Ferenc Fazekas Deterministic and stochastic vector differential equations applied in technical systems theory

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DETERMINISTIC AND STOCHASTIC VECTOR DIFFERENTIAL EQUATIONS APPLIED IN TECHNICAL SYSTEMS THEORY

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1. This paper will give a choice from different systems of technics, physics, fight, astronautic etc. treated mainly by vector differential equations (vDE) in our papers, books, bulletins [1-13], having various results. These can illustrate the mathematical and technical variety and complexity of such problems, of course, without the claim to totality. A firm use of matrix analysis & algebra will accompany the following research.

1.1. As well known [4,14], the models of (deterministic) danymic systems are often described in the state (S) space by its <u>SvDE</u> and by its output algebraic one OvAE (which remains here in background):

 $\frac{\dot{z}}{t} = \frac{f[z, x(t), t]}{t_0}, \quad y = g[z, x(t), t], \quad (1, 11a, b)$ $\langle t_0 \leq t, \ z_0(t_0) = z_0; \quad z_0(t) = ? \quad y_0(t) = ? \rangle$

where SvDE is supposed as satisfying the existency & unicity conditions. This generally <u>non-linear</u> (nl.) SvDE can have <u>special forms</u> occasionally, namely [4-14]

 $\underline{\dot{z}} = \underline{f}[\underline{z}, \underline{x}(t)] \text{ time-invariant (t.inv.)}, \underline{\dot{z}} = \underline{f}(z) \text{ autonomous, } \underline{\dot{z}} = \underline{\dot{z}} = \underline{A}(t)\underline{z} + \underline{B}x(t) \text{ linear (l.)}, \underline{\dot{z}} = \underline{A}\underline{z} + \underline{B}\underline{x}(t) \text{ l.t.inv.}, \underline{\dot{z}} = [-\underline{e}_{n}\underline{p}^{*}(t) + \underline{K}] z + \underline{e}_{n}x(t) \langle z_{1} = z^{(1-1)}, \underline{K} = [\delta_{1,j-1}], z^{(n)} + \frac{n-1}{2} \underline{p}_{1}(t)z^{(1)} = x(t) \rangle$

phase-vDE of a 1.DE. etc. and similarly the OvAE.

1.2. A system can have better and worse <u>models</u> [4a] so such SvDEs too, according to more, or less abstractions from the reality; but the truth of model to reality and the mathematical handling of SvDE is often compatible by compromise only.

1.3. As a hepl for the further treatment, let be mentioned our \underline{dy} <u>namic transform algorithm</u> (DTA) for a matrix $\underline{A} \cong \underline{A}_0$ in p(< r) steps
(in a spring: $\underline{A}_p = \underline{A}_0 - \frac{p^{-1}}{\substack{L\\q=0}} \gamma_q (\underline{a}_{1q}^{(q)} - \underline{e}_{kq}) (\underline{a}_{q}^{(q)} + \underline{e}^{(q)}) =$ [2,5] (D) $= \underline{A}_0 - (\underline{A}_L - \underline{E}_K) \underline{\Gamma}_{KL} (\underline{A}^K + \underline{E}^L) = \begin{bmatrix} \underline{\Gamma}_{KL} & \underline{\Gamma}_{KL} \underline{A}_{KJ} \\ -\underline{A}_{1L} \underline{\Gamma}_{KL} & \underline{\Omega}_{1J} \end{bmatrix}.$

- The rank $\rho(\underline{A})$ is r, if the p = rth step let vanish the free block: $\underline{\Omega}_{IJ} = \underline{O}_{1j} \dots (1.32)$ - For a regular \underline{A} having n = m = r, the p = nth step furnishes (with $\forall a_{kqkq} \neq 0$ at different $k_q \in K$) its inverse matrix: $\underline{A}_n = \underline{A}^{-1} = \underline{\Gamma} \dots (1.33)$ - This DTA is suitable in algebras (A) (D) to solve arbitrary 1.vAE [5], 1.programming [2] and is generalized

(gDTA) to solve nl.vAE [5] too. ...(1,34a-c) - If $\underline{A}' = \underline{A}_0 + \underline{bd}^* = \underline{A} + \underline{A}\beta \cdot \underline{\delta}^* \underline{A} = \underline{A}(\underline{E} + \underline{\beta}\underline{d}^*)$, then $\underline{A'}^{-1} \triangleq \underline{\Gamma}'$ can be found [15] in the form and with the scalar factor

$$\underline{\Gamma}' = \underline{\Gamma} - \underline{x}\underline{\beta}\underline{\delta}^{*}\underline{(\underline{E} - \underline{x}\underline{\beta}\underline{d}^{*})}\underline{\Gamma}, \quad x = (1 + \underline{d}^{*}\underline{\beta})^{-1} \langle \underline{d}^{*}\underline{\beta} + -1 \rangle, \quad (1.35)$$

as easy to control $(\underline{A}'\underline{\Gamma}' = \dots = \underline{E})$. - We have created a set of matrix algorithmic methods (MAM) [5,6,8,13] for various purpose; e.g. STA, SMA, OMA, TAD, IAA, ITA, OTA, FA, SOTA etc.

2. Let us make some remarks on the non-linear SDEs and their solutions!

2.1. Such ones can be solved *exactly* in exceptional cases only. -A such problem in [4e,14] is <u>the pursuiting motion of an averting roc-</u> <u>ket</u> R_1 in trace of an attacking on R_2 (in the vertical plane); namely at radial velocity v, distance $r = \overline{R_1 R_2}$ and incl. angle φ of R_1 , at horizontal velocity c of R_2 and at ratio $m \stackrel{c}{=} v/c > 1$ - the SDEs are as follow:

$$\underbrace{(\underline{0} \geq)}_{\underline{z}} \triangleq \begin{bmatrix} \dot{\mathbf{r}} \\ \dot{\mathbf{r}} \\ \dot{\mathbf{r}} \end{bmatrix} = \begin{bmatrix} c \frac{\cos\varphi}{r} & -\frac{\mathbf{v}}{\varphi} \\ -c \frac{\sin\varphi}{r} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{r} \\ \varphi \end{bmatrix} \triangleq \underline{A}(\underline{z})\underline{z} = \begin{bmatrix} c\cos\varphi & -\mathbf{v} \\ -c\sin\varphi + 0 \end{bmatrix} \triangleq \mathbf{f}(\underline{z}), \ \underline{z}(\mathbf{t}_0) = \begin{bmatrix} \mathbf{r}_0 \\ \pi/2 \end{bmatrix} (2.11)$$
 and the exact solution:

$$\frac{\mathbf{r}}{\mathbf{r}_{0}} = e^{\pi/2} \frac{\int_{0}^{\sqrt{2}} \frac{\cos(\psi' - \mathbf{v})}{\cos(\psi')} d\psi'}{\int_{0}^{\sqrt{2}} \frac{\operatorname{tg}^{\mathrm{m}\psi}}{\sin\psi}} \quad (m > 1).$$
(2.12)

2.2. A nl.SvDE is often solved approximately by (local) *lineari*zation around its equilibry (EL) points. - It is proposed in our [4e] for the growth vDE of two rival rasses with z_i populations

$$\underline{\dot{z}} \triangleq \begin{bmatrix} \dot{z}_1 \\ \dot{z}_2 \end{bmatrix} = \begin{bmatrix} K_1 z_1 - M_1 z_1 z_2 \\ -M_2 z_1 z_2 + K_2 z_2 \end{bmatrix} = \begin{bmatrix} K_1 z_1 (1 - z_2 / n_1) \\ K_2 z_2 (1 - z_1 / n_2) \end{bmatrix} \triangleq \underline{f(\underline{z})} \quad (\underline{M}_i > 0, K_i > 0; \\ (1 <) K_2 / K_1 = k, n_1 = K_1 / M_1)$$
(2.21)

with EL-situations (K = I,II)

$$\underline{\dot{z}}_{K} = \underline{f}(\underline{z}_{K}) = \underline{0}$$
 at $\underline{z}_{I} = \underline{0}$ and $\underline{z}_{II} = \begin{bmatrix} n_{2} \\ n_{1} \end{bmatrix}$. (2.22)

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The T_1 Taylor-polynom gives the approximate SvDE at z_K

$$\frac{\dot{z}_{K}}{\underline{z}_{K}} \approx \underbrace{\dot{0}}_{I} + \underbrace{F}(\underline{z}_{K}) d\underline{z}_{K} \stackrel{\neq}{=} \begin{bmatrix} (n_{1}-z_{2})M_{1}-M_{1}z_{1}\\ -M_{2}z_{2}(n_{2}-z_{1})M_{2} \end{bmatrix} \begin{bmatrix} dz_{1}\\ dz_{2} \end{bmatrix}_{K} ; \underbrace{F}_{I} \stackrel{=}{=} \begin{bmatrix} \kappa & 0\\ 0 & \kappa_{2} \end{bmatrix},$$

$$\underbrace{F}_{II} \stackrel{\hat{=}}{=} \begin{bmatrix} 0 & -M_{1}n_{2}\\ -M_{2}n_{1} & 0 \end{bmatrix}; \quad (\kappa = I, II; \underbrace{F}(\underline{z}_{K}) = \underbrace{F}_{K}),$$
then the eigen-values (1.at 3,2) and approximate solutions around \underline{z}_{K}

$$|\lambda\underline{E} - \underbrace{F}_{I}| = (\lambda - \kappa_{1})(\lambda - \kappa_{2}) = 0, \quad \lambda_{1} = \kappa_{1} > 0 \end{bmatrix} |\lambda\underline{E} - \underbrace{F}_{II}| \stackrel{\hat{=}}{=} \lambda^{2} - \kappa_{1}\kappa_{2} = 0,$$

2.3. The analytical difficulty of a nl. SvDE (1,11a) can sometimes invert to an algebraical facility by the *difference method*, as approximation, or if the problem itself has a structure of difference. -This is the case at <u>the bending and moving equations of a chain bridge</u> treated in our [4e,1]. The chain connected with links from rigid bars let be characterized at the end links by horizontal strains h/H = = h + Δ h, at link $\mathbf{x}_i = \sum_{j=1}^{i} \Delta \mathbf{x}_i$ = ia by hangs down $\mathbf{y}_i \in \mathbf{y}/\mathbf{y}_i + \mathbf{v}_i \in \mathbf{y} + \mathbf{y}$, loads p = P/a $\in \mathbf{p}$ (own)/ $\mathbf{q}_i = \mathbf{Q}_i/a \in \mathbf{q}$ (useful); its balance vEs (using the continuant matrix $\mathbf{c} = [\mathbf{c}^i]_{n-1}^1 = [0, \dots, 0, 1, 2, -1, \dots, 0]$, the vector $\mathbf{e} = [1]^1$ and the fact $-\frac{1}{2}\mathbf{c}\mathbf{y} = \begin{bmatrix} \Delta^2 \mathbf{y}_i \\ \Delta \mathbf{x}_i^2 \end{bmatrix}$ and remarking in (...) the (at $1/n = \mathbf{a} + \mathbf{0}$) correspondant DEs appear as follow [1]:

$$\frac{1}{a^2} \underbrace{\underline{C}}_{\underline{Y}} = \frac{1}{h} \underbrace{p} \left\langle -\underline{y'} + \underline{p} \right\rangle, \quad \frac{1}{a^2} \underbrace{\underline{C}}_{\underline{Y}} + \underline{v} = \frac{1}{H} \underbrace{p}_{\underline{Y}} \left\{ \underline{p} + \underline{q} \right\rangle \left\langle -(\underline{y'}' + \underline{v'}') = \frac{1}{H} \begin{bmatrix} \underline{p} + \underline{q} \\ \underline{p} + \underline{q} \\ \underline{r} & \underline{r} \\ \underline{r} \\ \underline{r} & \underline{r} \\ \underline{r} \\$$

The rigid <u>pendant bars</u> at x_i transfer the emotions v_i of the loaded chain to the beam and establish with m, the *relation* [1]

$$\frac{1}{\mathrm{EI}} \underset{=}{\mathrm{Km}} \stackrel{\diamond}{=} \frac{1}{\mathrm{EI}} \left(\underset{=}{\mathrm{E}} - \frac{1}{6} \underset{=}{\mathrm{C}} \right) \underset{=}{\mathrm{m}} = \frac{1}{a^2} \underset{=}{\mathrm{Cv}} \left\langle \begin{array}{c} \frac{1}{a^2} \underset{=}{\mathrm{EI}} \left(\underset{=}{\mathrm{m}} + \frac{a^2}{6} \underset{=}{\mathrm{m''}} \right) = -v'' \right\rangle.$$
(2,33)

With union of (2,31a-2,33), the basic vE of bending for the chain will be formed as follows:

$$\frac{1}{a^2} \underset{=}{\overset{L}(\underline{c})}{\underline{m}} \stackrel{\triangleq}{=} \frac{1}{a^2} \left[\underbrace{\underline{c}}_{\underline{c}} + \frac{\underline{Ha}^2}{\underline{EI}} (\underline{\underline{e}}_{\underline{c}} - \frac{1}{\underline{b}\underline{c}}) \right] \underline{\underline{m}} = (\underline{q} - \underline{\underline{q}}) + \underbrace{(\underline{q} - \underline{\Delta h}\underline{p})}_{\underline{c}} = \underline{q} - \underline{\underline{\Delta h}\underline{p}} \stackrel{\triangleq}{=} \underline{\underline{r}}$$

$$(\underline{L}(-\underline{d}^2)\underline{m} \stackrel{\triangleq}{=} -\underline{\underline{m'}'} + \frac{\underline{H}}{\underline{EI}} (\underline{m} + \frac{\underline{a}^2}{\underline{b} - \underline{m'}'}) \equiv (\underline{\underline{Ha}^2}_{\underline{EI}} - 1)\underline{m''} + \frac{\underline{H}}{\underline{EI}} \underline{\underline{m}} = \underline{q} (\underline{x}) - \underline{\underline{\Delta h}\underline{p}} \stackrel{\triangleq}{=} \underline{r} (\underline{x})) .$$

Its formal solution $\underline{m} = a^{2}\underline{L}^{-1}(\underline{C})\underline{r}$ can be surely realized, because the eigen-values -vectors are well known [16]:

$$\lambda_{i} = 2\cos \frac{i\pi}{n}, \ \underline{u}_{i} = \sqrt{\frac{2}{n}} [\sin \frac{ik\pi}{n}]_{n-1}^{k=1} (\underline{\underline{c}}\underline{\underline{u}} = \lambda_{i}\underline{\underline{u}}, \underline{\underline{u}}_{i}\underline{\underline{u}}_{j} = \delta_{ij}),$$
$$(\forall i \in \{1, 2, \dots, n-1\})$$
(2,35a,b)

then $L^{-1}(\lambda) \stackrel{\circ}{=} \left[\lambda + \frac{Ha^2}{EI}(1 - \frac{\lambda}{6})\right]^{-1}$ is rational function, finally the form $\underline{r} = \sum_{i} \left(\underline{u}_{i} \stackrel{\circ}{\underline{r}}\right) \underline{u}_{i} \stackrel{\circ}{=} \sum_{i} \rho_{i} \underline{u}_{i}$ is ready, so the s.c. <u>canonical formed</u> bending vector <u>m</u> can be writte as follows [4e]:

$$\underline{\mathbf{m}} = \sum_{i=1}^{n-1} \mathbf{L}^{-1}(\lambda_i) \underline{\mathbf{u}}_i(\underline{\mathbf{u}}_i^* \underline{\mathbf{r}}) \equiv \sum_{i=1}^{n-1} \frac{\mathbf{a}^2 \rho_i \underline{\mathbf{u}}_i}{\lambda_i + \frac{\mathbf{Ha}^2}{\mathbf{EI}} (1 - \frac{\lambda_i}{\mathbf{6}})} \cdot (2, 36)$$

Having it, the solving <u>emotion vector</u> appears so: $\underline{\mathbf{v}} = \frac{\mathbf{a}^2}{\mathrm{Er}_{\pm}} \mathbf{c}^{-1} \underline{\mathbf{k}} \underline{\mathbf{m}} = \frac{\mathbf{a}^2}{\mathrm{Er}} \mathbf{c}^{-1} (\underline{\mathbf{k}} - \frac{1}{6} \underline{\mathbf{c}}) \underline{\mathbf{m}} . \qquad (2,37)$

Remarkable, that the our upper procedure (1. in \underline{m}) is more simple and suitable as other ones (nl. in \underline{v}) [15].

2.4. Omitting the various numerical methods of Euler, Euler-Adams, Milne, predictor-corrector etc. [17], let be mentioned only the Runge -Kutta method (RKM) to solve the nl. SvDE (1,11a). Namely - for an interval [t,t+ τ] of length $\tau \stackrel{\land}{=} dt \sim 0$ and with signs $\underline{\xi} \stackrel{\land}{=} d\underline{x} \sim 0$, $\underline{\xi}_1 \stackrel{\land}{=} \frac{1}{2} d\underline{z} \equiv \underline{\dot{z}} dt = \underline{f}(\underline{z},\underline{x},t).\tau - \underline{a}$ procedure step advances as an algorithm of 4 substeps e.g. by our recurrent formula [4e].

$$\underline{\zeta}_{S} \stackrel{\circ}{=} \underline{\underline{F}}_{S} \equiv \tau \cdot \sum_{i=1}^{4} \underline{f}(\hat{\underline{Y}}_{i-1}^{*}) \ \mathbf{s}_{i} \equiv \tau \cdot \sum_{i=1}^{4} \underline{f}(\hat{\underline{Y}}_{0}^{*} + \mathbf{1}_{i}\hat{\mathbf{n}}_{i-1}^{*}) \ \mathbf{s}_{i}$$

$$(2,41)$$

$$\langle \underline{s} \stackrel{\circ}{=} \frac{1}{6} [1,2,2,1]^{*}, \ \underline{1} \stackrel{\circ}{=} [0,1/2,1/2,1]^{*}, \ \hat{\underline{Y}}_{0}^{*} \stackrel{\circ}{=} [\underline{z},\underline{x},t], \ \hat{\underline{n}}_{i-1}^{*} \stackrel{\circ}{=} [\underline{\zeta}_{i-1},\underline{\xi},\tau] \rangle$$

$$having - verificably [17] - the excellent accuracy /(with. suppl. 0-s)$$

$$|\underline{\zeta}_{\tau} - \underline{\zeta}_{S}| \stackrel{\circ}{=} |\underline{f}(\hat{\underline{Y}}_{0}^{*} + \hat{\underline{n}}_{1}^{*}) - \underline{f}(\hat{\underline{Y}}_{0}^{*} + \hat{\underline{n}}_{S}^{*})| \cdot |\tau| = c \cdot |\tau|^{5}.$$

$$(2,42)$$

2,5. Look at a problematics, where the former RKM is often used. - This is the motor vehicle as complex vibrating system treated in our [9]. Its (deterministic) model can be characterized very generally by the Lagrangean vDE of motion (one of second kind)

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\mathbf{q}}} - \frac{\partial L}{\partial \mathbf{q}} + \frac{\partial R}{\partial \dot{\mathbf{q}}} = \underline{P}(t) \text{ at } \delta \int L(\mathbf{q}, \dot{\mathbf{q}}, t') dt' = 0 \text{ for } \int L(\mathbf{q}, \dot{\mathbf{q}}, t') dt' = t_0 \qquad (2,51)$$

where (2,51a) is just the Eulerian vDE (as necessary condition) of the variation and extremum problem (2,51b,c), namely with T/U kinetic/ potential energy, L = T - U L-function, R dissipation, $\underline{q} = \frac{1}{1} [q_1] \frac{1}{f}$ vector of generalized coordinates, $\underline{p}(t) = \frac{1}{1} [p_1(t)] \frac{1}{f}$ external forces, f degree of freedom (e.g. at a car model can be f = 7). - The detailed form of (2.51) will be - after total derivation - as follows [10]:

$$\frac{\partial^2 \mathbf{L}}{\partial \underline{\mathbf{d}}} \frac{\partial^2 \mathbf{L}}{\partial \underline{\mathbf{d}}} \mathbf{\dot{\mathbf{q}}} \mathbf{\dot{\mathbf{q}}} + \frac{\partial^2 \mathbf{L}}{\partial \underline{\mathbf{d}}} \mathbf{\dot{\mathbf{q}}} \mathbf{\dot{\mathbf{q}}} + \frac{1}{\mathbf{q}^* \mathbf{q}} \left(\frac{\partial^2 \mathbf{L}}{\partial \underline{\mathbf{q}} \partial \mathbf{t}} - \frac{\partial \mathbf{L}}{\partial \underline{\mathbf{q}}} + \frac{\partial \mathbf{R}}{\partial \underline{\mathbf{d}}} \right) \mathbf{q}^* \mathbf{q} = \mathbf{p}(\mathbf{t}), \qquad (2,52a)$$

This is solved every now as *nl*.SvDE e.g. by the upper RKM [17,4e], then as *ql*. one: firstly $\underline{\dot{z}} = \underline{l}[\underline{z},\underline{x}(t)]$ by 4,1 to $\underline{z}_0(t)$, secondly $\underline{\dot{z}} = \underline{n}[\underline{z}_0(t), \underline{\dot{z}}_0(t), t]$ by integration to $\Delta \underline{z}_0(t)$ etc.

2,6. The stability of a nl. system is often contolled by Ljupanov's direct method [18].- We use here it for an <u>astronave</u> (N) with linear help-rocket (R) treated in [4e]. N is considered as a spin, whose known autonomous nl.SvDE - at R's l.vAE (of C > 0) - follows here: $\underline{\dot{z}} = \underline{z} \times \underline{I}_{1}^{-1} \underline{z} + \underline{I}_{1}^{-1} \underline{x}(t) \equiv \underline{z} \times \underline{I}_{1}^{-1} \underline{z} - \underline{I}_{1}^{-1} \underline{C}_{1} \underline{z} \equiv (\underline{z} \times \underline{I}_{1}^{-1} - \underline{I}_{1}^{-1} \underline{C}_{1} \underline{z}) \underline{z} \cong \underline{F}(\underline{z}) \underline{z}$ with the angular velocity's vector $\underline{z} \equiv \underline{\omega} \in E_{3}$ (through the centre of mass) and with the main inertia moments $(I_{1}, I_{2}, I_{3}) = \underline{I}_{1}$. Choosing $\underline{P}_{1} = \underline{I}_{1}^{2}$, the L.-function $V(\underline{z})$ and its derivative $W(\underline{z}) \cong \dot{V}(\underline{z}) =$ $= \underline{grad} * V(\underline{z}) \cdot \underline{\dot{z}}$ will be: $V(\underline{z}) = \underline{z} * \underline{P}_{1} \underline{z} > 0$ (for $\forall \underline{z} \neq 0$, def. pos.), $W(\underline{z}) = \underline{\dot{z}}^{*} \underline{P}_{1} \underline{z} + \underline{z} * \underline{P}_{1} \underline{\dot{z}} = \underline{z} * [\underline{F} * (z) \underline{P}_{1} + \underline{P}_{1} \underline{F}(\underline{z})] \underline{z} \cong \underline{z} * \underline{N}_{1} (\underline{z}) \underline{z} = -2\underline{z} * \underline{Q}, \underline{I}_{1} \underline{z} < 0$ $(\forall \underline{z} \neq 0)$, which last quadratic form is def. nagative. Consequently, the equilibry point $\underline{z} = \underline{0}$ has globally asymptotic stability.

3. The point 2 had shown, there is a natural gravitation into the direction of linearity at the SvDE, for its relative simplicity (e.g. for the superponability ets.).

3,1. The homogeneous form of general 1. SVDE (1,12c) $\underline{\dot{z}} = \underline{A}(t) \underline{z}$ can be solved simply in possession of a basic matrix (bM) $\underline{Z}(t) =$ $= \frac{1}{2}[\underline{z}_j(t)]_n (|\underline{Z}(t)| \stackrel{\diamond}{=} Z(t) \neq 0 \text{ for } \forall t \in T = [t_0,\infty))$, when its general and a particular solution appears as $\underline{Z}(t) = \underline{Z}(t)\underline{C}_0 = \underline{Z}(t)\underline{Z}^{-1}(t_0)\underline{Z}_0 \stackrel{\diamond}{=} (3,11)$

$$\stackrel{\simeq}{=} \underline{\underline{z}}(t, t_0) \underline{\underline{z}}_0 \langle \underline{z}(t_0) = \underline{z}_0 \rangle$$

with the (by $\underline{\tilde{z}}(t_0, t_0) = \underline{E}$) normed (n.) bM [6].- Having a phase SVDE (1,12e) $\underline{\tilde{z}} + \underline{P}(t)\underline{z} = \underline{0}$, or $L_n[\underline{z}] = z^{(n)} + \underline{p}^*(t)\underline{z} = 0$ with $\underline{z} = [z_1]_n^1 = [z^{(i-1)}]_n^1$, the bM is $\underline{z}(t)\underline{r}_n(t_0)$ (Green vector) (at $\underline{r}_n(t_0) \subset \underline{R}(t_0) \triangleq \underline{z}^{-1}(t)$) [6].- Our algorithm SoTA [13] advances e.g. from a $L_4[z] = 0$ - by transforms $z = z_1 \int udt$, $u = u_2 \int vdt$, $v = v_3 \int wdt$ - into $L_1[w] = 0$, giving the factorization $z_4(w) = c$. $z_1^4(t)u_2^3(t)v_3^2(t)w_4(t)$ etc. (3.13)

3,2. At a time-invariant hom. form of l.vDE (1,12d) $\frac{\dot{z}}{2} - \underline{A}\underline{z} = \underline{0}$, so at the ql. motor vehicle problem of 2,5 (at $\underline{x} = \underline{0}$, $\Delta \underline{\dot{z}} = \underline{0}$) too [9], exponential solutions $\underline{z} = e^{\lambda t}\underline{u}$ are supposed, which guides to the eigen-value problem [4c]

$$\underline{z}^{(t)} = \underbrace{\bar{\mathbf{k}}}_{\sigma=1} \underbrace{\underline{\mathbf{u}}}_{\sigma} e^{\underline{A}} \underbrace{\mathbf{c}}_{\sigma} = \underbrace{\underline{\mathbf{u}}}_{\sigma} e^{\underline{A}} \underbrace{\mathbf{c}}_{\sigma} = e^{\underline{A}} \underbrace{\mathbf{c}}_{\underline{C}} \text{ and } \underline{z}_{0}^{(t)} = \underbrace{\underline{\mathbf{u}}}_{\underline{C}} e^{\underline{\mathbf{k}}} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} = \underbrace{\underline{\mathbf{u}}}_{\underline{C}} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} = \underbrace{\underline{\mathbf{u}}}_{\underline{C}} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} = \underbrace{\underline{\mathbf{u}}}_{\underline{C}} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} = \underbrace{\underline{\mathbf{u}}}_{\underline{C}} \underbrace{\mathbf{c}}}_{\underline{C}} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} = \underbrace{\underline{\mathbf{u}}}_{\underline{C}} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}^{(t)}} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)} \underbrace{\mathbf{c}}}_{\underline{C}}^{(t)$$

with exponential nbM $\underline{\tilde{2}}(t,t_0) = \underline{\tilde{2}}(t-t_0)$. (L. at [40] for $\exists \beta_{\sigma} > 1$.) For the stability, all $\lambda_{\sigma} = \mu_{\sigma} + i\nu_{\sigma}$ must have $\mu_{\sigma} \stackrel{\circ}{=} \operatorname{Re} \lambda_{\sigma} < 0$. (3,23)

3,3. A problem of type 3,2 can be also very complicated one. This is illustrated by the rotating system of a rotor (R) and $n \sim n$ axles $(\underline{A_i})$, as a *turbine's* model reached by matrix method in our [7]. It was our lecture's theme at the Equadiff-6; so let be enough here to refer it only!

3,4. In the general case of 3,1 $\frac{1}{2} = A(t) \frac{1}{2}$, there is'nt generally an exponential bM $g(t) = e^{\int \underline{A}(\tau) d\tau} = e^{\hat{\underline{A}}(t)}$ (because $\underline{\dot{z}} \stackrel{\circ}{=} e^{\hat{\underline{A}}}_{\underline{A}\underline{C}} + \underline{A}e^{\hat{\underline{A}}}_{\underline{C}} \stackrel{\circ}{=}$ Az, gen.) To find a bM for the SvDE or for its matrix variant $\underline{\mathbf{z}}_{\mathbf{j}}(\tau) = \underline{\underline{A}}(\tau)\underline{\mathbf{z}}_{\mathbf{j}}(\tau) \text{ (for } \forall \mathbf{j} \in \{1, 2, \dots, n\}, \text{ so } \underline{\underline{\mathbf{z}}}(\tau) = \underline{\underline{A}}(\tau)\underline{\underline{\mathbf{z}}}(\tau) \text{ (3,41)}$ and $\tilde{\tilde{z}}_{(\tau,t_0)} = A(\tau) \tilde{z}_{(\tau,t_0)}$, the integral equation of Volterra-type [4c] $\widetilde{\underline{\mathbf{Z}}}(\mathsf{t},\mathsf{t}_{0}) = \underline{\underline{\mathbf{E}}}_{t} + \int_{t_{0}}^{t} \underline{\underline{\mathbf{E}}}(\tau) \widetilde{\underline{\mathbf{Z}}}(\tau,\mathsf{t}_{0}) d\tau$ (3, 42 - 43)will be solved by the Picard-iteration $(\forall t, t_0 \in T)$ $\widetilde{\underline{z}}_{0}^{(t,t_{0})} = \underline{\underline{E}}, \quad \widetilde{\underline{\underline{z}}}_{k+1}^{(t,t_{0})} = \underline{\underline{E}} + \int_{t_{0}}^{t} \underline{\underline{\underline{A}}}_{(\tau)}^{(t)} \underline{\underline{z}}_{k}^{(\tau,t)} d\tau \quad (k=0,1,2,\ldots,n,\ldots) \quad (3A4)$ <u>obtaing</u> so the Neuman-series $\tilde{\underline{z}}(t,t_0) = \underline{\underline{E}} + \sum_{K=\tau}^{N} \tilde{\underline{A}}^{K}(t,t_0) \approx \tilde{\underline{z}}(t,t_0)$. - A regular transform $z = U(t)v \langle U(t) \neq 0$ for $\forall t \in T \rangle$ and a 1. MDE $\dot{U}(t) = Q(t)U(t)$ sometimes guide to a diagonalized form [13] \dot{v} = $= \underbrace{\underline{v}}_{-1} (\underbrace{\underline{A}}_{-1} - \underbrace{\underline{0}}_{-1}) \underbrace{\underline{v}}_{-1} \stackrel{c}{=} \underbrace{\underline{A}}_{+} (t), \text{ so to an exp. nbM } \underbrace{\underline{\tilde{v}}}_{+} (t, t_{0}) =$ $= e^{\int_{0}^{\tau} \frac{1}{2} \sqrt{\tau} d\tau}$,... (3,45) if one can solve the eigen-value problem $[\underline{A}(t) - \underline{Q}(t) - \lambda_{j}(t)\underline{E}]\underline{u}_{j}(t) = \underline{0}, \forall j \in \mathbb{N} .$ 4. Let pass over to linear homogeneous systems. 4,1. In the general case of (1,12c), the solution (3,11a) of hom. 1.SvDE $\underline{z}(t) = \underline{z}(t)\underline{c}$ will be applied - by variation of constant \underline{c} into

 $\underline{c}(t) = ? - \text{ to the inhom. one [4] (at \underline{z}(t) = {}_{1}[\underline{z}_{j}(t)]_{n}, z(t) \neq 0).$ $(\underline{z} - \underline{A}\underline{z})\underline{c} + \underline{\underline{z}}\underline{c} = \underline{B}\underline{x}, \underline{c} = \underline{z}^{-1}\underline{B}\underline{x}, \text{ so } \underline{z}_{n}(t) \stackrel{c}{=} \underline{z}(t)\underline{c}(t) = \int_{0}^{t} \underline{z}(t,\tau)\underline{B}(\tau)\underline{x}(\tau)d\tau$ appears as (at t₀ with <u>0</u> conditioned) particular solutions. - In the phase case (1,12e) and (3,12), the (4,11c) formula is simplified [4] to the form (at \underline{R}(\tau) \stackrel{c}{=} \underline{z}^{-1}(\tau) \text{ and } \underline{z}_{n}(\tau,\tau) = \underline{e})

$$\underline{z}_{n}(t) = \int \underbrace{z}_{0}(t) \underline{R}_{\lambda}(\tau) \underline{e}_{n} x(\tau) d\tau = \int \underbrace{z}_{0}(t) \underline{r}_{n}(\tau) x(\tau) d\tau \triangleq t_{0}$$

$$= \int \underbrace{z}_{0}(t, \tau) x(\tau) d\tau.$$

$$(4,12)$$

4,2. In the time-invariant case of (1,12d) and 3,2, so at the motor vehicle problem of 2,5 and 3,2, the nb.M $\underline{\tilde{Z}}(t,\tau) = e^{\underline{A}(t-\tau)}$ let write the ordinary and eigen forms (with $\underline{s}_n(t) = \underline{U}^{-1}\underline{z}_n(t)$, $\underline{B}_{\underline{s}} = \underline{U}^{-1}\underline{B}$):

4,3. Let us treat - following [4e,18] - the dynamical optimalization of a linear control system on the basis of quadhatical criterium (QC). - Here must minimalize a Ljapunov-function of QC V(\underline{z}^{0}) beside the 1.SvDE⁰ $\underline{\dot{z}}^{0} = \underline{A}\underline{z}^{0} + \underline{B}\underline{x}^{0}(t)$ at an optimal feed-back $\underline{x}^{0}(t) = -\underline{K}\underline{z}^{0}(t)$ (with $\underline{K} = ?$): $V(\underline{z}^{0}) \stackrel{e}{=} \int_{\underline{z}}^{\infty} (\underline{z}^{0} \stackrel{*}{\underline{P}}\underline{z}^{0} + \underline{x}^{0} \underline{Q}\underline{x}^{0}) d\tau = \int_{\underline{z}}^{\infty} \underline{z}^{\hat{v}0} (\underline{P} + \underline{K} \stackrel{*}{\underline{Q}}\underline{K}) \underline{z}^{0} d\tau = Min!$ (4,31) ($\underline{P} = \underline{P} \stackrel{*}{,} \underline{z} \stackrel{*}{\underline{P}}\underline{z}^{0} \ge 0$ (+s.def.); $\underline{Q} = \underline{Q} \stackrel{*}{,} \underline{x} \stackrel{*}{\underline{v}^{0}} \underline{Q} \stackrel{*}{\underline{v}^{0}} > 0$ at $\underline{x}^{0} \pm \underline{0}$ (+def.), so \underline{Q}^{-1}). Supposing $V(\underline{z}^{0}) = \underline{z} \stackrel{*}{\underline{v}^{0}} \underline{R} \stackrel{*}{\underline{v}^{0}} > 0$ at $\underline{R} = \underline{R} \stackrel{*}{\underline{R}} = ?$ and $\underline{z}^{0} \pm \underline{0}$ (+def.), its derivative has double form: $W(\underline{z}^{0}) = -\underline{z} \stackrel{*}{\underline{v}^{0}} (\underline{P} + \underline{K} \stackrel{*}{\underline{Q}} \underline{K}) \underline{z}^{0} = \underline{z} \stackrel{*}{\underline{v}^{0}} \underline{R} \stackrel{*}{\underline{z}^{0}} = (4,32)$ $= \underline{z} \stackrel{*}{\underline{v}^{0}} (\underline{A} - \underline{B}\underline{K}) \stackrel{*}{\underline{x}} \frac{\underline{z}^{0}}{\underline{z}} (< 0 \text{ for } \forall \underline{z}^{0}]$: asympt. stab. supposed), similarly the coefficient matrix too: $\underline{W}(\underline{K}) \stackrel{*}{\underline{C}} (\underline{A} - \underline{B}\underline{K}) \stackrel{*}{\underline{R}} + \underline{R} (\underline{A} - \underline{B}\underline{K}) = -(\underline{P} + \underline{K} \stackrel{*}{\underline{Q}} \underline{K}),$ (4,33) where from $\partial \underline{W} / \partial \underline{K} \stackrel{*}{\underline{K}} = -\underline{B} \stackrel{*}{\underline{R}} \underline{R} = -\underline{Q} \stackrel{K}{\underline{K}}$ follows $\underline{K} = \underline{Q}^{-1} \underline{B} \stackrel{K}{\underline{B}} (\underline{B} = ?),$ as the optimal feed-back matrix. With this $\underline{K}(\underline{R})$, one obtains

 $\langle \underline{\mathbf{W}}[\underline{\mathbf{K}}(\underline{\mathbf{R}})] - \underline{\mathbf{W}}[\underline{\mathbf{K}}(\underline{\mathbf{R}})] \equiv \rangle \underline{\mathbf{Q}}(\underline{\mathbf{R}}) \stackrel{2}{=} (\underline{\mathbf{R}}\underline{\mathbf{A}} + \underline{\mathbf{A}}^*\underline{\mathbf{R}}) - \underline{\mathbf{R}}\underline{\mathbf{B}}\underline{\mathbf{Q}}^{-1}\underline{\mathbf{B}}^*\underline{\mathbf{R}} + \underline{\mathbf{P}} = \underline{\mathbf{Q}} \quad (=-\underline{\mathbf{R}}),$ as a degenerated Riccatian MDE ($\underline{\mathbf{R}}$ const.) being nl. (quadr.) MAE and with its solution $\underline{\mathbf{R}}$ (e.g. by our gDTA of (3,14c) [5]) the optimal control in final form: $\underline{\mathbf{x}}^0(\mathbf{t}) = -\underline{\mathbf{Q}}^{-1}\underline{\mathbf{B}}^*\underline{\mathbf{R}} \underline{\mathbf{z}}(\mathbf{t}) = -\underline{\mathbf{K}}\underline{\mathbf{z}}^0(\mathbf{t}).$

5. Finally, let us turn shortly to the stochastic systems!

5,1. To avoide the complications of stochastic analysis, there is advantageous to transform linearly an arbitrary $\xi(t)$ into its random basic product (Rbp) [8](with ordinary coordinate factor $\underline{X}(t)$)

$$\xi(t) = \underline{m}_{\xi}(t) + \sum_{l=1}^{\infty} x_{l}(t) \xi_{l} = \underline{m}_{\xi}(t) + \underbrace{x}_{l}(t) \xi_{j}(t) \stackrel{\circ}{=} \xi(t) - \underline{m}_{\xi}(t)$$
(5,11)
$$(\underline{m}_{X} \infty) (\overline{\infty})$$

where $\underline{\mathbf{m}}_{\xi}(t) \cong \mathbf{M}[\underline{\xi}(t)]; \underline{\mathbf{m}}_{\xi} \cong \mathbf{M}(\underline{\xi}) = 0$, $\underline{\mathbb{C}}_{\xi\xi} \cong \mathbf{M}(\underline{\xi}\underline{\xi}^{*}) = \mathbf{M}\langle \xi_{j}^{2} \rangle \cong \underline{\mathbb{V}}_{\xi} \langle \mathbf{V}_{\xi} = \|\mathbf{J}\sigma_{j}^{2} > 0 \rangle$, so $\underline{\xi}(t)$ consist of incovariant components (white noises) $\underline{\mathbf{x}}_{1}(t) \boldsymbol{\xi}_{1}$ and has the covariance functions: $\underline{\mathbb{C}}_{\xi\xi}$, $(t,t') = \underline{\mathbf{x}}(t) \underline{\mathbb{V}}_{\xi} \overline{\underline{\mathbf{x}}}^{*}(t')$, $\underline{\mathbf{x}}(t) = \underline{\mathbb{C}}_{\xi}(t) \underline{\mathbb{V}}_{\xi}$. - One uses it in an <u>finite</u> (approximate) form (but with former $\underline{\mathbf{x}}_{1}(t)$).

$$\underbrace{\xi(t)}_{(m)} = \underbrace{\Sigma}_{1=1} \underbrace{X}_{1}(t) \xi_{1} + \underbrace{\rho}_{\mu}(t) = \underbrace{X}_{(t)}(t) \underbrace{\xi}_{\xi} + \underbrace{\rho}_{\mu}(t) \approx \underbrace{X}_{1}(t) \underbrace{\xi}_{\mu}(t) \underbrace{\xi}_{\mu} = \underbrace{V}_{\mu}(t) \\ (m \times_{\mu})(\mu) - \underbrace{X}_{\mu}(t) \underbrace{Y}_{\mu} \underbrace{\overline{X}}^{\mu}(t)$$

and the suitable random vector ξ can be realized e.g. by our algorithms ITA (probable) or OTA (statistical) [8], e.g. giving for $\underline{T}_{\mu} = \frac{1}{1} \{ t_{\lambda} (\langle t_{\lambda+\mu} \} the exact (sample) values \xi (T_{\mu}) = \xi X_{\lambda} (T_{\mu}), \dots (5, 15) but for t \in T - T_{\mu}$ the approximates ones only.

5,2. This method can be used also at our motor vehicle problem of

4,2, namely with the Rbp-form (5,11) and with (4,22), its stoch. one:

$$\underline{\zeta}_{\mathbf{n}}(t) = \int_{-\infty}^{t} e^{\underline{A}(t-\tau)} \underline{B}[\underline{m}_{\xi}(\tau) + \underline{X}(\tau)\underline{\xi}] d\tau = m_{\zeta}(t) + \underline{Z}(t)\underline{\xi}. \quad (5,21a,b)$$

From our [10], its covariance is for a general and stac. & ergodic case: \$T\$

$$\underline{\underline{C}}_{\xi\xi}, (t,t') = \underline{\underline{Z}}(t) \underline{\underline{V}}_{\xi} \overline{\underline{Z}}^{*}(t'), \ \underline{\underline{C}}_{\xi\xi}, (\tau) = \frac{1}{2T_{\infty}} \int_{-T_{\infty}}^{\infty} \underline{\underline{C}}(t) \underline{\underline{C}}^{*}(t+\tau) dt, \quad (5,22)$$

then the spectral density matrix and its inverse, by Fourier- $\ensuremath{ \int }$ &-inverse

$$\begin{split} \underbrace{S}_{\xi\zeta}(\omega) &= \int_{-\infty}^{+\infty} e^{-i\omega\tau} \underbrace{C}_{\xi\xi\zeta}(\tau) d\tau, \\ \underbrace{C}_{\xi\xi\zeta}(\tau) &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{i\tau} \underbrace{S}_{\xi\zeta}(\omega) d\omega, \\ \text{whose approximate form (at } \tau = 0) \\ \underbrace{C}_{\xi\xi\zeta}(0) &\approx \frac{1}{2\pi} \int_{u_1}^{u_2} \underbrace{S}_{\xi\zeta}(\omega) d\omega, \\ \underbrace{S}_{\xi\zeta}(\omega) d\omega,$$

5,3. At the end, let us mention the Markov-chains treated by matrix analysis in our bulletin [11] with problems of mass service, demography, random walk etc., then our investigations [13] on parametrical and noisy Gaussian process and white noise, which promis an advance at the optimalization of noisy control systems.

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