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# TWISTOR SPINORS AND SOLUTIONS OF THE EQUATION (E) ON RIEMANNIAN MANIFOLDS 

by
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## 1. Introduction

Let $M^{n}$ be a Riemannian spin manifold and denote by $R$ its scalar curvature. The conformally invariant twistor operator $\mathcal{D}$ acting on sections $\psi$ of the spinor bundle $S$ is defined by the covariant derivative and the projection onto the kernel of the Clifford multiplication. The kernel $\operatorname{Ker}(\mathscr{D})$ of this operator is the space of all spinor fields $\psi$ satisfying the first order differential equation

$$
\nabla_{X} \psi+\frac{1}{n} x \cdot D \psi=0
$$

where $D$ denotes the Dirac operator (see [3],[4]). Killing spinors, i.e. spinor fields satisfying for some complex number $\lambda \in \mathbb{C}$ the equation

$$
\nabla_{x} \psi=\lambda x \cdot \psi
$$

are special solutions of the twistor equation. A. Lichnerowicz (see [5]) studied the so-called equation (E) for spinor fields:

$$
\nabla_{X}(D \psi)+\frac{R}{4(n-1)} X \cdot \psi=0 .
$$

In particular, A. Lichnerowicz proved that if a connected Riemannian manifold admits a non-trivial solution of the equation ( $E$ ) then its scalar curvature $R$ is constant. Consequently, in case of a compact Riemannian manifold the space of all twistor spinors coincides with $\operatorname{Ker}(E)$ and with the space.
of all Killing spinors (see [2]). The aim of this note is to compare the kernel of the twistor operator with the space $\operatorname{Ker}(E)$ of all solutions of equation (E).

## 2. A relation between $\operatorname{Ker}(E)$ and $\operatorname{Ker}(D)$.

We consider a Riemannian spin manifold ( $M^{n}, g$ ) with a nontrivial solution of equation (E). Then its scalar curvature $R$ is constant.

Proposition 1:

$$
\operatorname{Ker}(E)=\operatorname{Ker}\left(D^{2}-\frac{n R}{4(n-1)}\right) \cap D^{-1}(\operatorname{Ker}(D))
$$

Proof: If $\psi \in \operatorname{Ker}(E)$ we obtain from equation $\langle 3\rangle$

$$
D^{2} \psi=\sum_{i=1}^{n} e_{i} \cdot \nabla_{e_{i}}(D \psi)=-\sum_{i=1}^{n} e_{i} \cdot \frac{R}{4(n-1)} e_{i} \cdot \psi=\frac{n R}{4(n-1)} \psi,
$$

i.e. $\psi$ belongs to the kernel of the operator $D^{2}-\frac{n R}{4(n-1)}$.

Furthermore, we have

$$
\begin{aligned}
& \nabla_{x}(D \psi)+\frac{1}{n} x \cdot D(D \psi)=\bar{V}_{X}(D \psi)+\frac{1}{n} x \cdot \frac{n R}{4(n-1)} \psi= \\
& \quad=\nabla_{X}(D \psi)+\frac{R}{4(n-1)} x \cdot \psi=0
\end{aligned}
$$

and, consequently, $D \psi$ is a twistor spinor. Conversely, if $\psi \in \operatorname{Ker}\left(D^{2}-\frac{n R}{4(n-1)}\right) \cap D^{-1}(\operatorname{Ker}(D))$ then we have $D^{2} \psi=\frac{n R}{4(n-1)} \psi$ as well as $\nabla_{X}(D \psi)+\frac{1}{n} X \cdot D(D \psi)=0$ and $\psi$ belongs to $\operatorname{Ker(E).}$

Theorem 1: Let ( $M^{n}, g$ ) be a Riemannian spin manifold with constant scalar curvature $R \neq 0$. Then the map

$$
\operatorname{Ker}(E) \partial \psi \longrightarrow D \psi \in \operatorname{Ker}(D)
$$

is an isomorphism.
Proof: By proposition 1 the map is well-defined. If $D \psi=0$
and $\psi \in \operatorname{Ker}(E)$ the equation $\langle 3\rangle$ implies

$$
\frac{R}{4(n-1)} x \cdot \psi=0
$$

Since $R \neq 0$ we conclude $\psi=0$, ide. the map $\operatorname{Ker}(E) \rightarrow \operatorname{Ker}(\mathscr{\infty})$
is infective. On the other hand, given a twistor spinor $\varphi \in \operatorname{Ker}(D)$ we consider.

$$
\psi=\frac{4(n-1)}{n R} D \varphi .
$$

Then $D \psi=\frac{4(n-1)}{n R} D^{2} \varphi=\frac{4(n-1)}{n R} \frac{n R}{4(n-1)} \varphi=\varphi$
and $\nabla_{X}(D \psi)+\frac{R}{4\left(\frac{n-1)}{} X \cdot \psi=\right.}$
$=\nabla_{X}\left(\frac{4(n-1)}{n R} D^{2} \varphi\right)+\frac{1}{n} x \cdot D \varphi=$
$=\nabla_{X}(\varphi)+\frac{1}{n} X \cdot D \varphi=0$.

This means that $\psi$ belongs to $\operatorname{Ker}(E)$ and is the preimage of $\varphi$, i.e. the map $\operatorname{Ker}(E) \rightarrow \operatorname{Ker}(\underset{\infty}{ })$ is surjective.

Corollary: Let ( $M^{n}, g$ ) be a connected Riemannian manifold with scalar curvature $R \neq 0$. Then $\left[\begin{array}{l}n \\ \frac{2}{2}\end{array}\right]+1$

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dim Ker(E) = dim Ker (\mathcal{D }\leqslant2
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Proof: see [3].

Proposition 2: If $\psi \in \operatorname{Ker}(E)$ and $R \neq 0$ then

$$
\nabla_{X} \psi=\frac{2(n-1)}{R(n-2)}\left(\frac{R}{2(n-1)} x-R i c(x)\right) \cdot D \psi
$$

Proof: If $\psi$ belongs to Ker(E) then by proposition 1 $\mathcal{L} \psi$ is a twister spinor. The general Lichnerowicz formula

$$
\begin{aligned}
& \nabla_{X}(D \varphi)=\frac{n}{2\left(\frac{n-2)}{(n-1)}\right.}\left(\frac{R}{2(n \cdot \varphi-R i c(x) \cdot \varphi)}\right. \\
& \text { valid for any twistor spinor } \varphi(\operatorname{see}[4]) \text { now yields } \\
& \nabla_{X}\left(D^{2} \psi\right)=\frac{n}{2(n-2)}\left(\frac{R}{2(n-1)} x-R i c(X)\right) \cdot D \psi \cdot
\end{aligned}
$$

Moreover, since $\psi \in \operatorname{Ker}(E)$ it satisfies the equation

$$
D^{2} \psi=\frac{n R}{4(n-1)} \psi
$$

and we obtain

$$
\frac{n R}{4(n-1)} \nabla x \psi=\frac{n}{2(n-2)}\left(\frac{R}{2(n-1)} x-\operatorname{Ric}(x)\right) \cdot D \psi
$$

## Theorem 2:

a) Let ( $M^{n}, g$ ) be an Einstein manifold with scalar curvature $R \neq 0$. Then $\operatorname{Ker}(E)$ and $\operatorname{Ker}(\mathcal{D})$ coincide, i.e.

$$
\operatorname{Ker}(E)=\operatorname{Ker}(D)
$$

b) Let $\left(M^{n}, g\right)$ be a connected Riemannian manifold such that $\operatorname{Ker}(E) \cap \operatorname{Ker}(, Q) \neq\{0\}$. Then $\left(M^{n}, g\right)$ is an Einstein space.

Proof: Suppose first that $M^{n}$ is an Einstein space, $\operatorname{Ric}(X)=\frac{R}{n} X$. If $\psi \in \operatorname{Ker}(E)$ we obtain from proposition 2
$\nabla_{X} \psi=\frac{2(n-1)}{R(n-2)}\left(\frac{R}{2(n-1)} x-R i c(x)\right) \cdot D \psi=-\frac{1}{n} x \cdot D \psi$
and $\psi$ is a twistor spinor. Conversely, if $\psi \in \operatorname{Ker}(D)$ we use the Lichnerowicz formula 〈4〉

$$
\nabla_{x}(D \psi)=\frac{n}{2(n-2)}\left(\frac{R}{2(n-1)} x-R i c(x)\right) \cdot \psi
$$

which reduces in an Einstein space to

$$
\nabla_{X}(\mathrm{D} \varphi)+\frac{\mathrm{R}}{4(\mathrm{n}-1)} \mathrm{x} \cdot \psi=0
$$

This means that in an Einstein space every twistor spinor is a solution of equation (E). We consider now an arbitrary Riemannian manifold as well as a non-trivial solution $\psi \in \operatorname{Ker}(E) \cap \operatorname{Ker}(D)$. Using the formulas $\langle 3\rangle$ and $\langle 4\rangle$ we obtain the condition

$$
-\frac{R}{4(n-1)} x \cdot \psi=\frac{n}{2(n-2)}\left(\frac{R}{2(n-1)} x-\operatorname{Ric}(x)\right) \cdot \psi
$$

and, finally,

$$
\operatorname{Ric}(x) \cdot \psi=\frac{R}{n} x \cdot \psi
$$

Since $\psi$ is a twistor spinor, the zeroes of $\psi$ are isolated points (see [3]) and therefore we conclude

$$
\operatorname{Ric}(x)=\frac{R}{n} x,
$$

i.e. $M^{n}$ is an Einstein space.

## 3. An example

Let $M^{2}$ be the simply connected Riemannian surface with constant scalar curvature $R$. Then $M^{2}$ admits two Killing spinors $\varphi$ and $\psi$. According to the decomposition of the spinor bundle $\mathrm{S}=\mathrm{S}^{+} \oplus \mathrm{S}^{-}$on $\mathrm{M}^{2}$ we decompose the Killing spinors into $\psi=\psi^{+}+\psi^{-}, \varphi=\varphi^{+}+\varphi^{-}$. We consider the three-dimensional Riemannian manifold $M^{3}=M^{2} \times R^{1}$. The general solution of equation ( $E$ ) on $M^{3}=H^{2} \times R^{1} \quad$ (case $R<0$ ) is given by

$$
\begin{aligned}
\psi(x, t) & =\left\{A_{0} \cos \left(\frac{1}{2} t\right)+A_{1} \sin \left(\frac{1}{2} t\right)\right\} \psi^{+}(x)+ \\
& +\left\{A_{0} \sin \left(\frac{1}{2} t\right)-A_{1} \cos \left(\frac{1}{2} t\right)\right\} \psi^{-}(x)+ \\
& +\left\{B_{0} \cos \left(\frac{1}{2} t\right)+B_{1} \sin \left(\frac{1}{2} t\right)\right\} \varphi^{+}(x)+ \\
& +\left\{-B_{0} \sin \left(\frac{1}{2} t\right)+B_{1} \cos \left(\frac{1}{2} t\right)\right\} \varphi^{-}(x) .
\end{aligned}
$$

A similar formula we can get on $M^{3}=S^{2} \times R^{1} \quad$ (case $R>0$ ):

$$
\begin{aligned}
\psi(x, t) & =\left\{A_{0} \cosh \left(\frac{1}{2} t\right)+A_{1} \sinh \left(\frac{1}{2} t\right)\right\} \psi^{+}(x)- \\
& -i\left\{A_{0} \sinh \left(\frac{1}{2} t\right)+A_{1} \cosh \left(\frac{1}{2} t\right)\right\} \psi^{-}(x)+ \\
& +\left\{B_{o} \cosh \left(\frac{1}{2} t\right)+B_{1} \sinh \left(\frac{1}{2} t\right)\right\} \varphi^{+}(x)+ \\
& +i\left\{B_{0} \sinh \left(\frac{1}{2} t\right)+B_{1} \cosh \left(\frac{1}{2} t\right)\right\} \varphi^{-}(x) .
\end{aligned}
$$

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