$ is called normalized if with every point $Q \in X_{n-1}$ there is associated:

1) a line $P_{I}$ which has the point $Q$ as the only intersection with the tangent space $T_{n-1}$, and
2) a linear space $P_{n-2}$ that belongs to $T_{n-1}$, but it does not contain the point $Q$.

We call them normals of the first and second types, $P_{I}$ and $P_{I I}$.
In order to have a polar normalization, $P_{I}$ and $P_{I I}$ must be polar with respect to the absolute $X_{n-1}$.

We enlarge the semi-Euclidean space ${ }_{2} E_{1}^{3}$ to a projective space $P^{3}$. Here ${ }_{k} E_{l}^{n}$ denotes a $n$-dimensional semi-Euclidean space with the metric tensor of rank $k$, and $l$ is the number of negative inertia index in a quadric form. We consider homogeneous coordinates $\left(y_{0}: y_{1}: y_{2}: y_{3}\right)$ in $P^{3}$, where $x_{i}=\frac{y_{i}}{y_{3}}, i=0,1,2$. Thus $S^{2}(1): x_{0}^{2}-x_{1}^{2}=1$ describes the hyperquadric in $P^{3}$ :

$$
\begin{equation*}
y_{0}^{2}-y_{1}^{2}-y_{3}^{2}=0 . \tag{20}
\end{equation*}
$$

Here the projective basis $\left(E_{0}, E_{1}, E_{2}, E_{3}\right)$ is chosen in the following way. The vertex $E_{0}$ of basis is inside the hyperquadric. The other vertices $E_{1}, E_{2}, E_{3}$ are on its polar plane, $y_{0}=0$. The line $E_{0} E_{3}$ crosses the hyperquadric at poles $N(1: 0: 0: 1)$, $N^{\prime}(1: 0: 0:-1)$. Vertices $E_{1}, E_{2}$ lie on the polar of the line $E_{0} E_{3}$. The vertex of the hyperquadric coincides with the vertex $E_{2}$.

The stereographic map of the projective plane $P^{2}: y_{0}=0$ to the hyperquadric (20) from the pole $N(1: 0: 0: 1)$ is shown on the picture. Let $U\left(0: y_{1}: y_{2}: y_{3}\right) \in P^{2}$. If $y_{3}=0$, then the line $U N$ belongs to the tangent plane $T_{N}: y_{0}-y_{3}=0$ of the hyperquadric (20) at the point $N$ and in this case the intersection point of the line $U N$ with the hyperquadric is not uniquely determined. If $y_{3} \neq 0$, then the intersection point of the line $U N$ with the hyperquadric is unique. So, we choose the line $E_{1} E_{2}$ :

$: y_{3}=0$ as the line at infinity. In the area $y_{3} \neq 0$ we consider the Cartesian coordinates $x_{1}=\frac{y_{1}}{y_{3}}, x_{2}=\frac{y_{2}}{y_{3}}$. Then the plane $\alpha: y_{0}=0, y_{3} \neq 0$ becomes a plane with an affine structure $A^{2}$. It is possible to introduce the structure of semi-Euclidean plane ${ }_{1} E^{2}$ with the linear element

$$
\begin{equation*}
d s_{0}^{2}=d x_{1}^{2} \tag{21}
\end{equation*}
$$

The hyperquadric and the plane $\alpha$ do not intersect each other or intersect in two imaginary parallel lines

$$
\begin{equation*}
x_{1}^{2}=-1 . \tag{22}
\end{equation*}
$$

The restriction of the stereographic projection to the plane $\alpha$ maps the point $U\left(0: x_{1}: x_{2}: 1\right)$ into the point $X_{1}$

$$
\begin{equation*}
X_{1}\left(-1-x_{1}^{2}: 2 x_{1}: 2 x_{2}: 1-x_{1}^{2}\right) \tag{23}
\end{equation*}
$$

So, the Cartesian coordinates $x_{i}$ can be used as the local coordinates at the hyperquadric except the point of its intersection with the tangent plane $T_{N}$.

We construct an autopolar normalization of the hyperquadric. As a normal of the first type we take lines with the fixed center $E_{0}$ and as a normal of the second type we take their polar lines which belong to the plane $\alpha$ and cross the vertex $E_{2}$ of the hyperquadric. The line $E_{0} X_{1}$ intersects the plane $\alpha$ at the point

$$
X\left(0: 2 x_{1}: 2 x_{2}: 1-x_{1}^{2}\right)
$$

In this normalization the polar of the point $X$ intersects the plane $\alpha$ on the normal $P_{I I}$. Thus for any point $X$ in the plane $\alpha$ there corresponds a line which does not cross this point. It means that the plane $\alpha$ is also normalized. The normalization of $\alpha$ is defined by an absolute quadric (22).

We consider the derivative equations of this normalization. If we take normals of the first type with fixed center $E_{0}$, then the derivative equations ([7], p. 204) have the form:

$$
\begin{gather*}
\partial_{i} X=Y_{i}+l_{i} X \\
\nabla_{j} Y_{i}=l_{j} Y_{i}+p_{j i} X \tag{24}
\end{gather*}
$$

The points $X, Y_{i}, E_{0}$ define a family of projective frames. Here $Y_{i}$ are generating points of the normal $P_{I I}$.

We can calculate the values $(X, X),\left(X, Y_{i}\right)$ on the plane $\alpha$ using the quadric form, which is in the left part of equation (20). So, $(X, X)=-\left(1+x_{1}^{2}\right)^{2}$.

Let us find coordinates of the metric tensor on the plane $\alpha$. Hence, we take the Weierstrass standardization

$$
(\widetilde{X}, \widetilde{X})=-1, \quad \widetilde{X}=\frac{X}{1+x_{1}^{2}}
$$

Then the coordinates of the metric tensor are the scalar products of partial derivatives $g_{i j}=-\left(\partial_{i} \widetilde{X}, \partial_{j} \widetilde{X}\right)$ :

$$
\left(g_{i j}\right)=\left(\begin{array}{cc}
\frac{4}{\left(1+x_{1}^{2}\right)^{2}} & 0 \\
0 & 0
\end{array}\right)
$$

We got the conformal model of the polar normalized plane $\alpha: y_{0}=0, y_{3} \neq 0$ with a linear element

$$
\begin{equation*}
d s^{2}=\frac{d x_{1}^{2}}{\left(1+x_{1}^{2}\right)^{2}} \tag{25}
\end{equation*}
$$

It means that the following theorem is true:
Theorem. The non-Euclidean plane $\alpha$ is conformally equivalent to semi-Euclidean plane ${ }_{1} E^{2}$.

The points $X$ and $Y_{i}$ are conjugated with respect to the polar (20) and $\left(X, Y_{i}\right)=0$. From this equation and the derivative equations (24) we can get the non-zero connection coefficients:

$$
\Gamma_{11}^{1}=\Gamma_{12}^{2}=\Gamma_{21}^{2}=-\frac{2 x_{1}}{1+x_{1}^{2}}, \quad \Gamma_{11}^{2}=\frac{2 x_{2}}{1+x_{1}^{2}}
$$

The sums $\Gamma_{k s}^{s}=\partial_{k} \ln \frac{c}{\left(1+x_{1}^{2}\right)^{2}}(c=$ const $)$ are gradients, so the connection is equiaffine. Curvature tensor has the following non-zero elements:

$$
R_{121 \cdot}^{2}=-R_{211 \cdot}^{2}=-\frac{4}{\left(1+x_{1}^{2}\right)^{2}}
$$

Ricci curvature tensor $R_{s k}=R_{i s k}{ }^{i}$. is symmetric: $R_{11}=\frac{4}{\left(1+x_{1}^{2}\right)^{2}}$. Metric $g_{i j}$ and curvature $R_{r s k}{ }^{i}$ tensors are covariantly constant in this connection: $\nabla_{k} g_{i j}=0, \nabla_{l} R_{r s k}{ }^{i}=0$.

This can be summarized into a proposition:
Proposition. The autopolar normalization of the hyperquadric (20) constructed above defines an equiaffine connection on it with symmetric Ricci curvature tensor and covariantly constant metric and curvature tensors.

The infinitesimal linear operators for the quadric are

$$
\left\{\begin{array}{l}
L_{1}=y_{0} \frac{\partial}{\partial y_{1}}+y_{1} \frac{\partial}{\partial y_{0}},  \tag{26}\\
L_{2}=y_{0} \frac{\partial}{\partial y_{3}}+y_{3} \frac{\partial}{\partial y_{0}}, \\
L_{3}=y_{1} \frac{\partial}{\partial y_{3}}-y_{3} \frac{\partial}{\partial y_{1}} .
\end{array}\right.
$$

Solving geodesic equations we find parametric solutions

$$
\left\{\begin{array}{l}
x_{1}=\tan (\omega t+\phi)  \tag{27}\\
x_{2}=\left(c_{1} e^{2 i \omega t}+c_{2} e^{-2 i \omega t}\right) \sec ^{2}(\omega t+\phi)
\end{array}\right.
$$

where $c_{1}, c_{2}, \omega, \phi$ are integration constants. Eliminating the parameter $t$ we can rewrite these equations in a simple form

$$
x_{2}=A\left(x_{1}^{2}-1\right)+B x_{1},
$$

where $A$ and $B$ are arbitrary constants. We see that the solution represents parabolas and lines in $x_{1} x_{2}$-plane.

Let us consider the bundle of this plane by the double numbers subalgebra. We write the equations of fibers of semi-Euclidean sphere $S^{2}(1)$ in homogeneous coordinates:

$$
\left\{\begin{array}{l}
\left(y_{0}-y_{1}\right) v-y_{2}=0  \tag{28}\\
y_{0}^{2}-y_{1}^{2}-y_{3}^{2}=0
\end{array}\right.
$$

This 1-parametric family of curves fibers the hyperquadric and it defines a bundle on it. The image of these fibers under the stereographic projection from the pole $N$ to the plane $\alpha$ is:

$$
x_{2}=-v / 2 \cdot\left(x_{1}+1\right)^{2} .
$$

It is 1-parametric family of parabolas (compare with (19)).

## Conclusion

General program is to study non-Euclidean spaces generated by unital associative algebras. In this paper we give an example of pseudo-Euclidean space and then we present two ways how to construct models of its fibration, which are conformal and projective models of pseudo-Euclidean sphere. For both models we get metric and images of a fibration. For the second we use a normalization method to construct an equiaffine connection, then we find infinitesimal linear operators and geodesics.

We would obtain the similar results for the space of right cosets by the Lie subgroup $H_{2}$ (it is the subgroup of invertible dual numbers) and the bundle of the group $G$ by $H_{2}$. However, $H_{2}$ is a normal divisor of the group $G$. Therefore, the spaces of right and left cosets coincide.

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## References

[1] Bourbaki, N.: Elements of Mathematics, Algebra I, Paris, Hermann, Reading, Mass., AddisonWesley, 1989.
[2] Encyclopaedia of Mathematical Sciences, vol. 18, Algebra II. Noncommutative Rings, identities/ A. I. Kostrikin, I. R. Shafarevich, eds. Berlin, New York: Springer Verlag, 1991.
[3] Husemoller, N.: Fibre Bundles, New York: McGraw-Hill, 1966.
[4] Kirichenko, V. F., Arseneva, O. E.: Differential Geometry of generalized almost quaternionic structures, eprint arXiv: dg-ga/9702013.
[5] Malakhal'tsev, M. A.: A class of manifolds over the algebra of dual numbers, (Russian) Trudy Geom. Sem. Kazan. Univ., No. 21 (1991), 70-79; MR1195520 (94f:53053).
[6] Morimoto, A.: Prolongation of connections to bundles of infinitely near points, J. Differ. Geom., No. 11(4) (1976), 479-498.
[7] Norden, A. P.: Prostranstva affinnoi sviaznosti, (Russian) [Spaces with Affine Connection], Moscow: Nauka, 1976; MR0467565 (57 \#7421).
[8] Norden, A. P.: A generalization of the fundamental theorem of the theory of normalization, (Russian) Izv. VUZ Mat., No. 2(51) (1966), 78-82; MR0196663 (33 \#4850).
[9] Pavlov, E. V., Hopteriev, H. T.: Manifolds with an algebraic structure and CH-mapping, Bulg. Acad. of Science press, T. 35, No. 2 (1982), 141-144; MR0666267 (83i:53041).
[10] Redei, L.: Foundations of Euclidean and Non-Euclidean Geometries According to F. Klein, Oxford, New York, Pergamon Press, 1968.
[11] Rosenfeld, B. A.: A History of Non-Euclidean Geometry: Evolution of the Concept of a Geometric Space; translated by Abe Shenitzer with the editorial assistance of Hardy Grant, New York: Springer Verlag, 1988.
[12] Rosenfeld, B. A.: Geometry of Lie Groups, Boston: Kluwer, 1991.
[13] Shiroкоv, A. P.: Spaces over algebras and their applications, J. Math. Sci., Vol. 108, No. 2 (2002), 232-248.
[14] Shirokov, A. P.: Neevklidovy prostranstva, (Russian) [Non-Euclidean Spaces], Izdat. Kazan. univ., 1997; Zbl 0933.51011.
[15] Study, E., Cartan, E.: Nombres complexes, Encyclopedie des sciences mathematiques pures et appliquees, t. 1. Vol. 1. (1908) 329-468.
[16] Shurygin, V. V.: Manifolds over local algebras that are equivalent to jet bundles, (Russian) Izv. VUZ Mat., No. 10 (1992), 68-79; translation in Russian Math. (Izv. VUZ Mat.), 36 (1993) No. 10, 66-77.
[17] Vishnevskii, V. V.: Integrable affinor structures and their plural interpretations, J. Math. Sci., 108 (2002), No. 2, 151-187.
[18] Vishnevski, V. V., Shirokov, A. P., Shurygin, V. V.: Prostranstva nad algebrami, (Russian) [Spaces over Algebras] Kazanskii Gosudarstvennyi Universitet, Kazan, 1985; MR0928390 (89m:53002).
[19] Yaglom, I. M.: A Simple Non-Euclidean Geometry and its Physical Basis: an Elementary Account of Galilean Geometry and the Galilean Principle of Relativity; translated by Abe Shenitzer with the editorial assistance of Basil Gordon, New York: Springer, 1979.
[20] Yano, K.: Differential Geometry on Complex and Almost Complex Spaces, Pergamon. Press, New York, 1965.


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[^1]:    ${ }^{1}$ Irreducible means indecomposable into a direct sum of algebras.

